

Precision Cosmology in the E-ELT Era

Jochen Liske



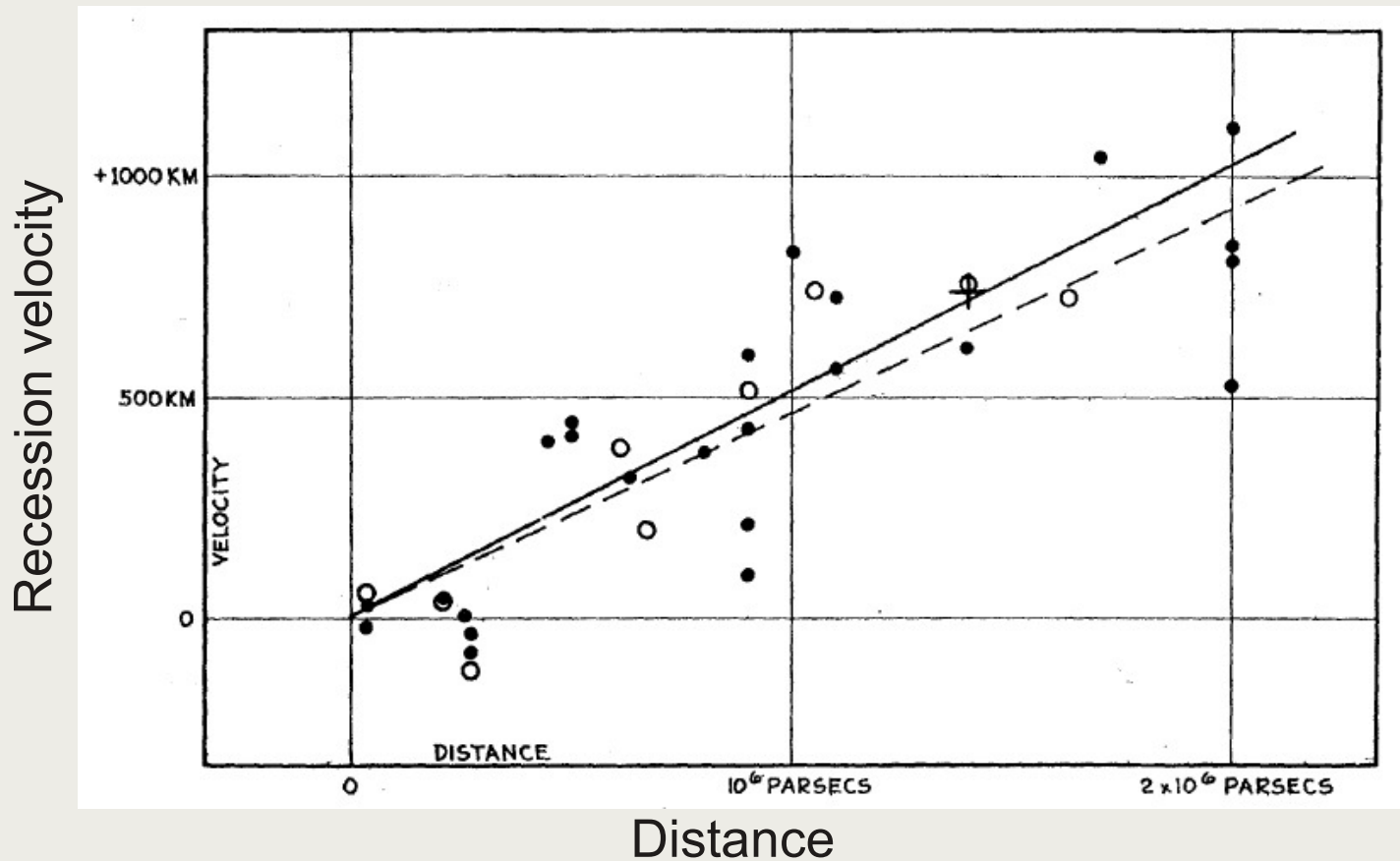
Precision Measurements for not very precise Cosmology in the E-ELT Era

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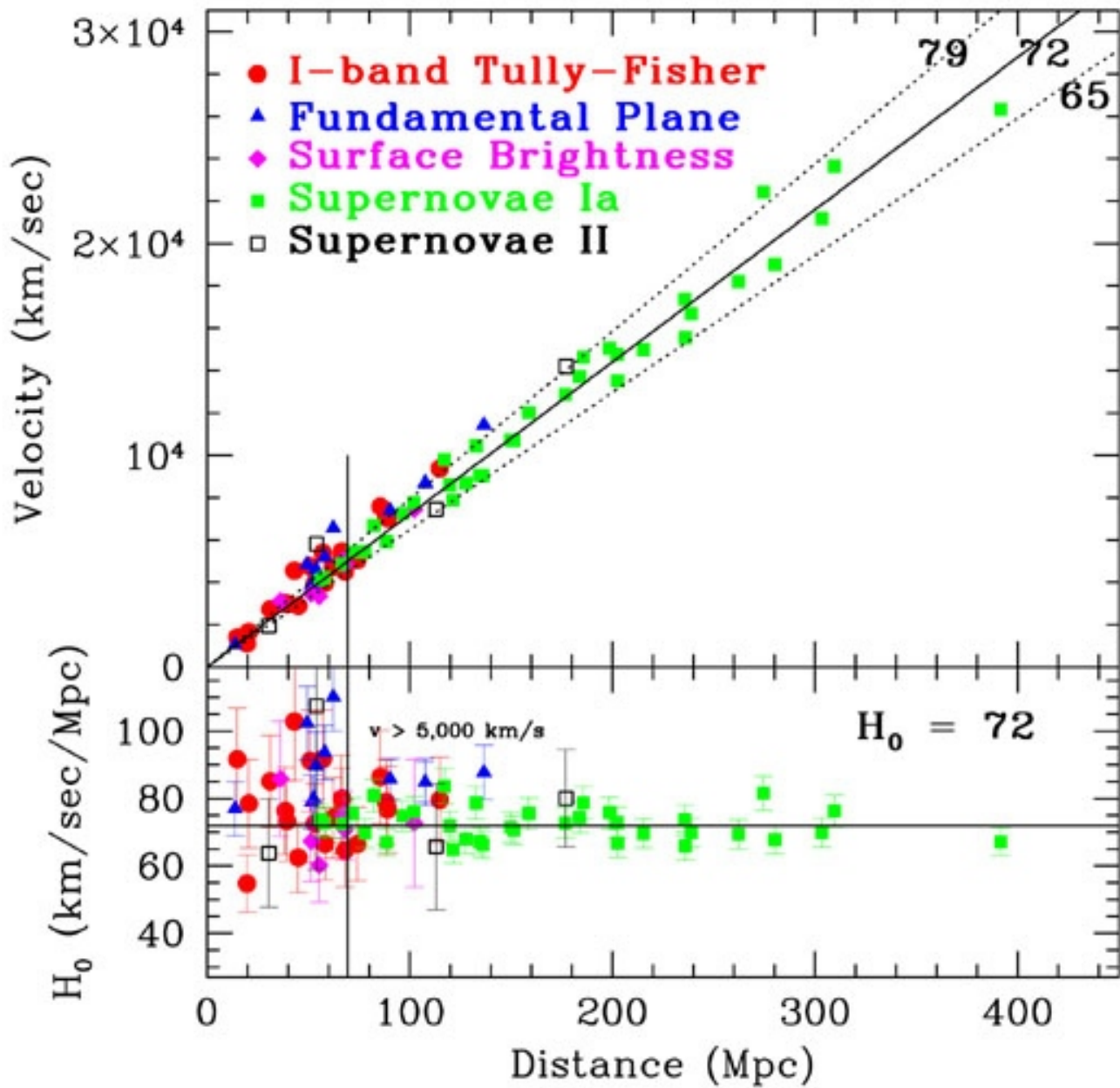


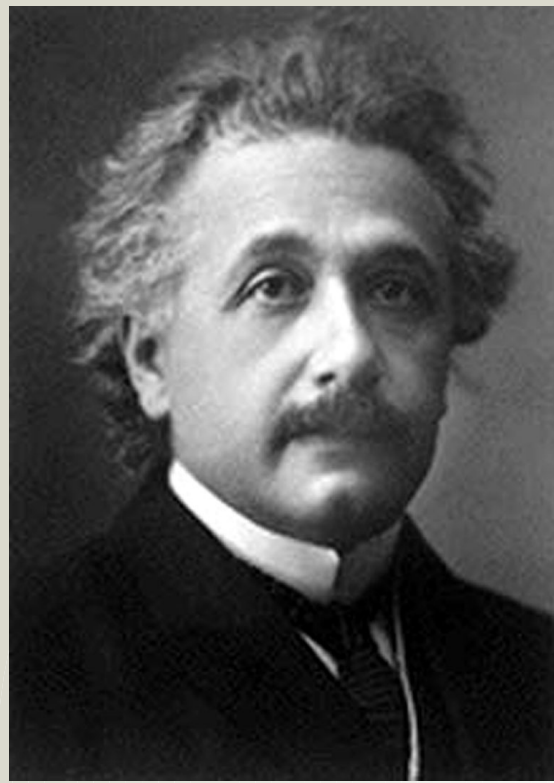
1929: Universal Expansion

The original Hubble diagram (Hubble 1929):



All distant galaxies are found to recede from us.
Hubble's Law: $v = H_0 d$ → The Universe expands!

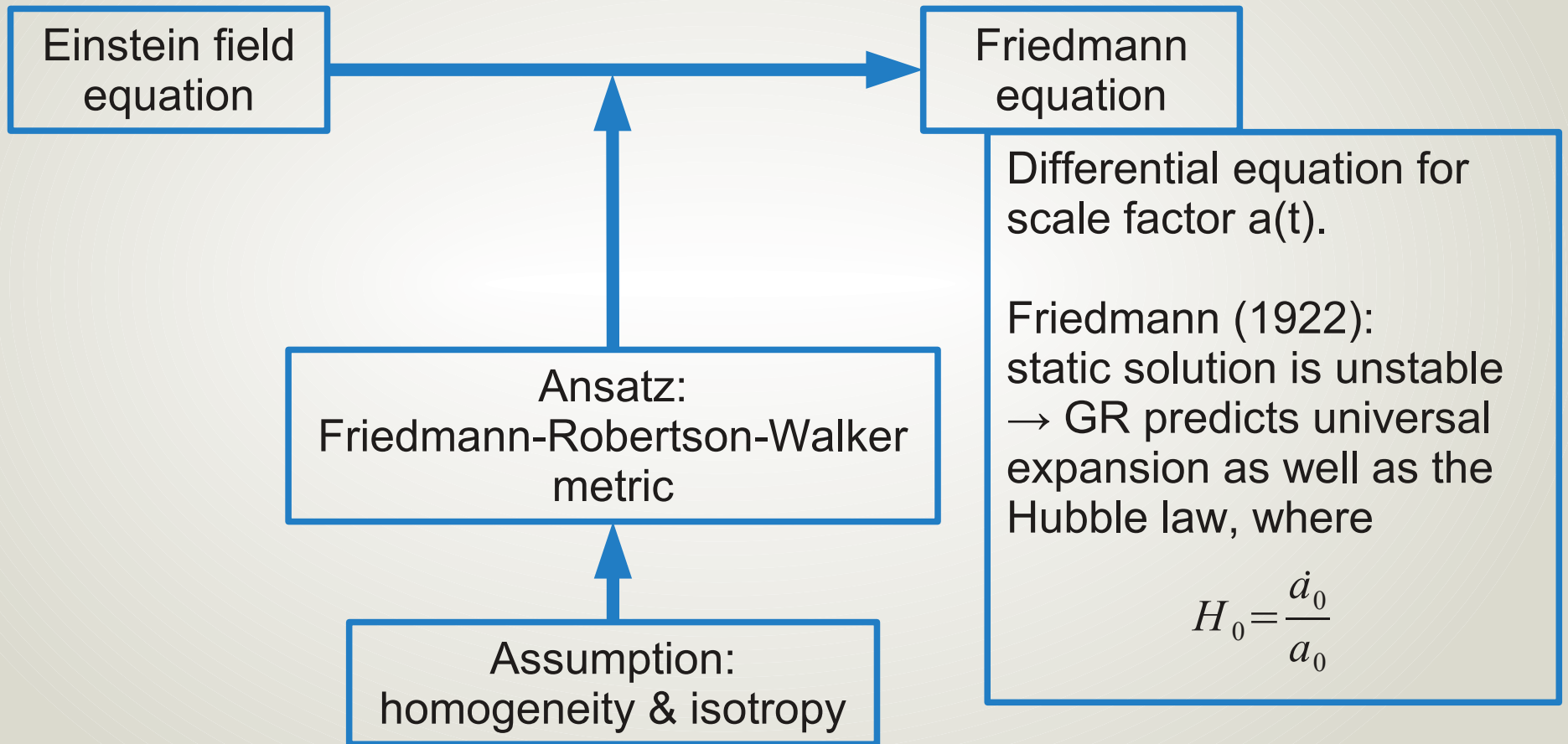




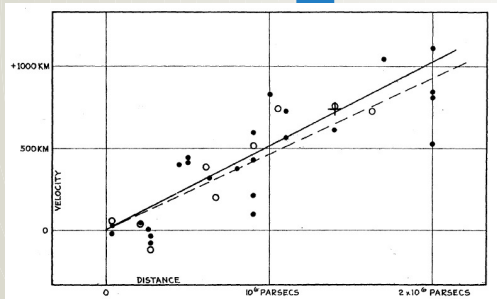
1916: Einstein's Theory of General Relativity



Relativistic Cosmology

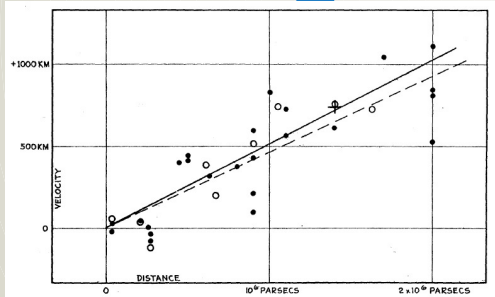


Relativistic Big Bang Cosmology

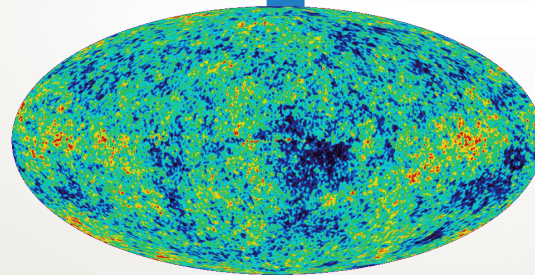


Expansion

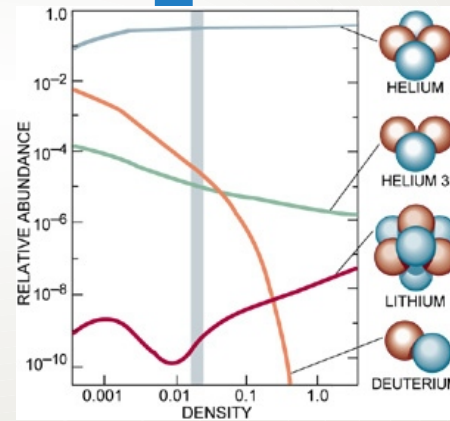
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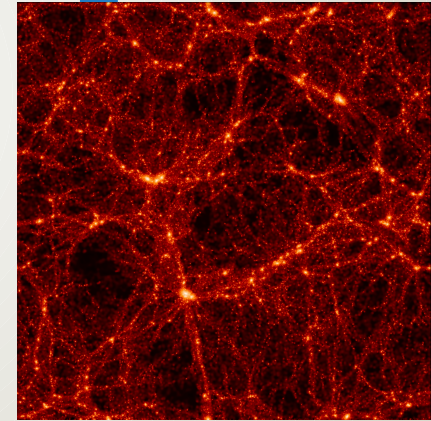
Expansion



Cosmic
Microwave
Background



Abundance
of light
elements



Structure
formation

Relativistic Cosmology

FRW metric:

$$ds^2 = -c^2 dt^2 + a^2(t)[d\chi^2 + \Sigma^2(\chi)(d\theta^2 + \sin^2 \theta d\phi^2)]$$

$$\Sigma(\chi) = \begin{cases} \sin \chi & k = +1 \\ \chi & k = 0 \\ \sinh \chi & k = -1 \end{cases}$$

Friedmann equation:

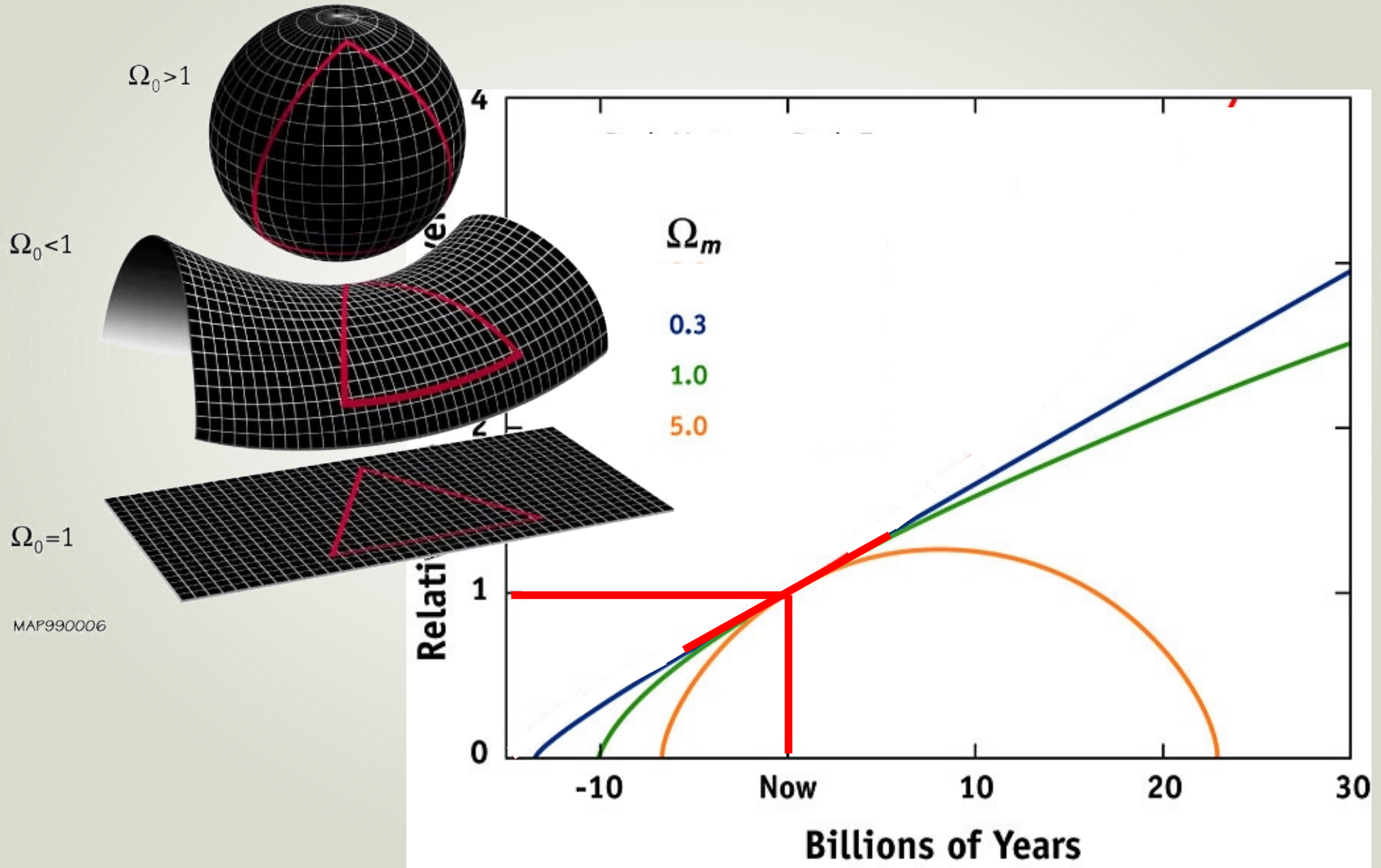
$$H(z) = H_0 \left[\sum_i \Omega_i (1+z)^{3(1+w_i)} + \Omega_k (1+z)^2 \right]^{\frac{1}{2}}$$

Equation of state:

$$p_i = w_i c^2 \rho_i$$

$$H = \frac{\dot{a}}{a}, \quad 1+z = \frac{a_0}{a}$$

Relativistic Cosmology



Which of the solutions of the Friedmann equation corresponds to reality?

$$G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

Or in other words:

What is the stress-energy tensor of the universe?

For each mass/energy component i , what is Ω_i , w_i (and what is H_0)?

Density parameter



Equation of state parameter



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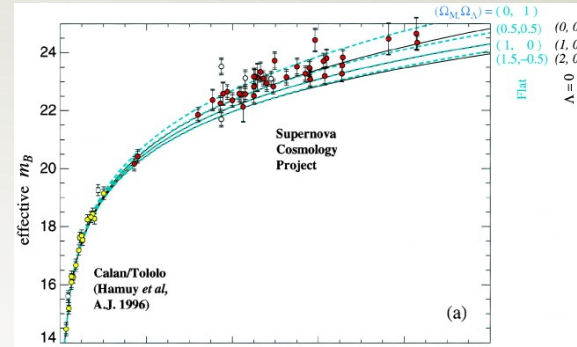
How can these be measured?

- Geometry
- Expansion history
- Clustering, evolution and dynamics of density perturbations

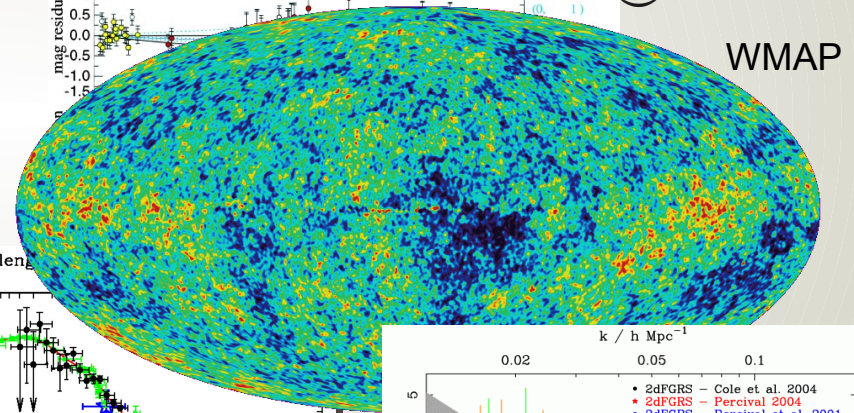
Precision Cosmology

Past decades: development of a wide array of observations to constrain the cosmological model:

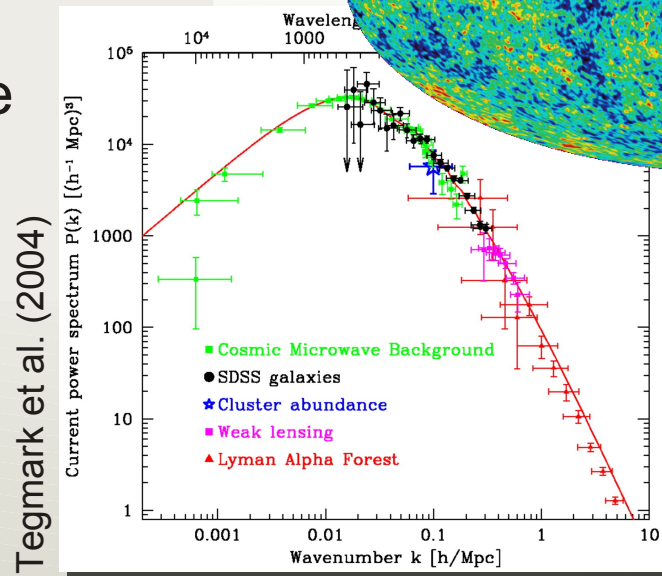
- Cosmic Microwave Background
- Supernovae type Ia
- Large scale structure of galaxies and intergalactic medium
- Galaxy cluster abundance
- Weak lensing



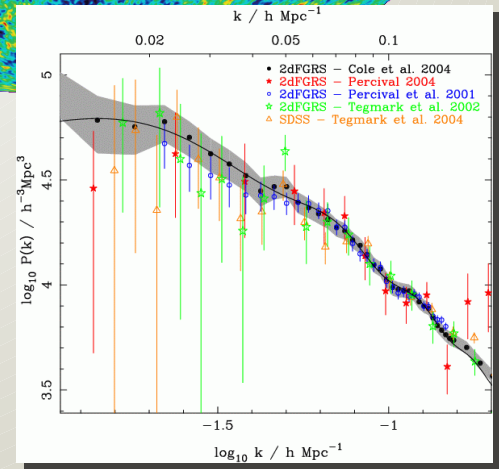
Perlmutter et al. (1999)



WMAP



Tegmark et al. (2004)



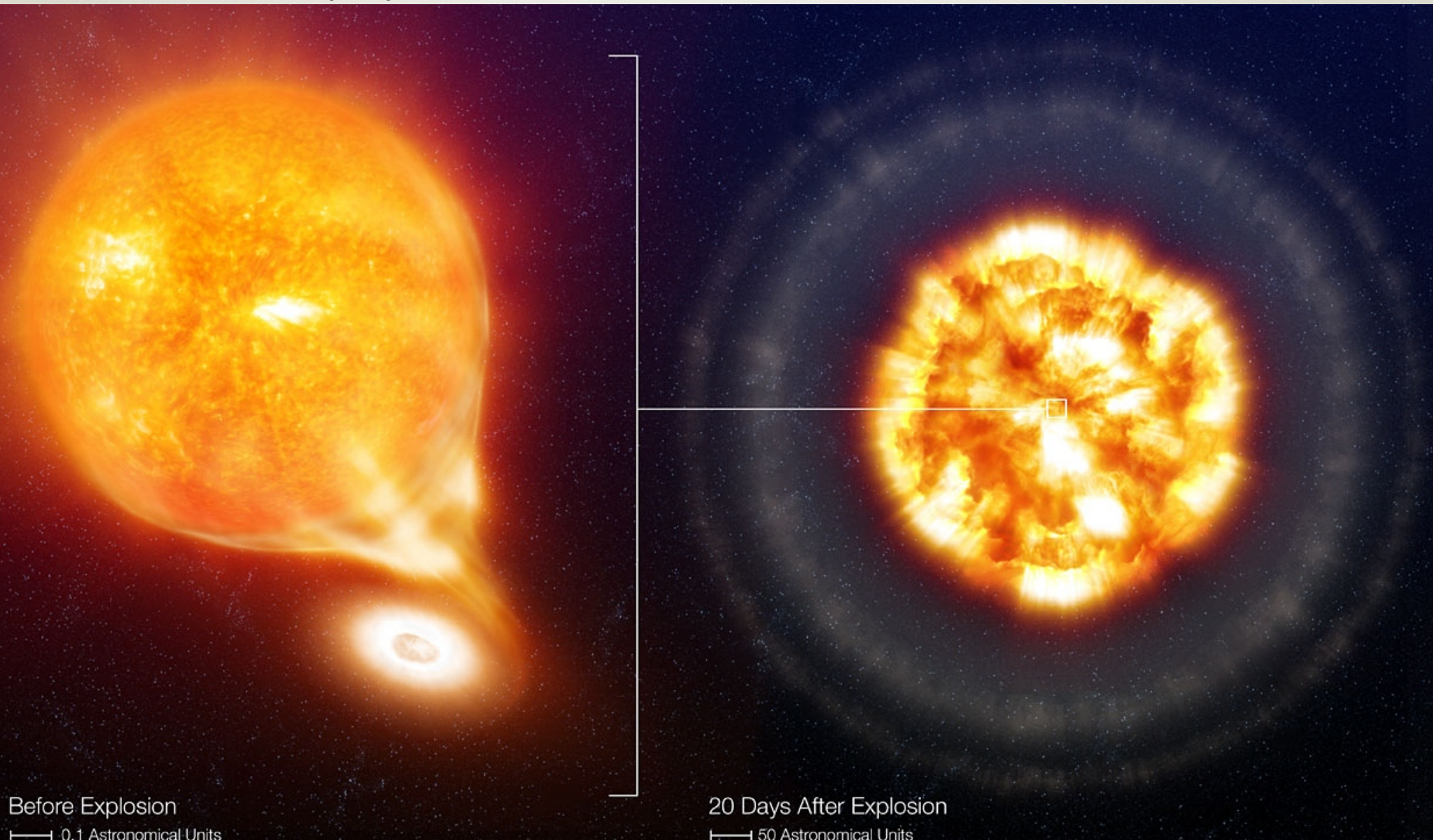
Cole et al. (2005)

Surprise: Accelerated Expansion

- Good evidence from SNIa that a period of decelerated expansion was followed 'recently' by a period of acceleration.

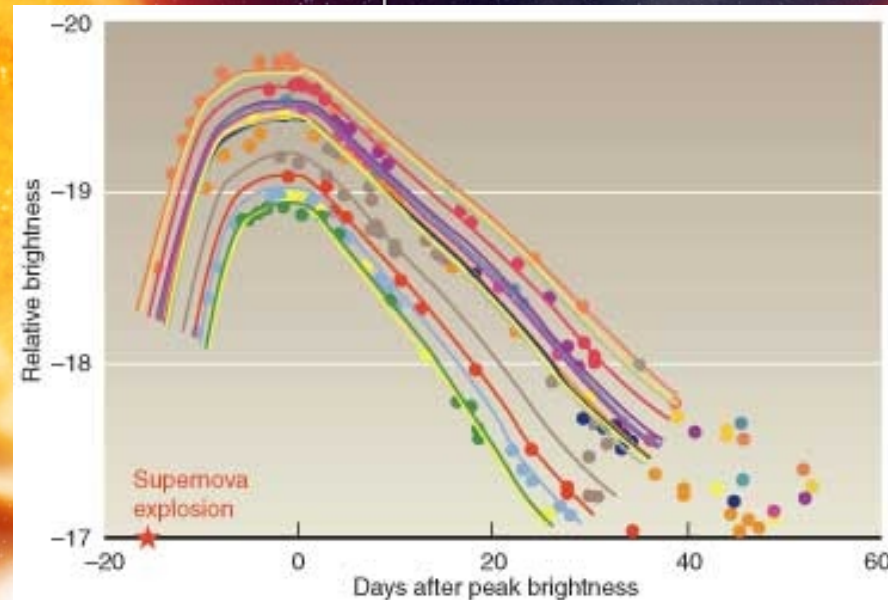
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Before Explosion

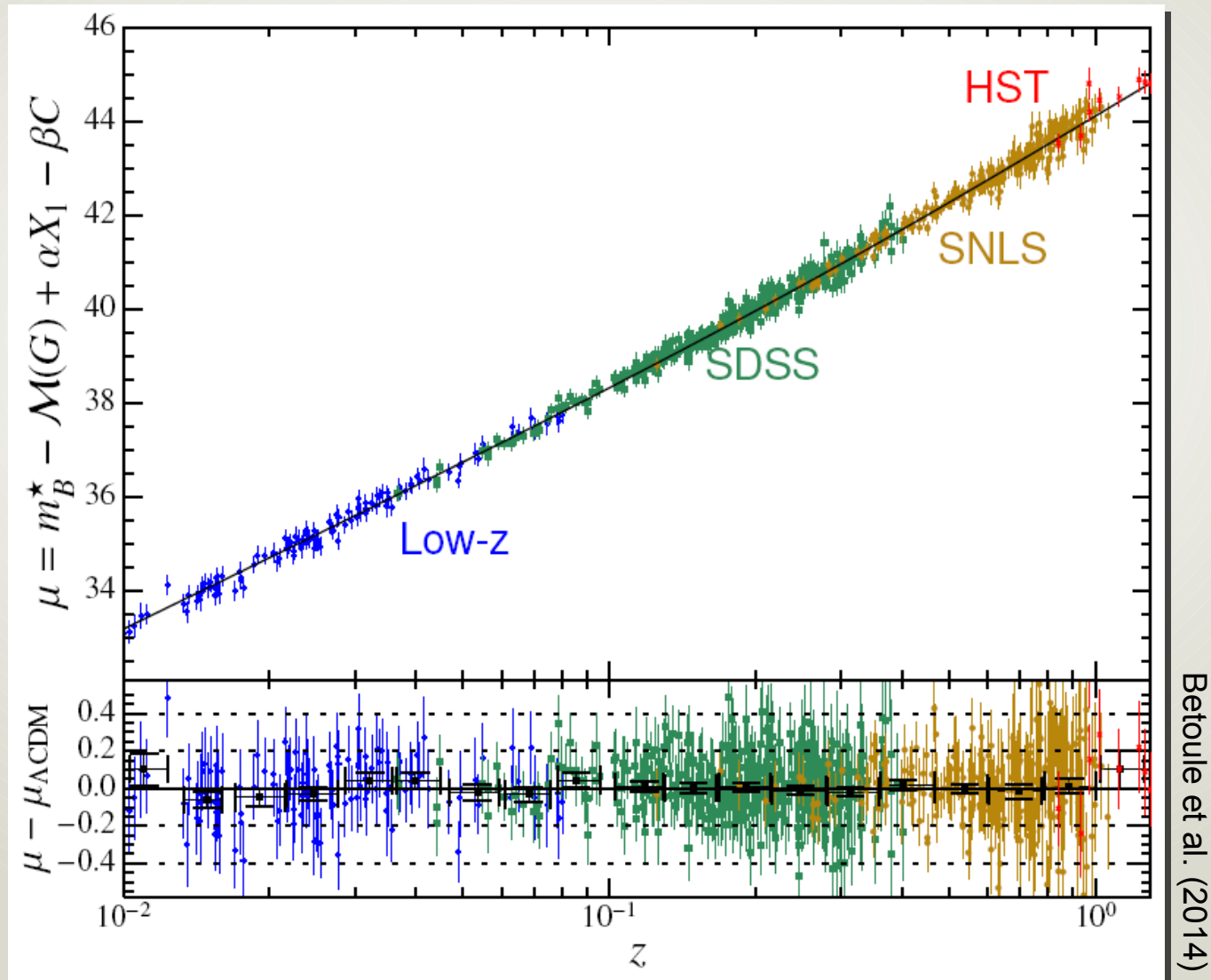
— 0.1 Astronomical Units

20 Days After Explosion

— 50 Astronomical Units

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Surprise: Accelerated Expansion

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- The source of the acceleration is entirely unknown. Most explanations so far proposed require new physics.

Dark energy:

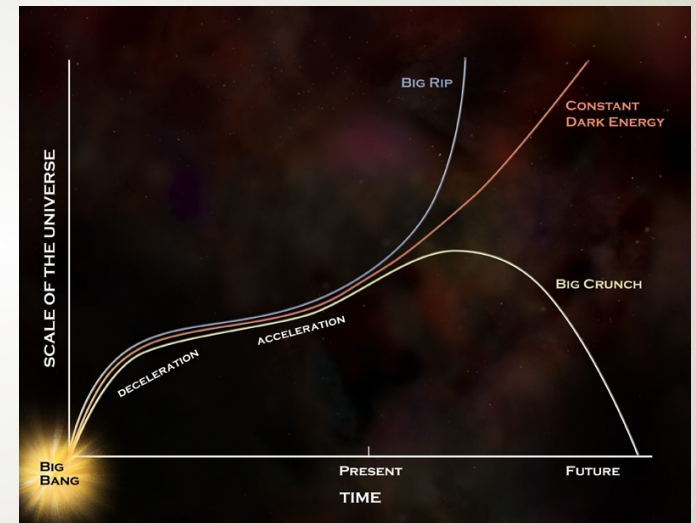
- Cosmological constant $w = -1$
- Quintessence $-1 < w(z) < 0$
- Phantom energy $w(z) < -1$
- ...

Modification of gravity:

- $f(R)$
- Non-minimal couplings
- Braneworld scenarios (DGP, Cardassian, ...)
- ...

Modification of Copernican Principle:

- Inhomogeneous models without DE can reproduce past light-cone observations of FRW models with DE (LTB, void models, ...)
- Backreaction (averaging and evolution do not commute)



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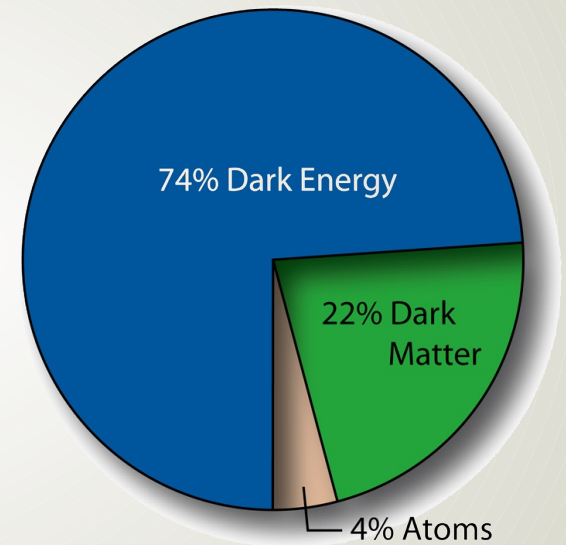
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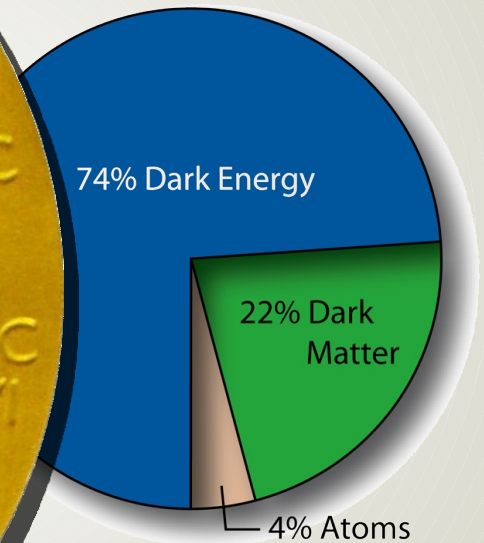
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Nobel Prize for Physics 2011



Saul Perlmutter



Brian Schmidt



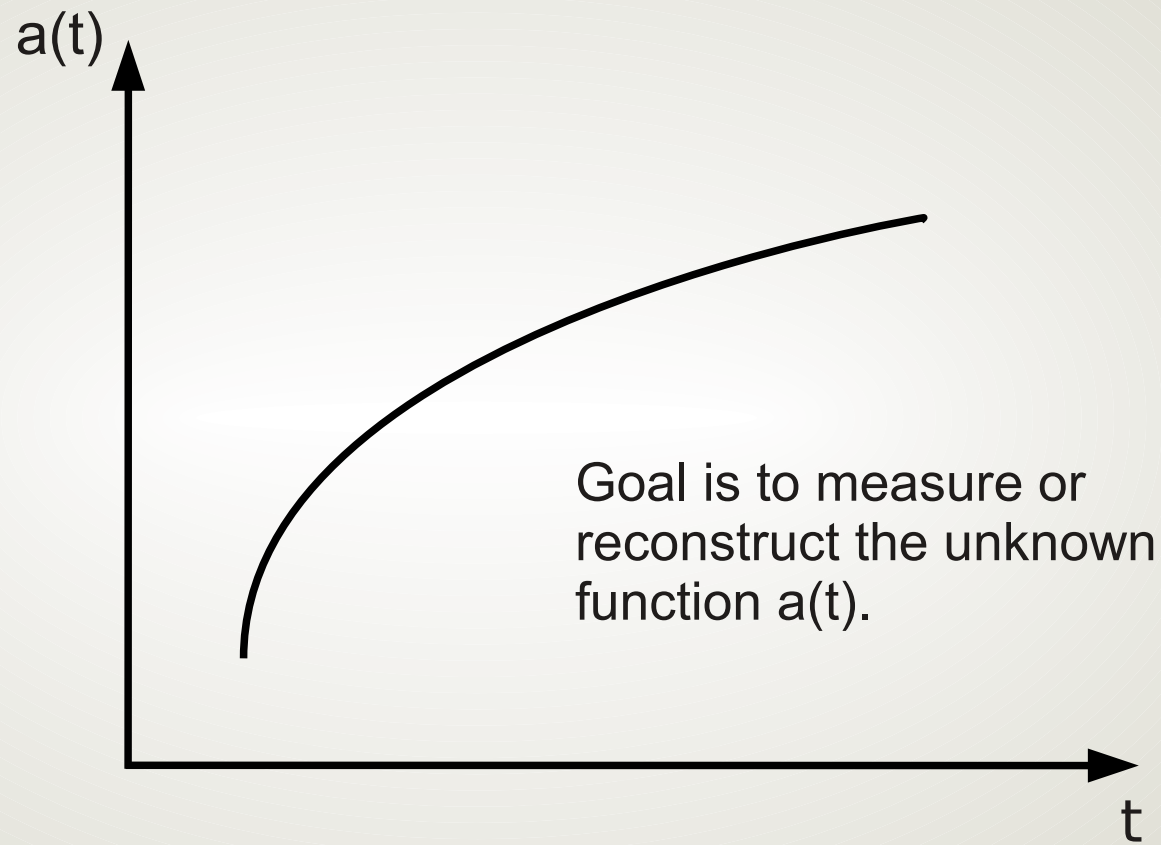
Adam Riess

Accelerated Expansion

→ Intense interest in the expansion history.

Accelerated Expansion

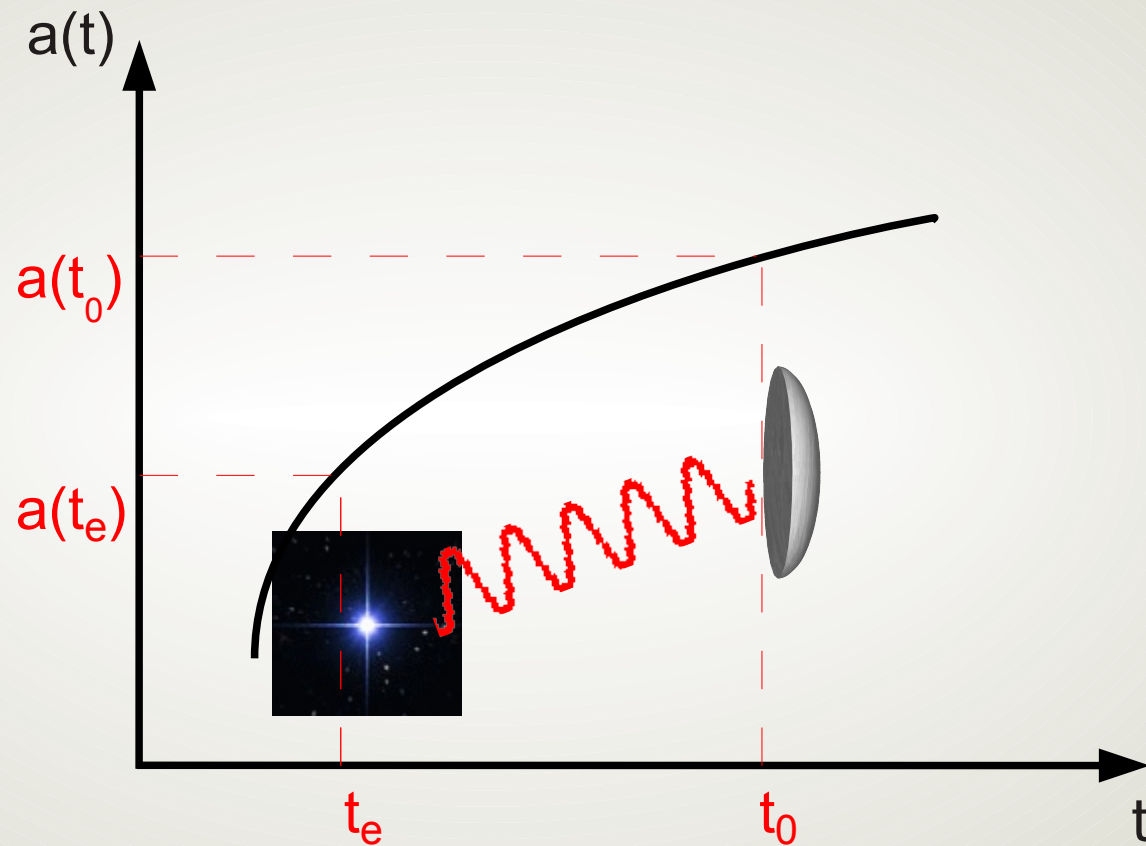
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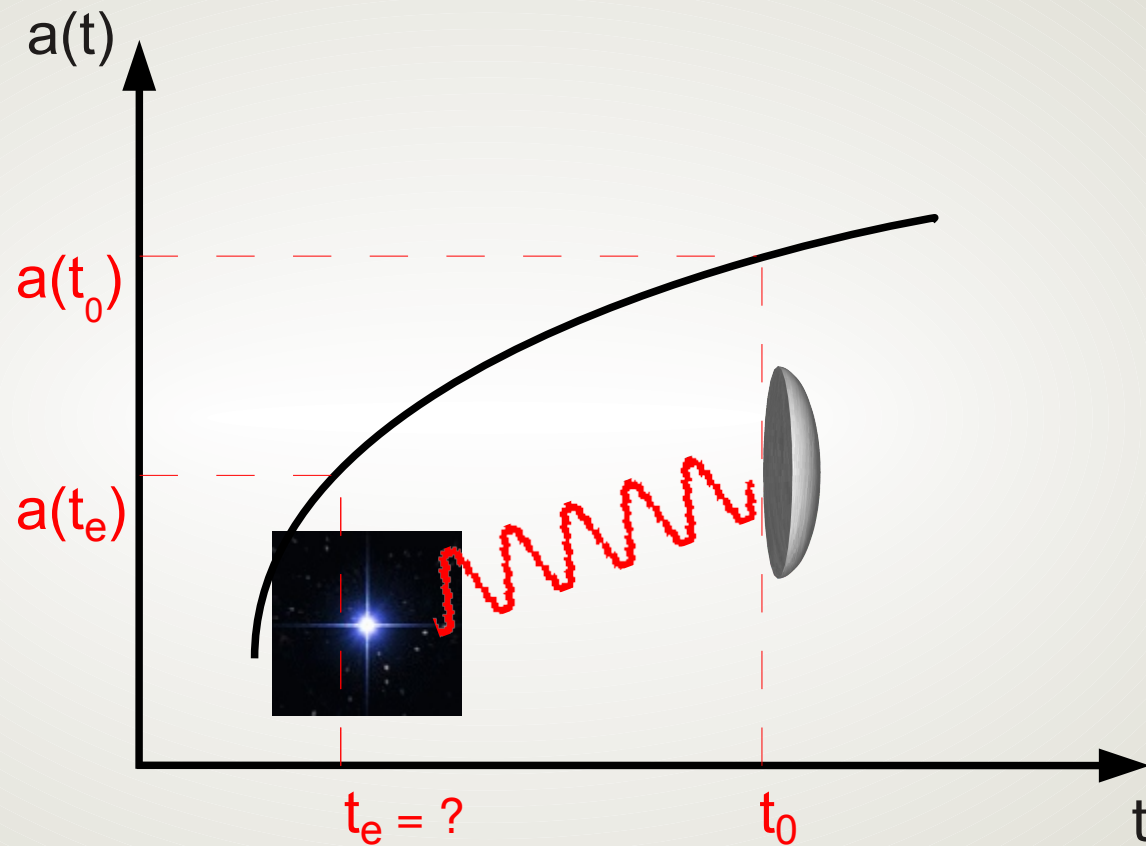
$$1+z = \frac{a(t_0)}{a(t_e)}$$



Accelerated Expansion

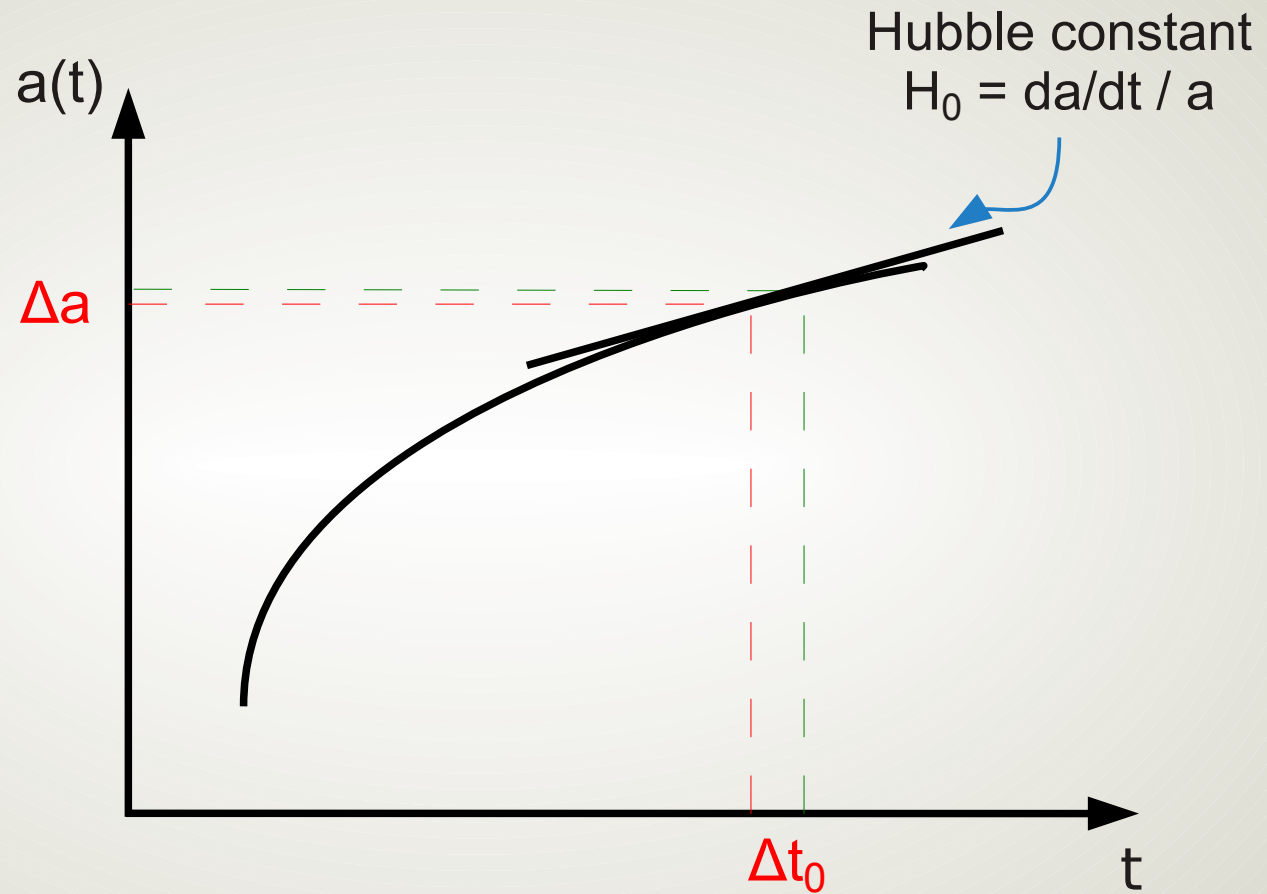
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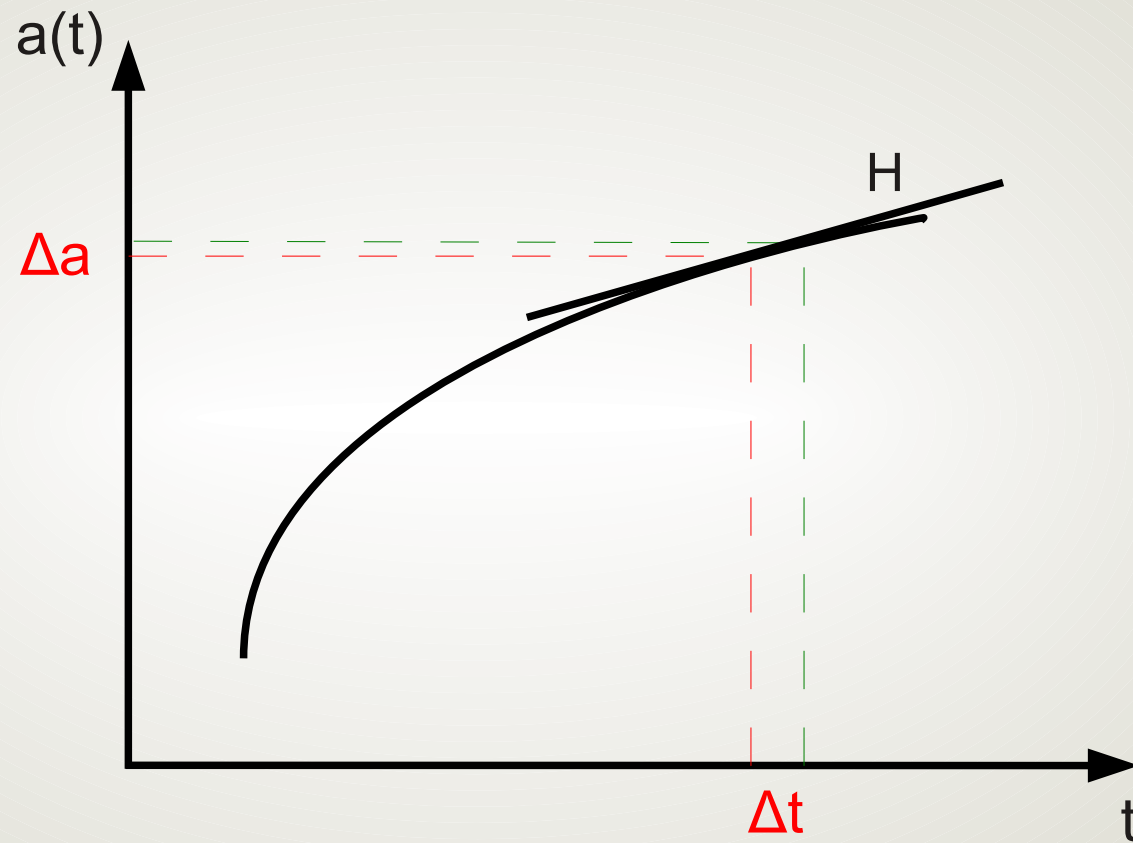
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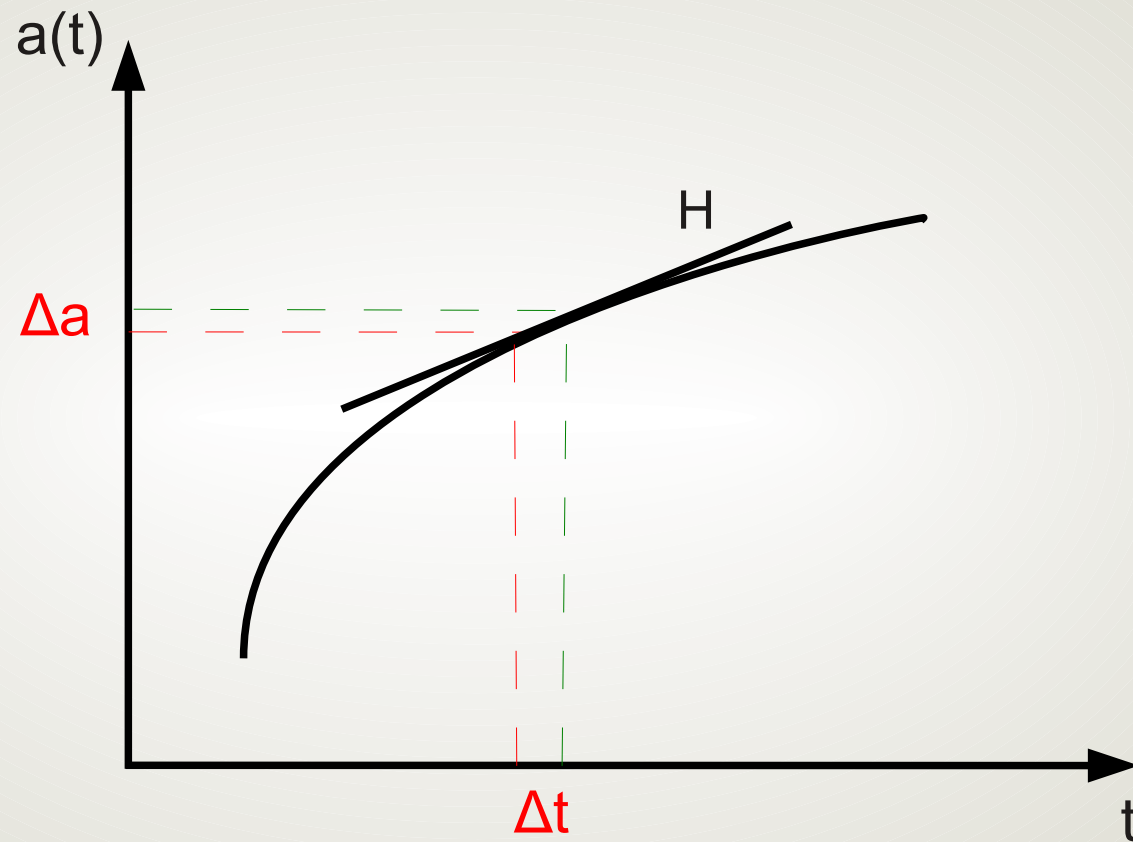
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A measurement of $H(z)$ allows the reconstruction of $a(t)$.

Accelerated Expansion

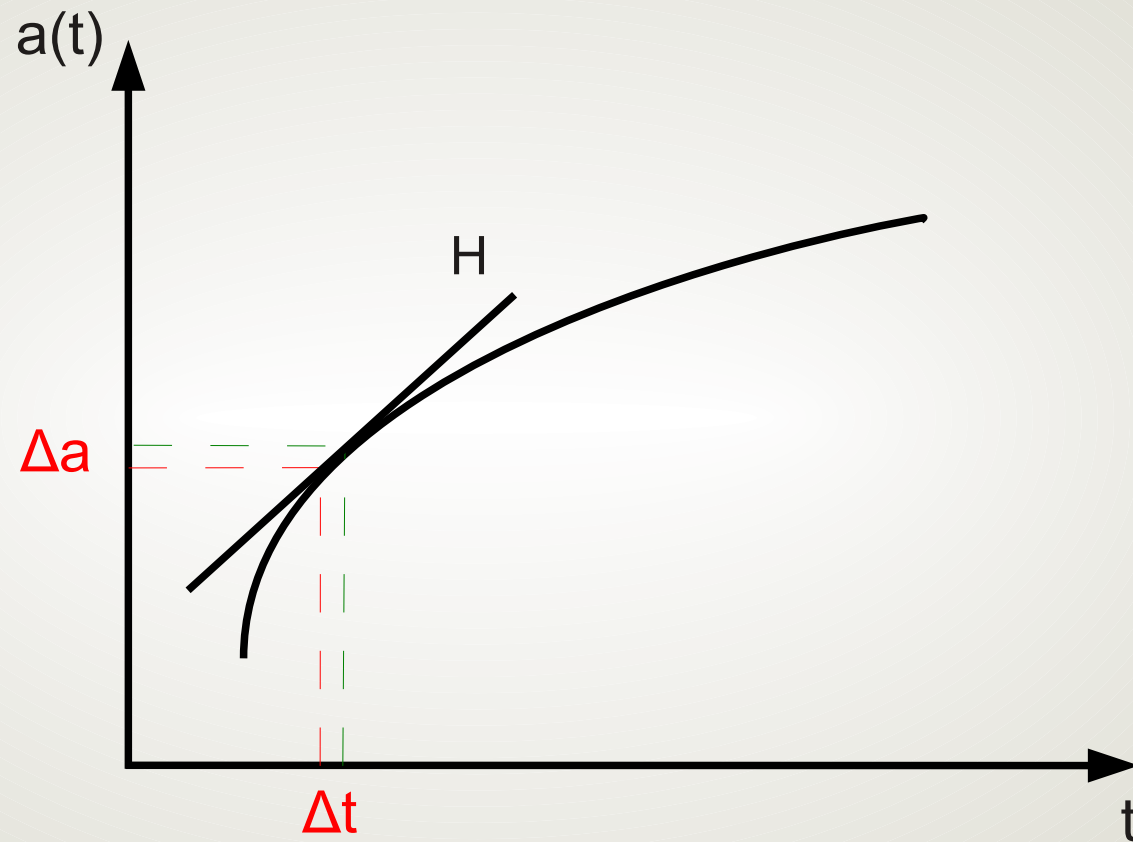
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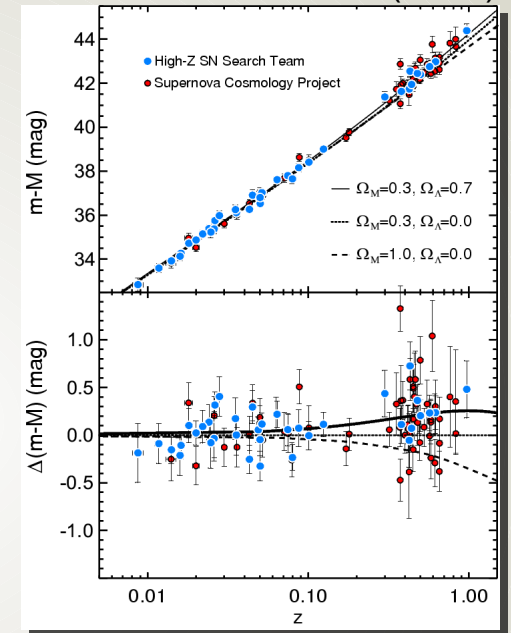
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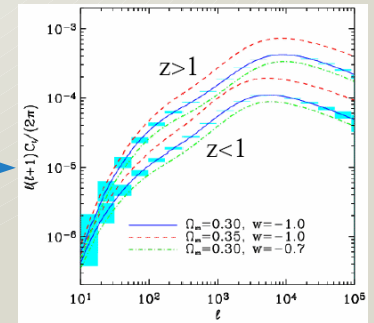
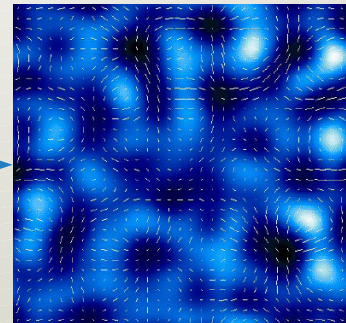
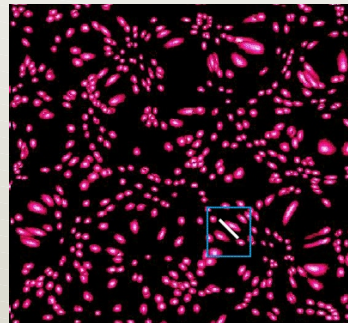
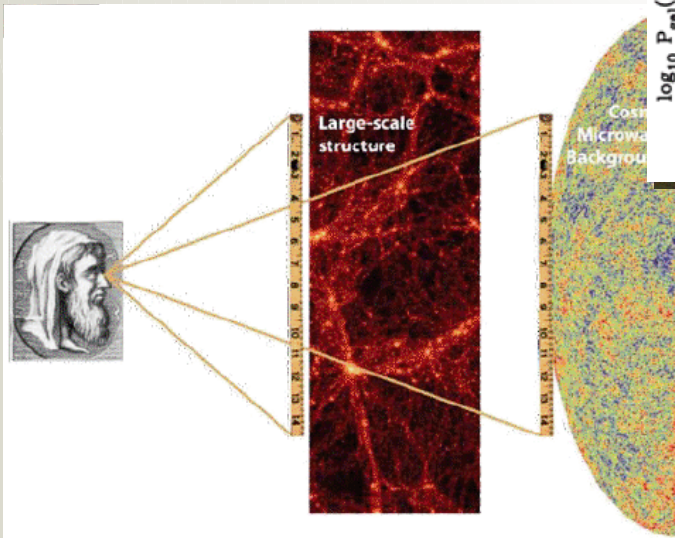
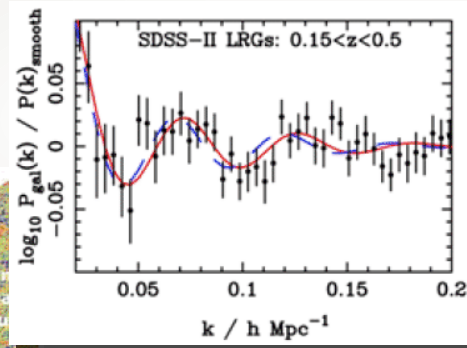
→ Intense interest in the expansion history.
Best current methods of measuring $H(z)$:

- SNIa
- Weak lensing
- Baryon Acoustic Oscillations (BAO)
- Redshift Space Distortions (RSD)

Perlmutter & Schmidt (2003)



Percival et al. (2010)

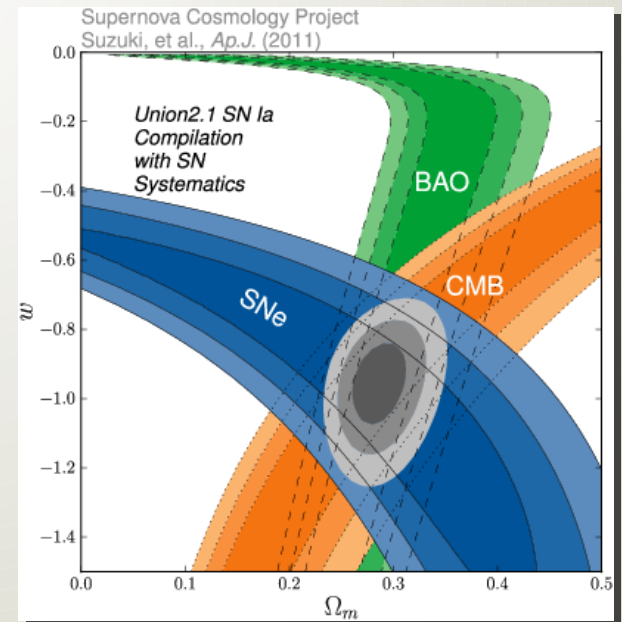
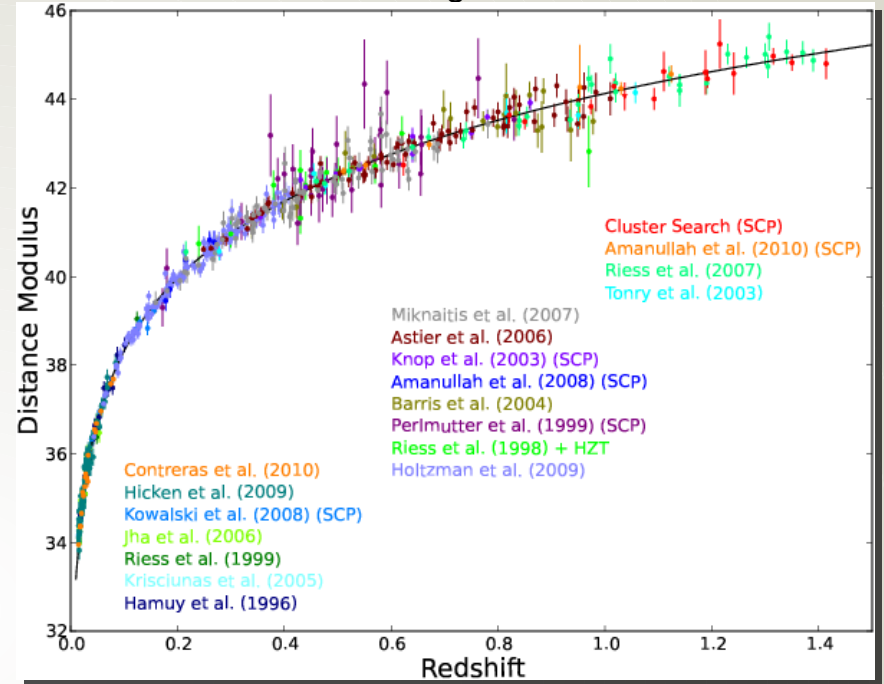


by S. Lilly / A. Refregier

Supernovae Ia

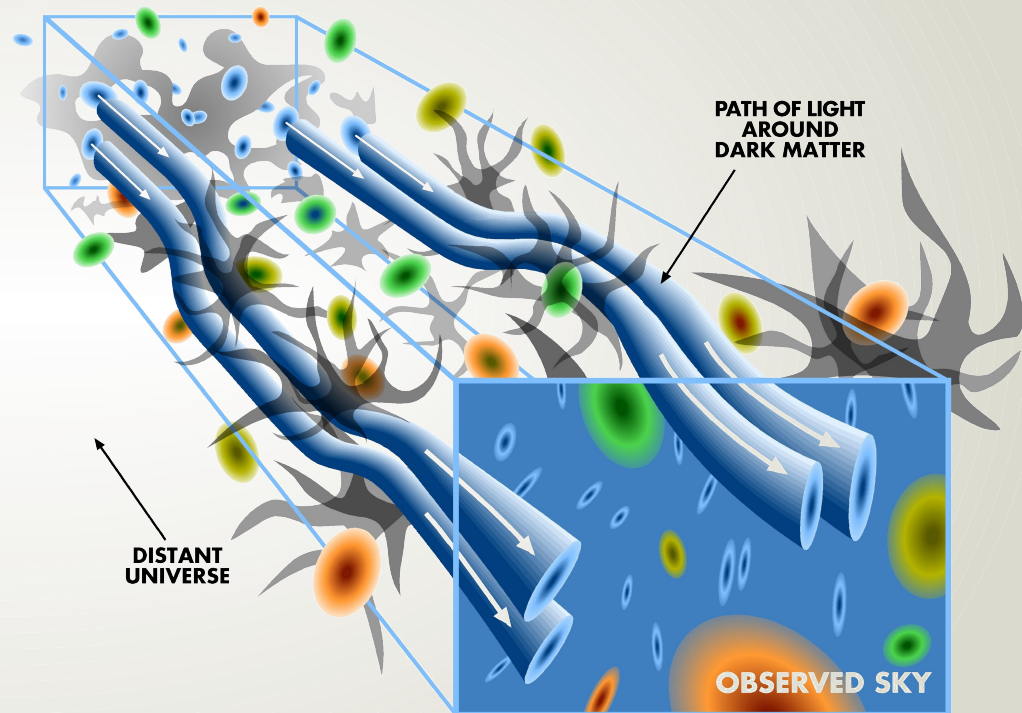
- SNe Ia are standardisable candles which hence provide $D_L(z) \sim \int 1/H(z)$.
- Current datasets give ~ 850 SNe Ia to $z \sim 1.5$ and constrain w to within $\sim 10\%$.
- Many new experiments running or planned but going to high redshifts is hard (no SNe Ia at $z > 2$). Secondary parameters? Evolution?

SCP Union2.1 Hubble diagram:



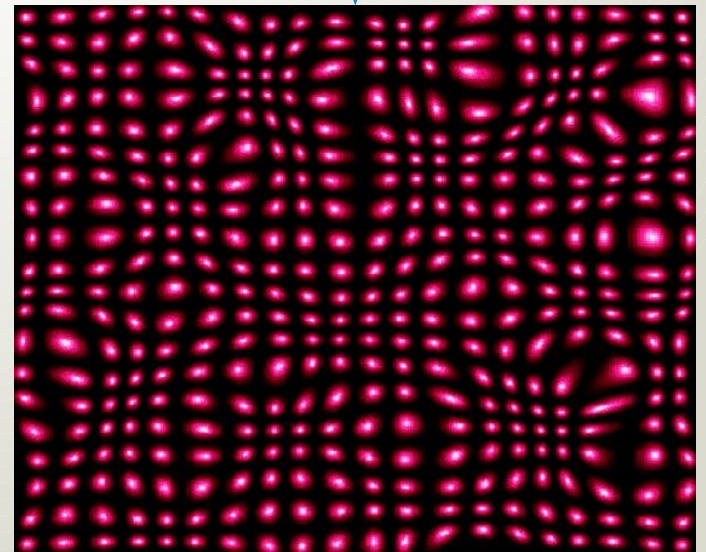
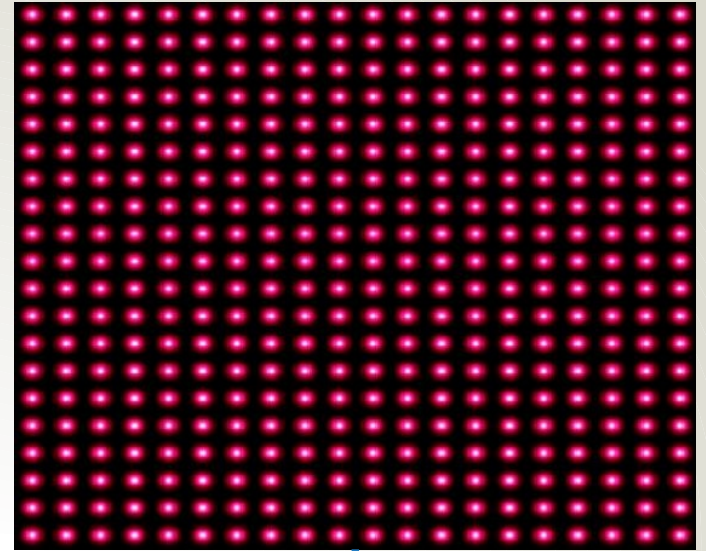
Weak Lensing, Cosmic Shear

- Gravitational lensing by large-scale structure distorts the images of background galaxies.



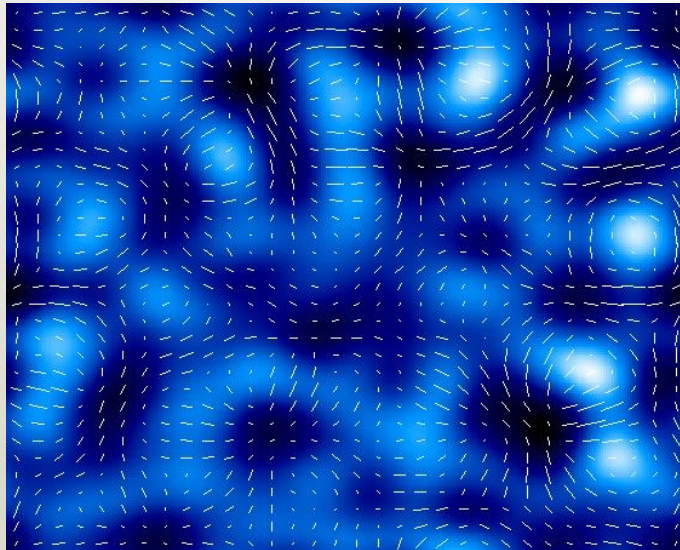
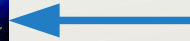
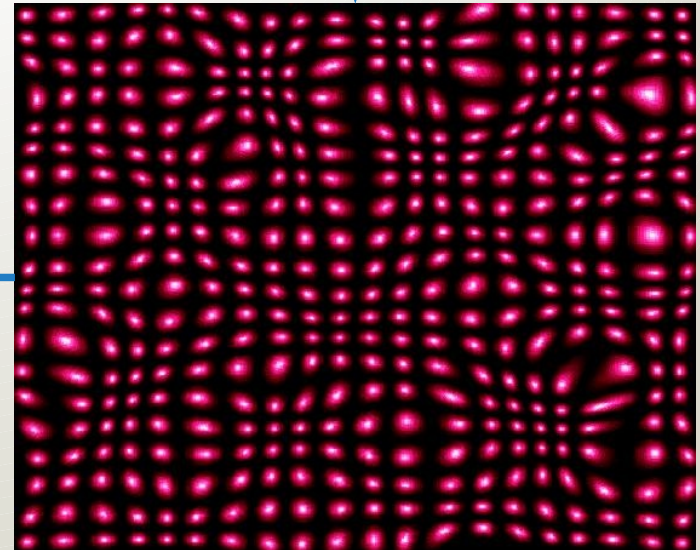
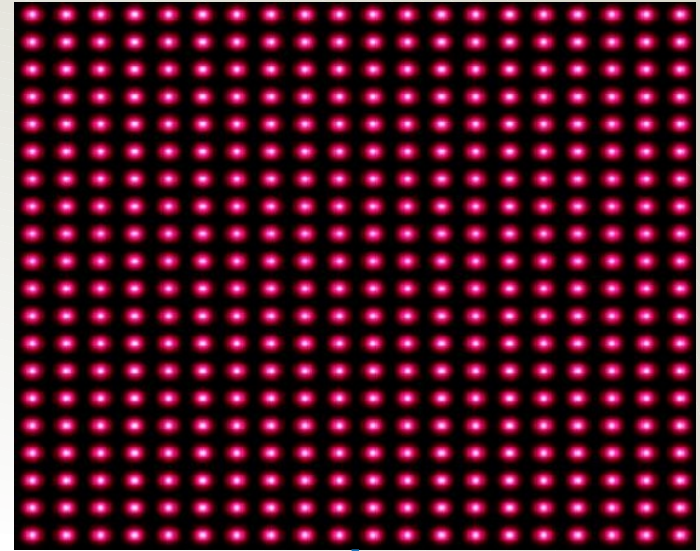
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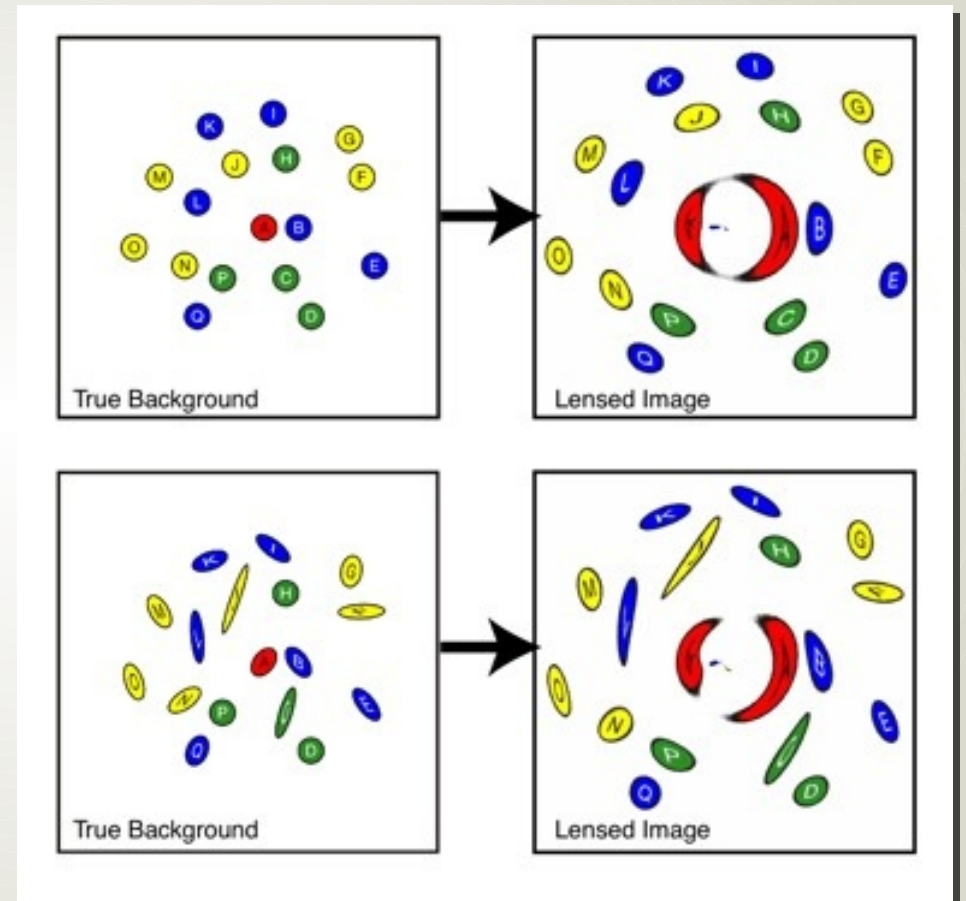
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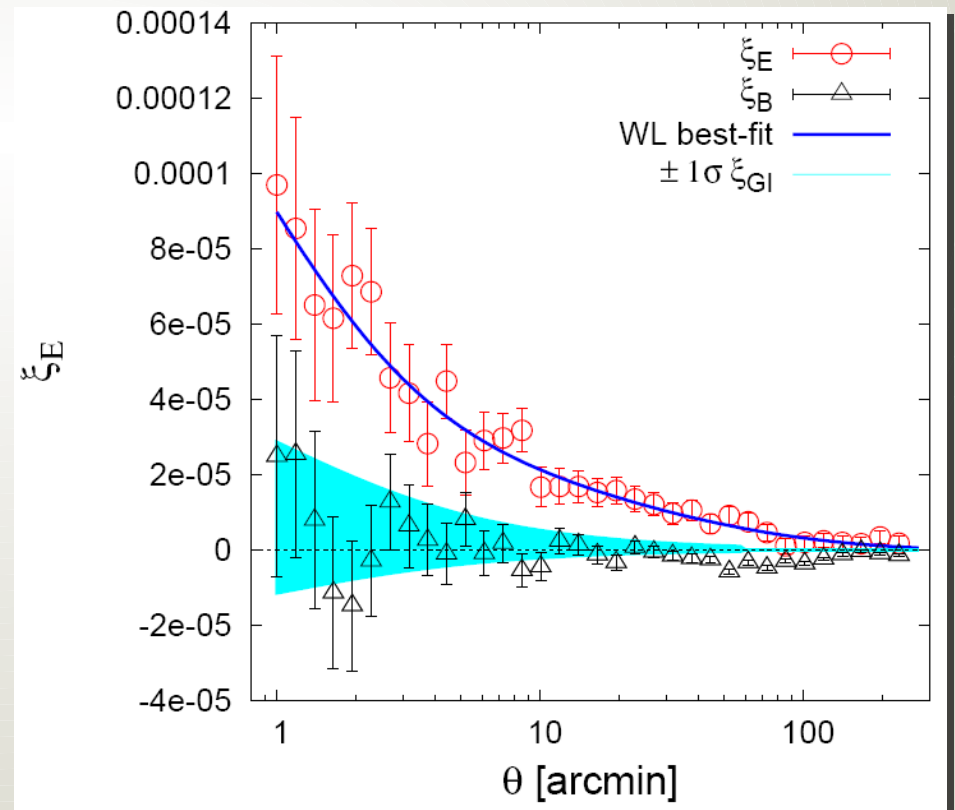
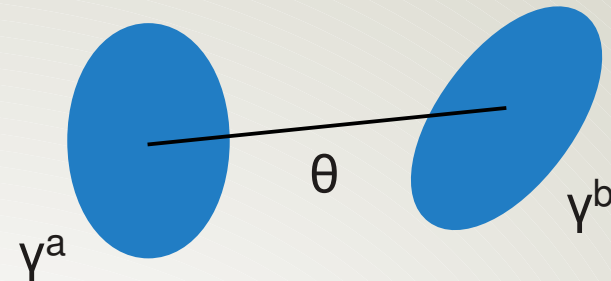
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- Intrinsic ellipticities of galaxies are much larger than shear and act as 'shape noise' → need to combine many galaxies to obtain a signal.



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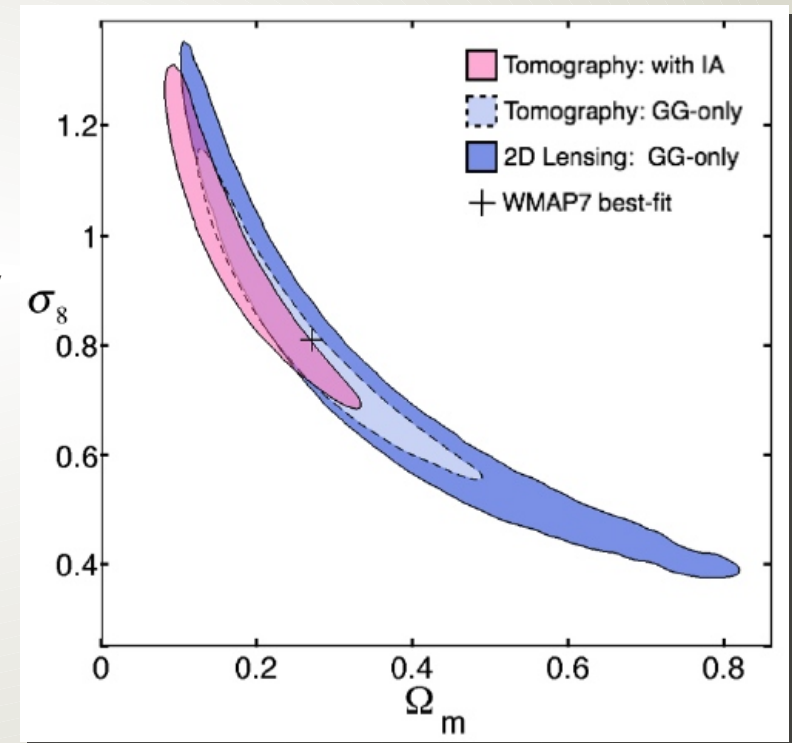
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- Use ellipticities of large samples of galaxies to estimate shear correlation function (or power spectrum).



Weak Lensing, Cosmic Shear

$$P_{\kappa}(l, \chi_s) = \frac{9H_0^4 \Omega_m^2}{4c^4} \int_0^{\chi_s} d\chi \frac{(\chi_s - \chi)^2}{\chi_s^2} \frac{P_{\delta}(l/\chi, \chi)}{a(\chi)^2}$$

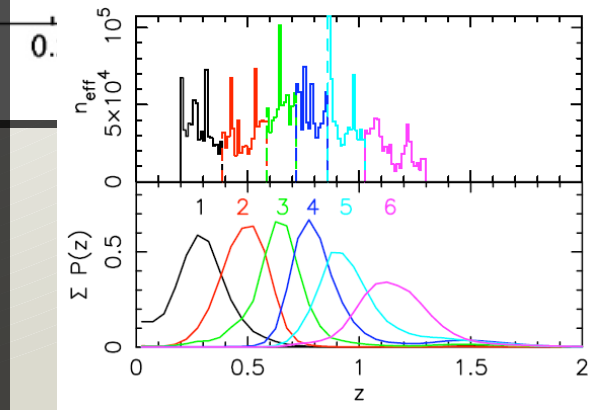
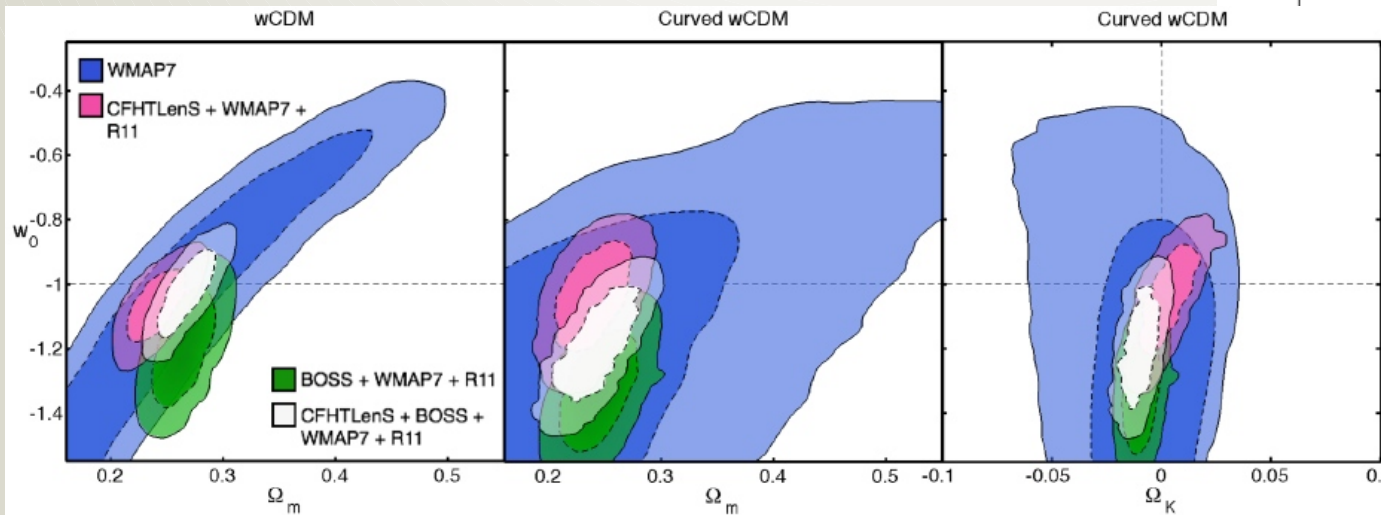
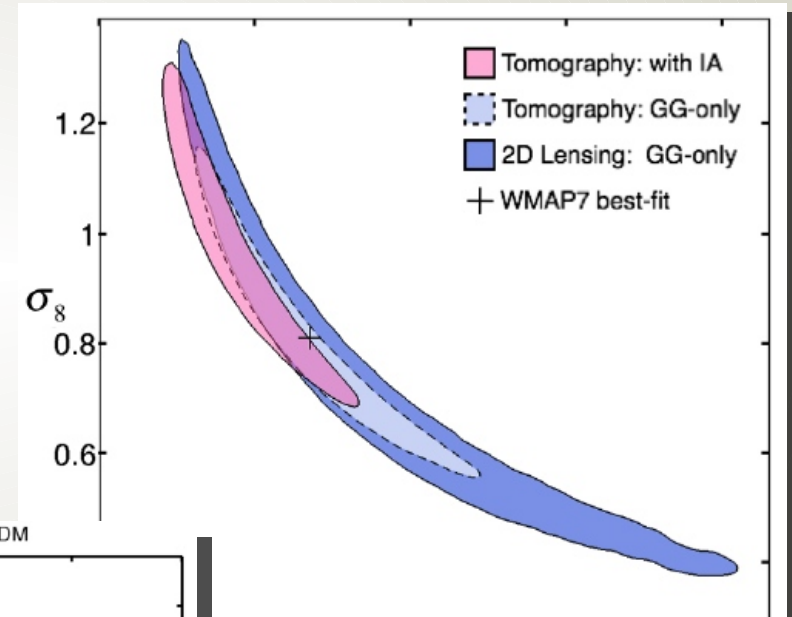
- The shear power spectrum is sensitive to:
 - Matter density Ω_M
 - Amplitude of DM power spectrum σ_8
 - Growth of structure \rightarrow DE, break degeneracy between DE and modified gravity
 - Source distances \rightarrow DE
 - Expansion history \rightarrow DE



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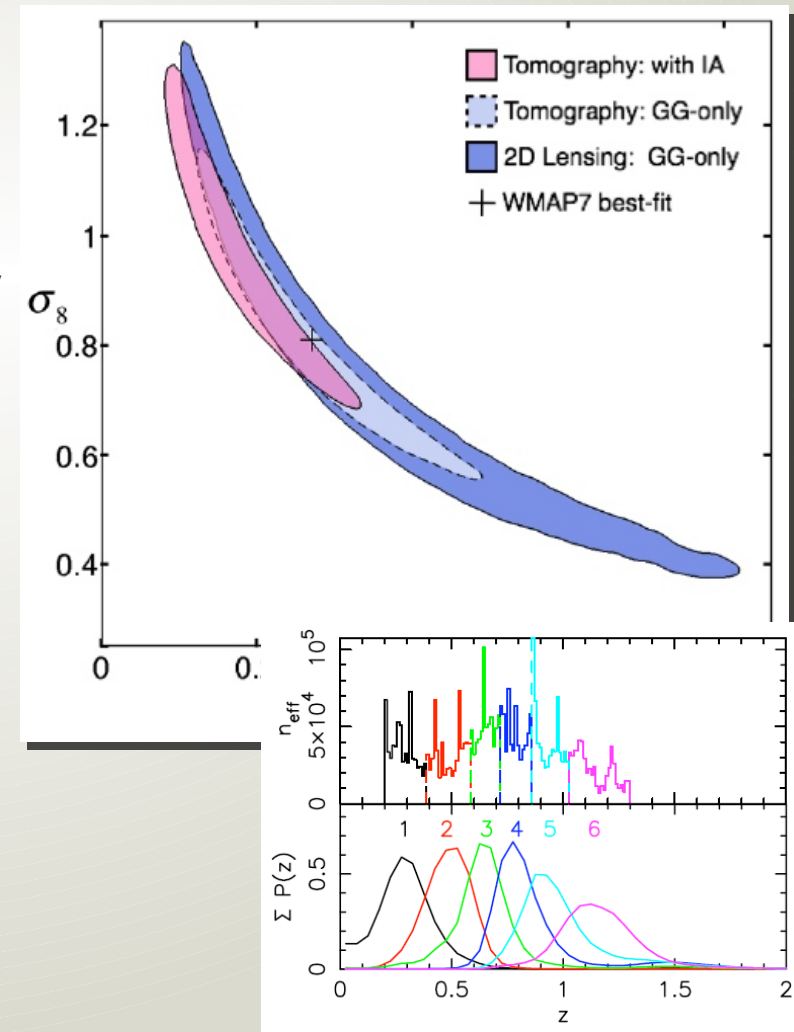
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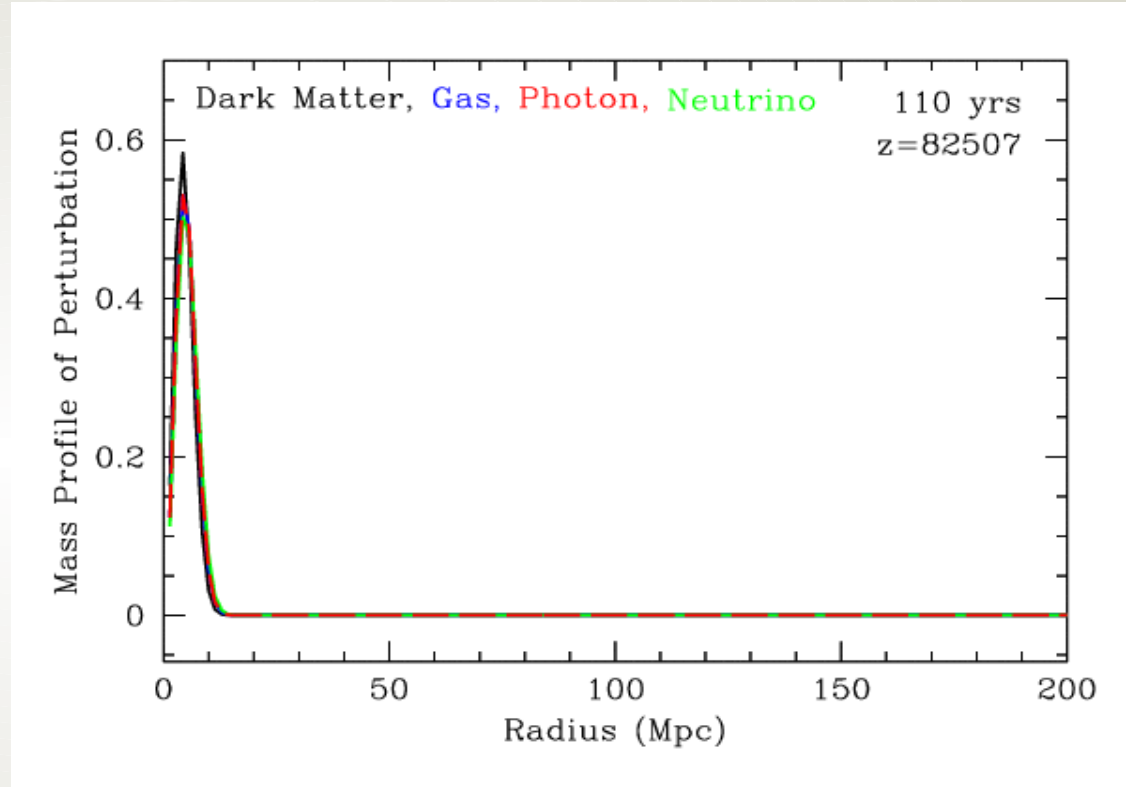
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- This is hard! Need:
 - huge imaging surveys
 - in multiple bands (for photo-z)
 - excellent control of PSF in at least one band
 - shape measurements
 - deal with intrinsic galaxy alignments



Baryon Acoustic Oscillations

Origin of acoustic peaks in CMB and galaxy power spectra (from D. Eisenstein and W. Hu)

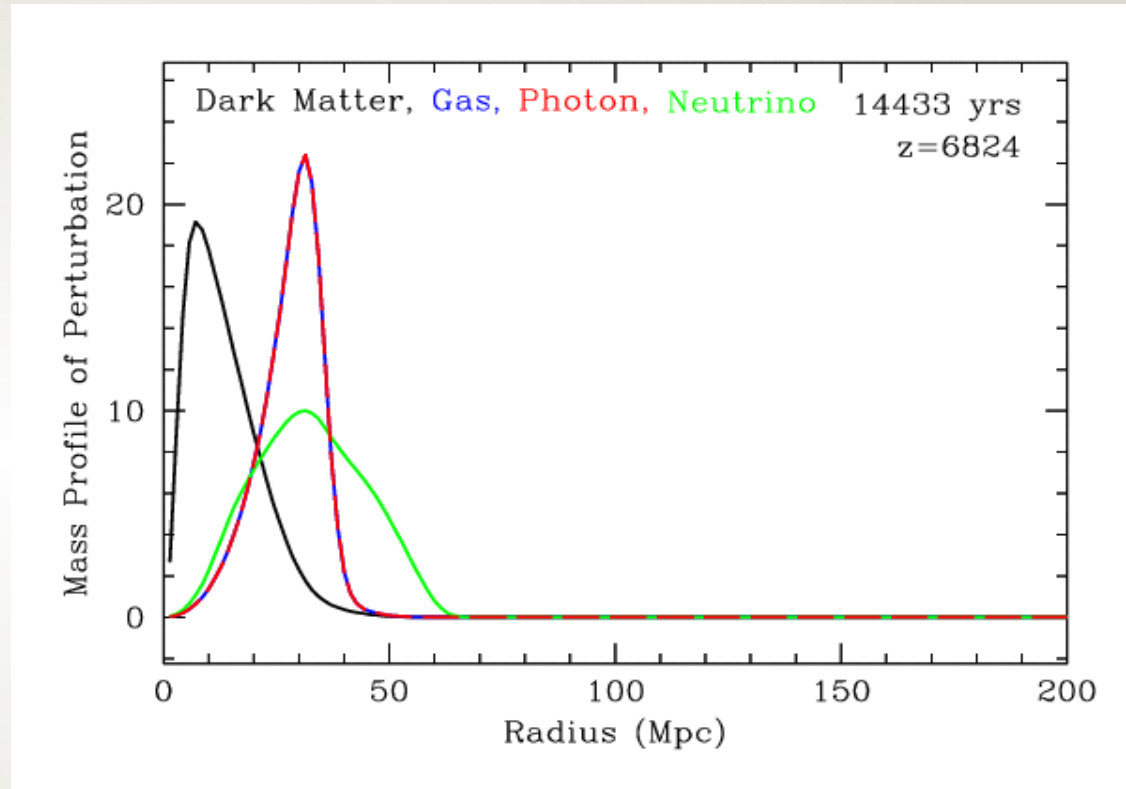
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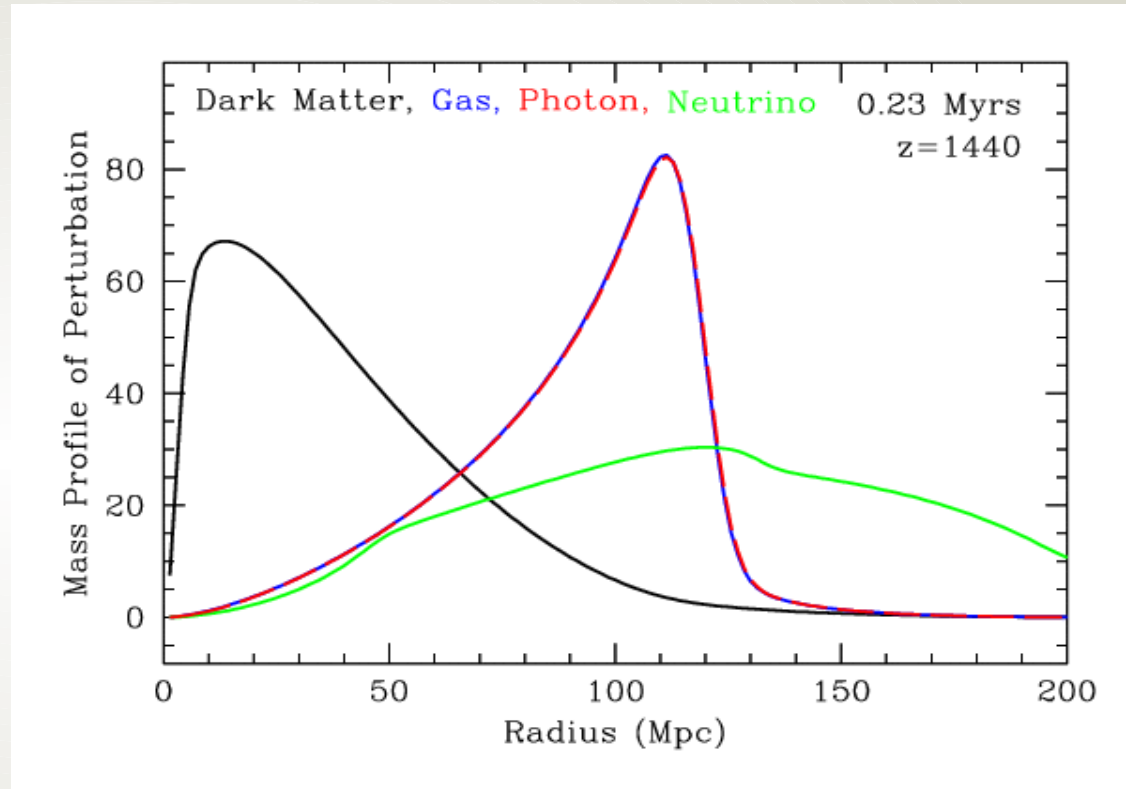
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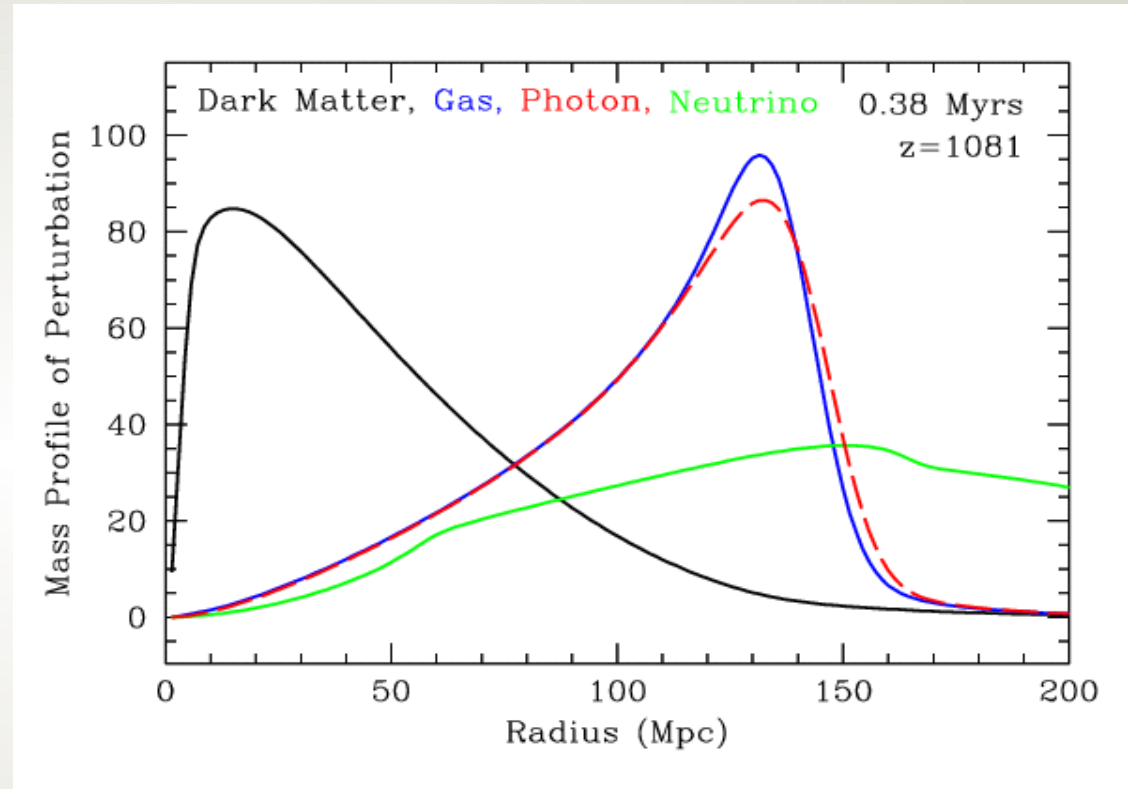
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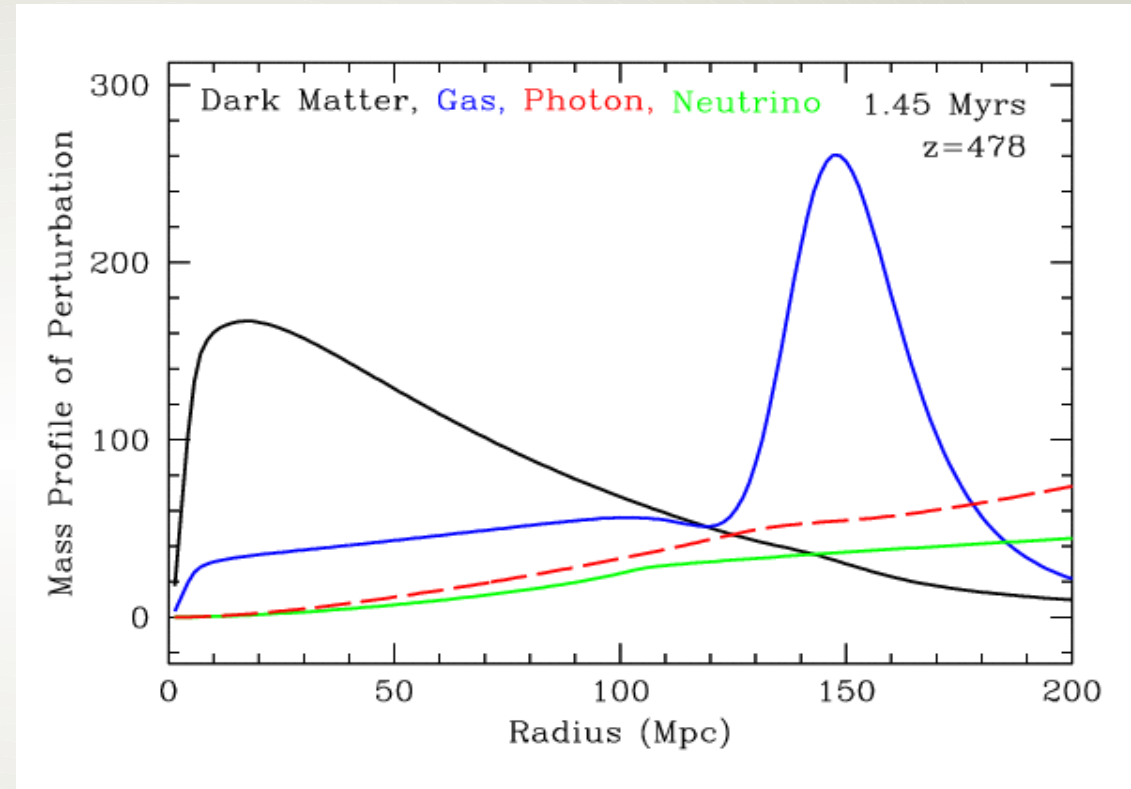
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- At $z \sim 1000$ p, e^- combine \rightarrow photons decouple from baryons $\rightarrow c_s$ drops dramatically \rightarrow baryon perturbation is frozen in.



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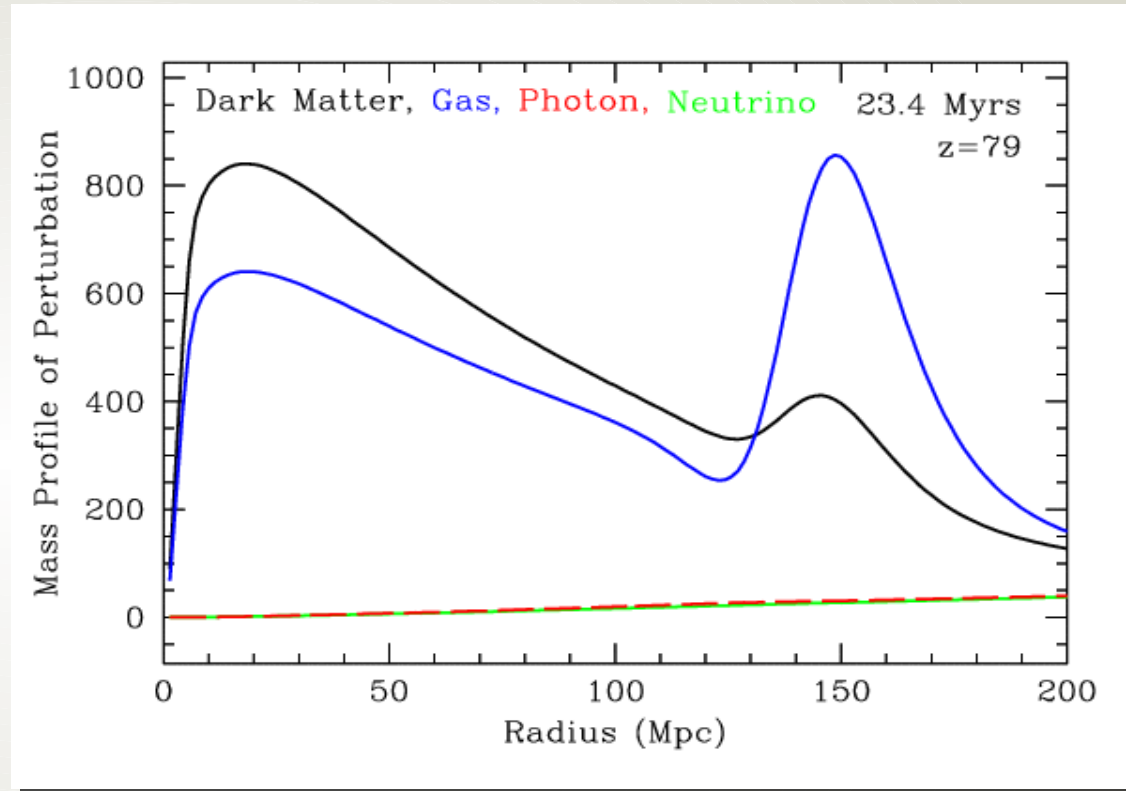
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- At $z \sim 1000$ p, e^- combine \rightarrow photons decouple from baryons $\rightarrow c_s$ drops dramatically \rightarrow baryon perturbation is frozen in.
- Photons free-stream and disperse.



Baryon Acoustic Oscillations

Origin of acoustic peaks in CMB and galaxy power spectra (from D. Eisenstein and W. Hu)

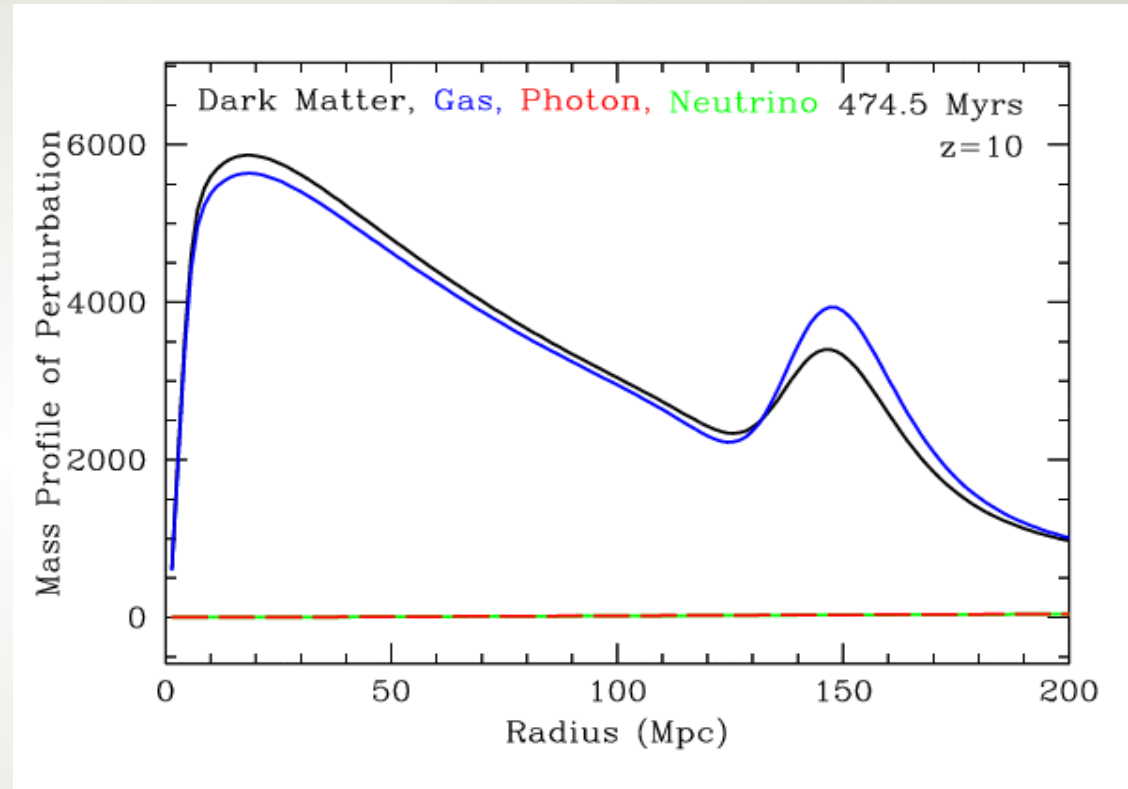
- Consider initial point-like density perturbation in the early Universe.
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 - CDM attracts more CDM
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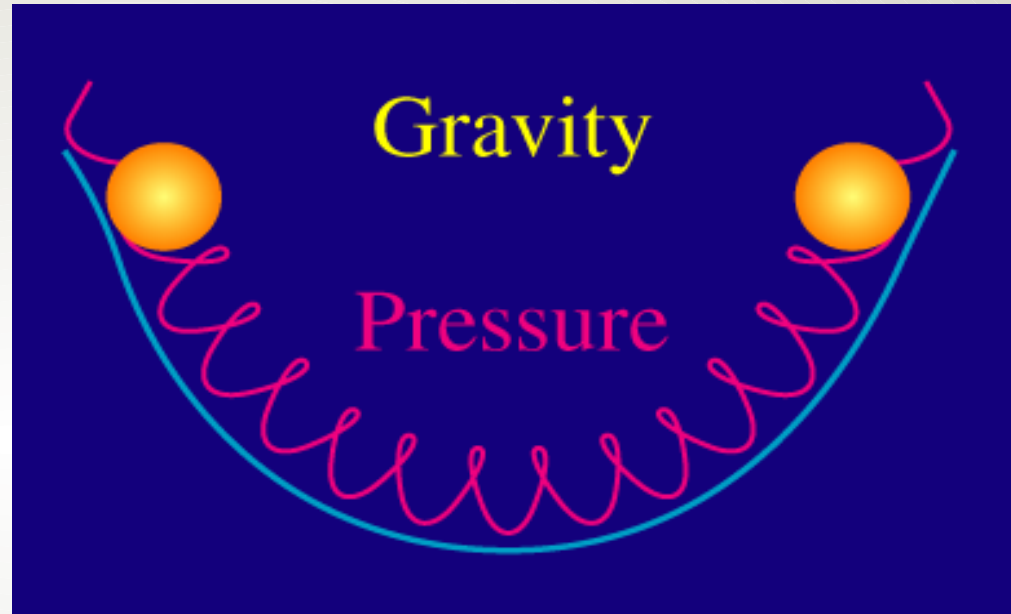


- An overdensity of both baryons and CDM remains at the location of the initial density perturbation as well as at a distance of $c_s * t_{\text{recomb}}$ \rightarrow these act as seeds for galaxy formation \rightarrow a preferred scale is imprinted on the galaxy distribution.

Baryon Acoustic Oscillations

Origin of acoustic peaks in CMB and galaxy power spectra (from D. Eisenstein and W. Hu)

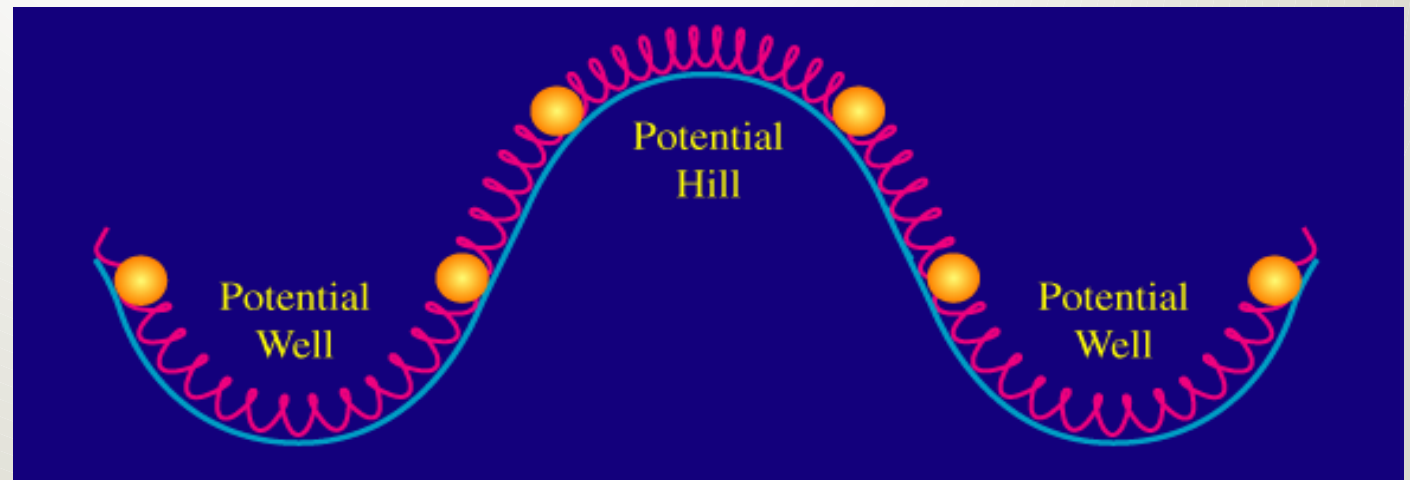
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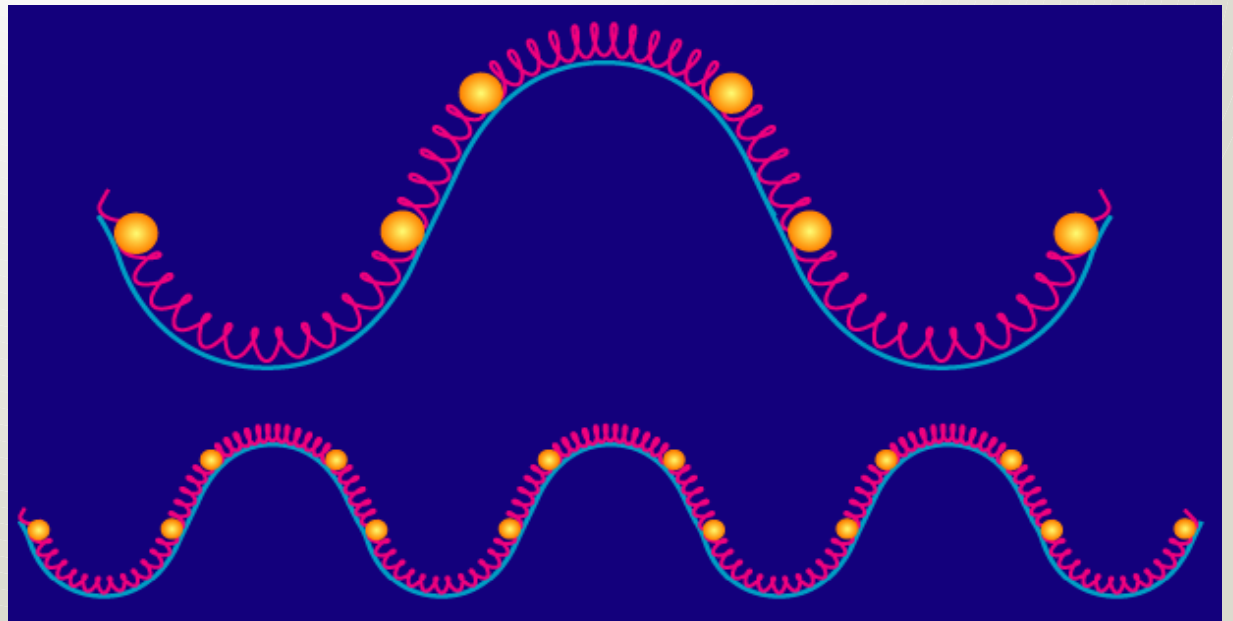
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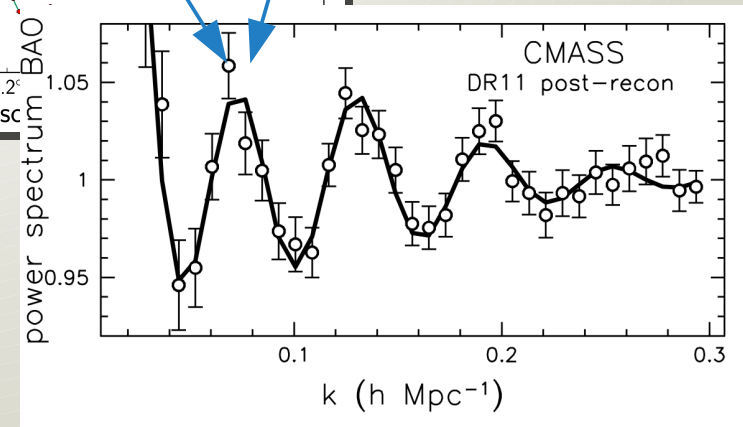
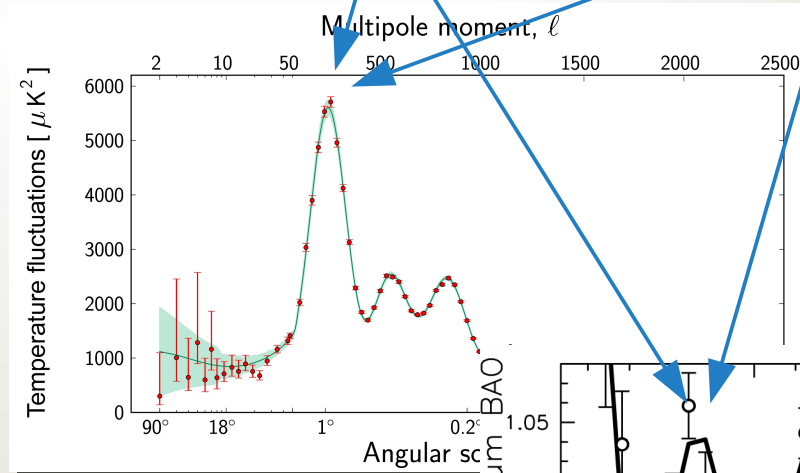
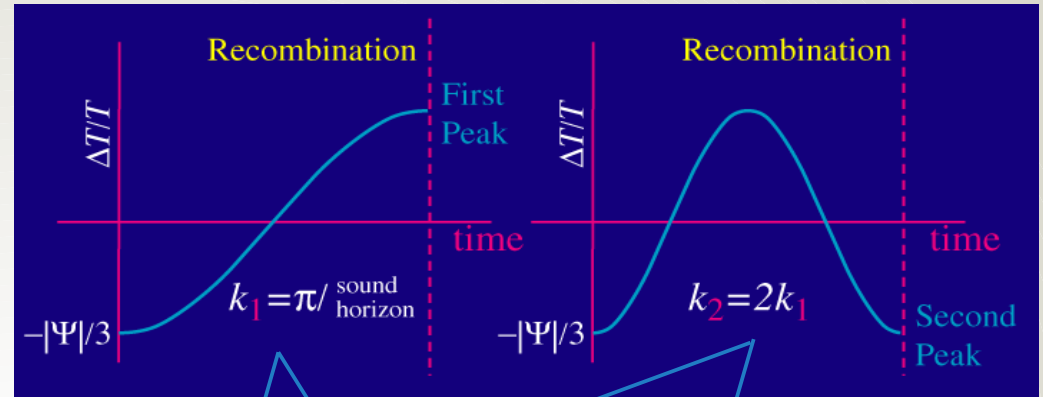
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- In fact, there's a spectrum of perturbations with some power spectrum.



Baryon Acoustic Oscillations

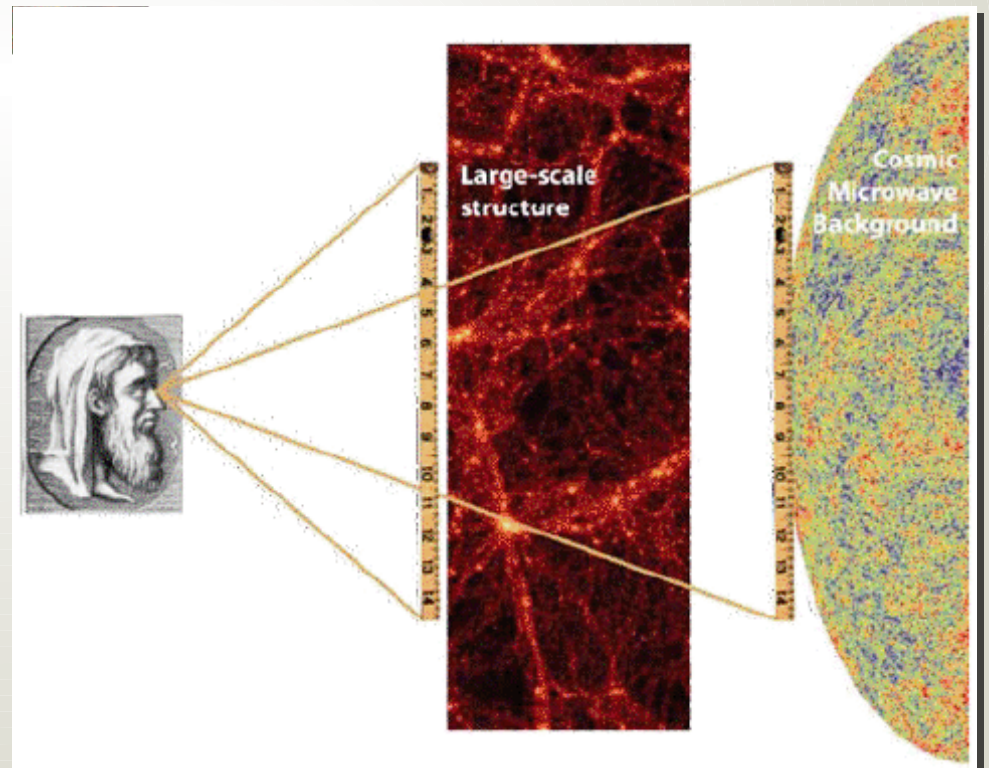
Origin of acoustic peaks in CMB and galaxy power spectra (from D. Eisenstein and W. Hu)

- But it's a wave! So far only considered a single crest.
- And there are many perturbations. So far only considered a single one.
- In fact, there's a spectrum of perturbations with some power spectrum.
- All modes that are multiples of $c_s * t_{\text{recomb}}$ are enhanced.



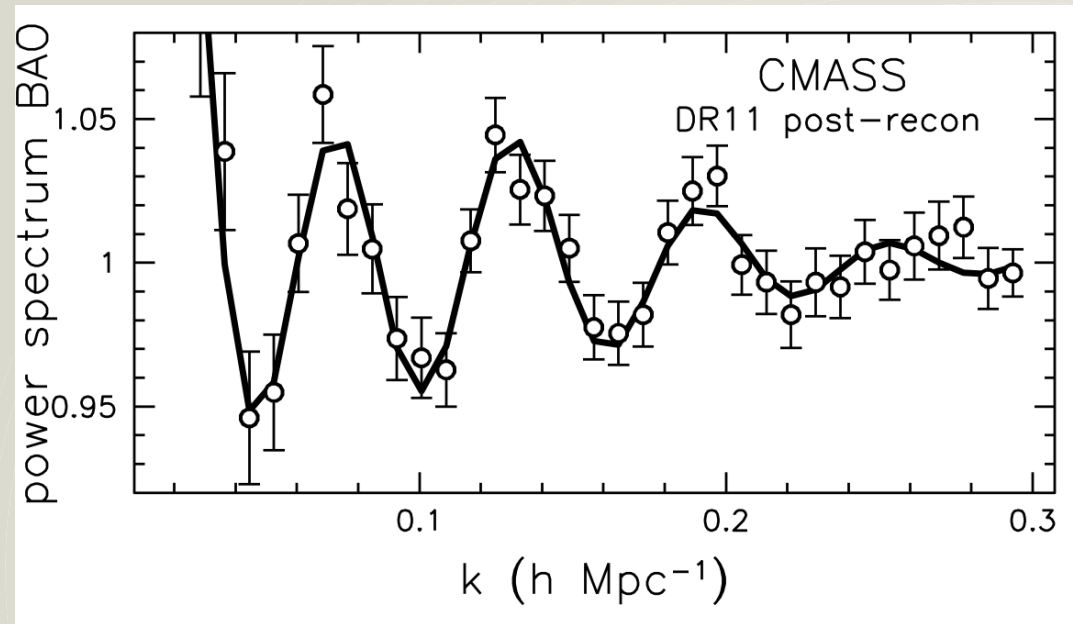
Baryon Acoustic Oscillations

- Geometrical large-angle standard ruler test.
- The ruler itself is based on clean, linear-regime physics at the recombination epoch which is very sensitively probed by the CMB and well understood.
- Provides $D_A(z)$, $H(z)$, $D_V(z)$ (Alcock-Pacinski test).
- Not sensitive to galaxy evolution, dust, etc.
- Does not require precise measurements. Basic galaxy photometry and spectroscopy is enough,
- Works best at $1 < z < 3$.
- Get RSD for free.
- Requires huge samples, i.e. Surveys: volumes of $> 1 \text{ Gpc}^3$
- Needs spectroscopy.
- Works best at $1 < z < 3$.

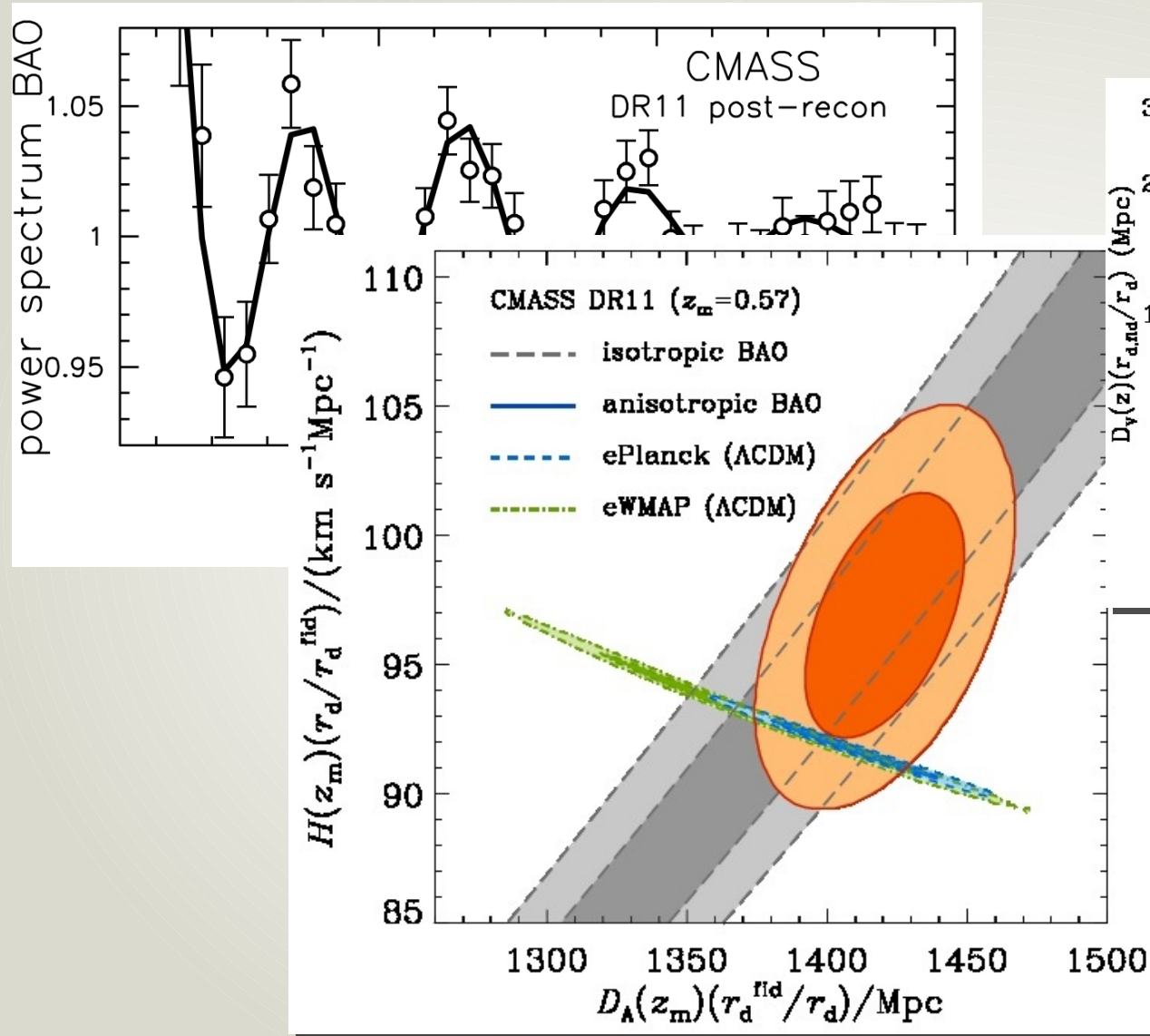


BAO Current Results

BOSS, Anderson et al. (2014)

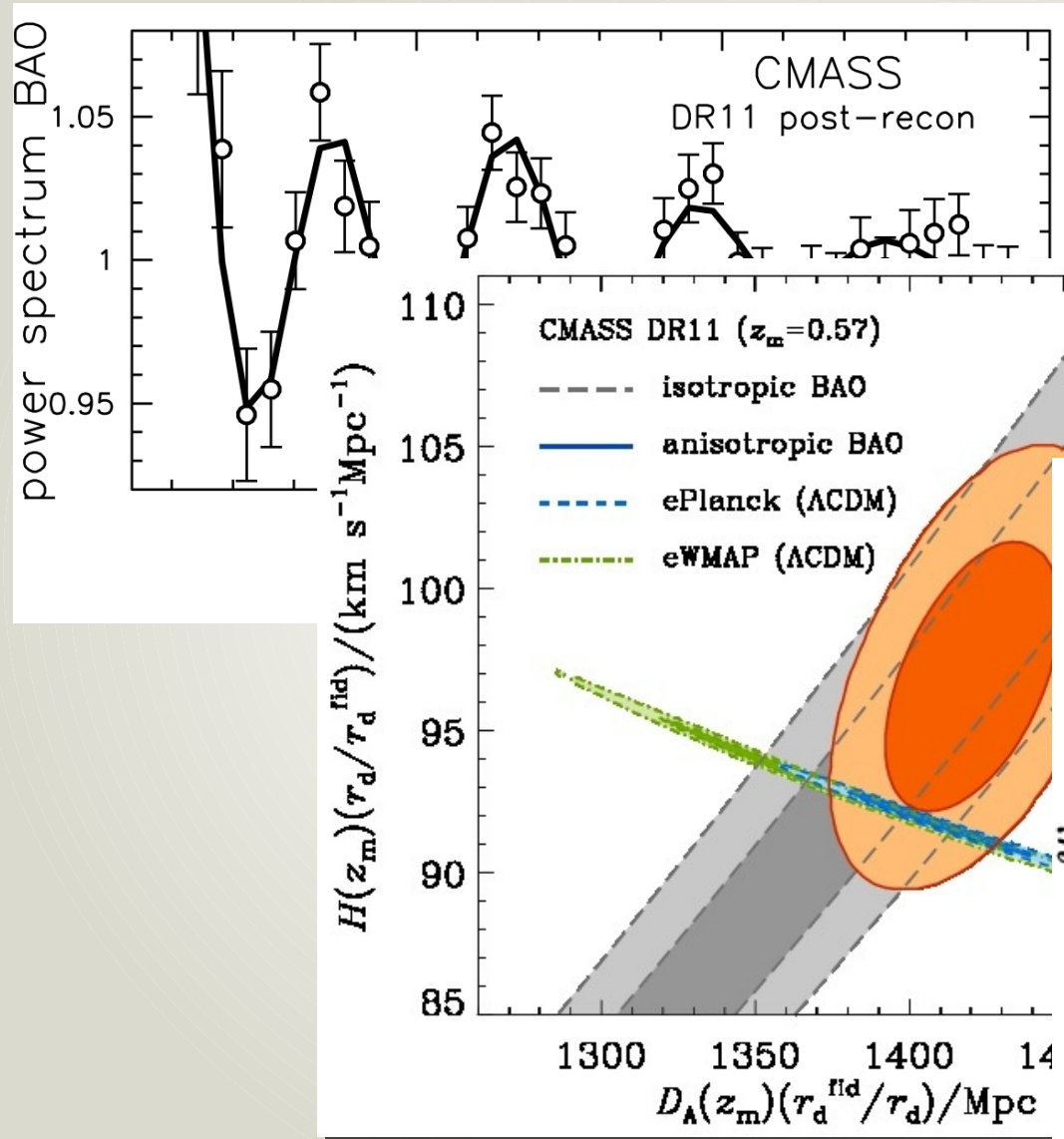


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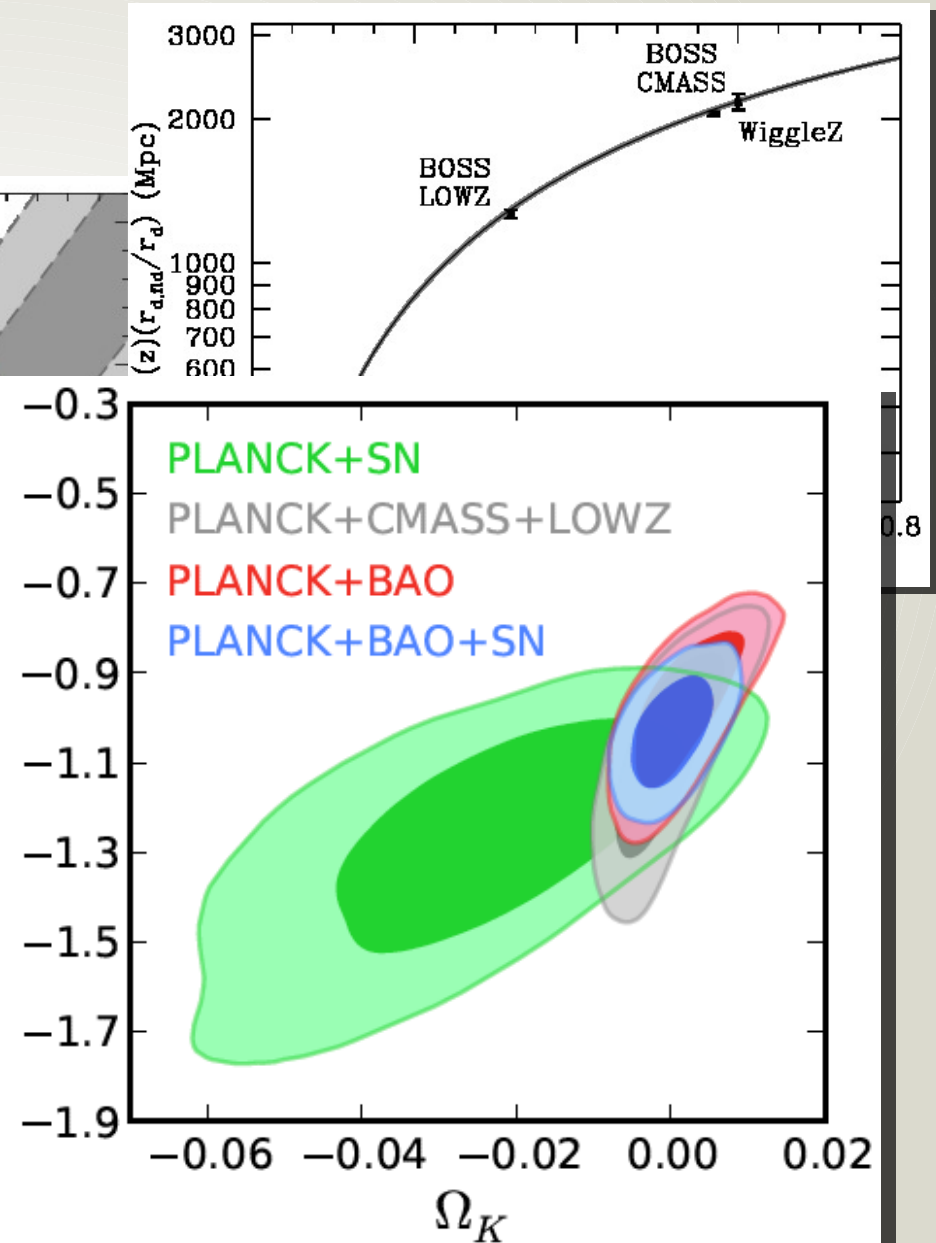


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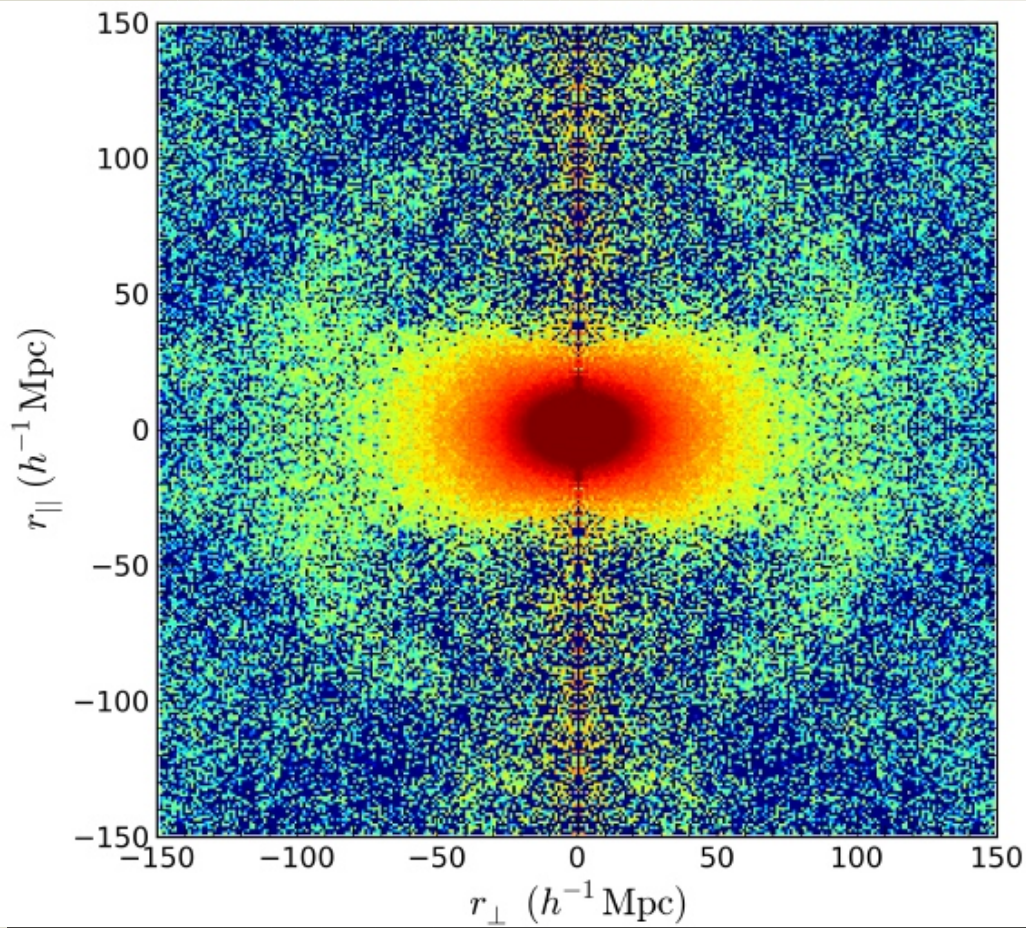
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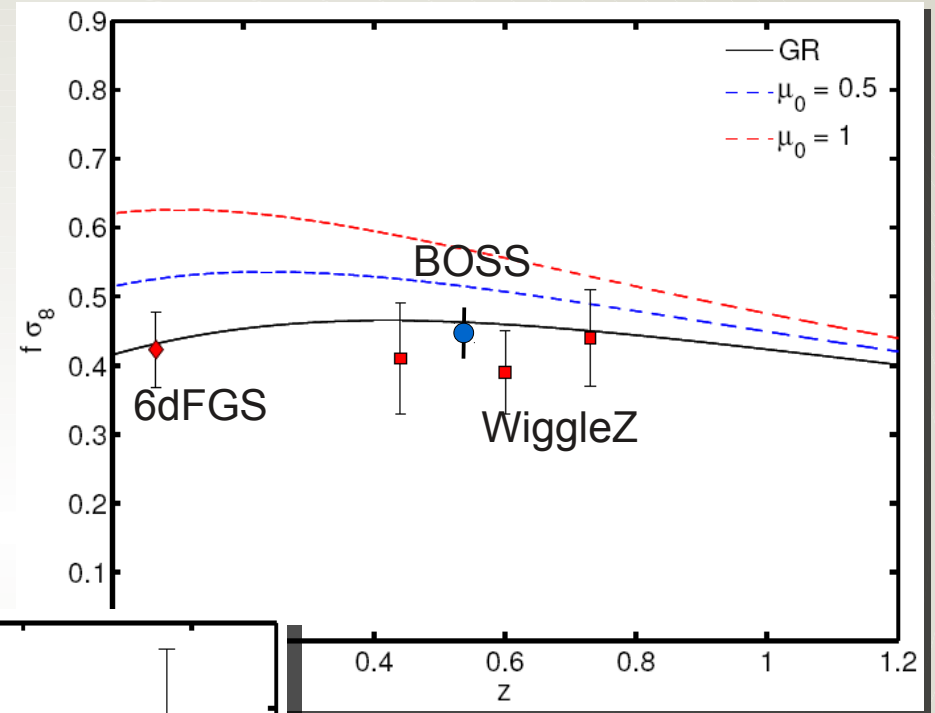
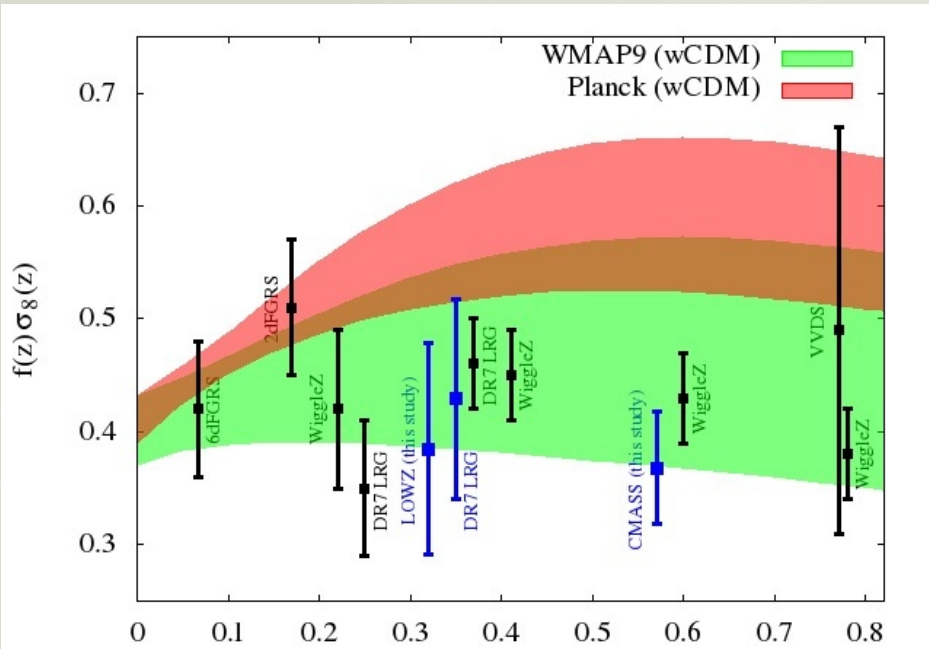
Redshift Space Distortions



Samushia et al. (2013)

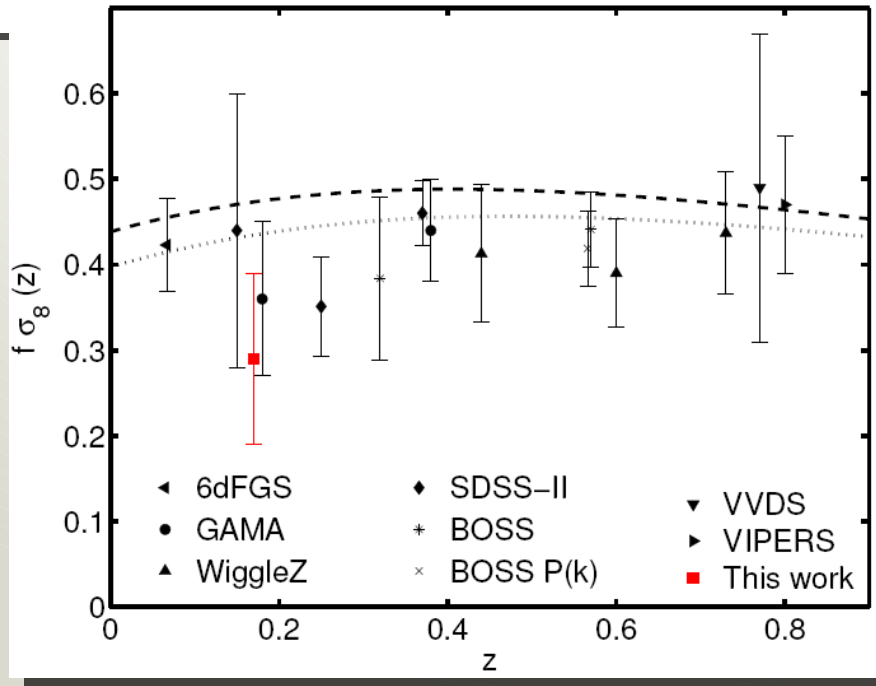
- Measured redshifts include not only the Hubble flow but also peculiar velocities:
 - on small scales: finger-of-God effect in collapsed structures
 - on large scales: infall into high-density regions and outflow from low-density regions (Kaiser effect)
- Creates anisotropy between the LOS and transverse correlation functions.
- Anisotropy constrains $\sigma_8^* d\ln G/d\ln a$, i.e. the growth of structure.
- Breaks the degeneracy between DE and modified gravity models with the same $H(z)$.
- Again need big redshift surveys, but get them 'for free' with BAO surveys.

RSD Current Results



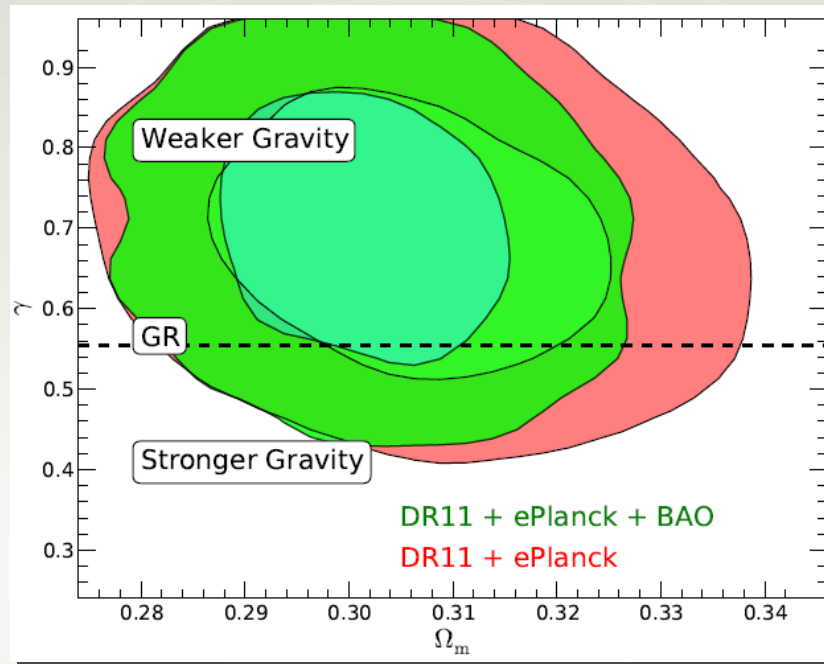
Chuang et al. (2014)

Simpson et al. (2015)

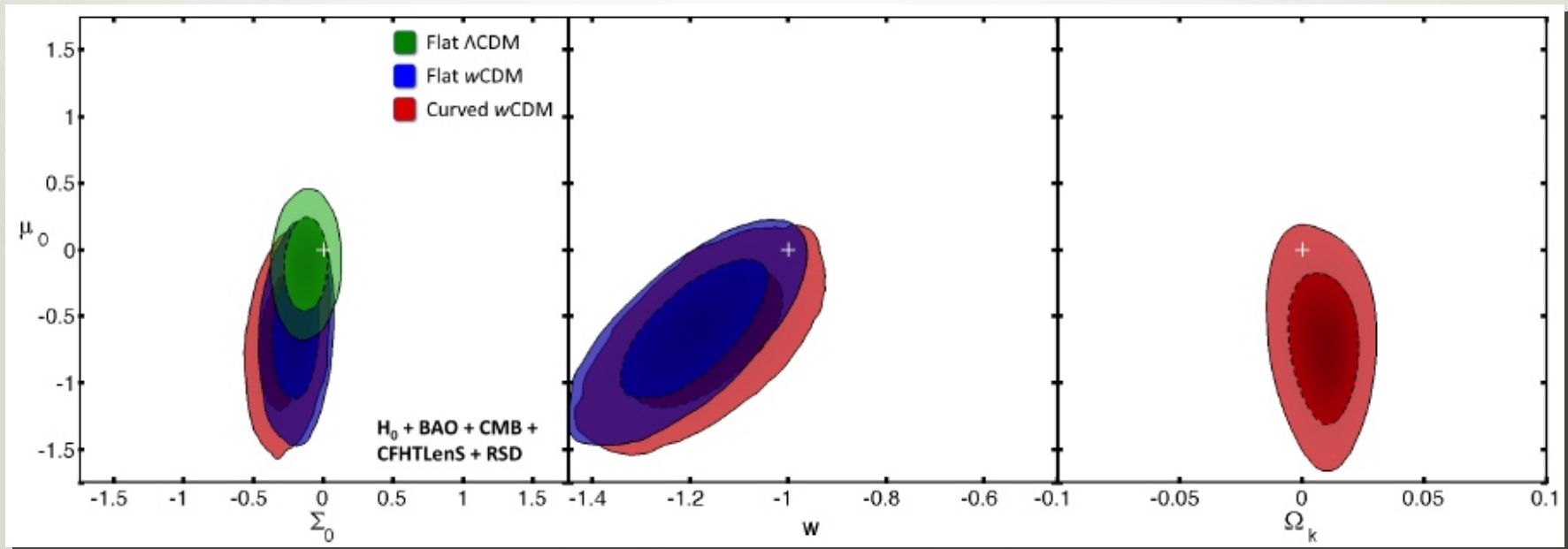


Simpson et al. (2013)

RSD Current Results



Samushia et al. (2013)



Simpson et al. (2013)

Accelerated Expansion

- Intense interest in the expansion history.
Best current methods of measuring $H(z)$:
 - SNIa
 - Weak lensing
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- Many, many, many (really very many) surveys ongoing and planned surveys to probe any combination of the above (plus some more). These will constrain w and MG at the level of $\sigma_w \sim 0.01$ and $\sigma_\gamma \sim 0.01$.

DES
OzDES
CFHTLenS
Pan-STARRS
eBOSS
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These methods are essentially geometric in nature and/or probe the dynamics of localised density perturbations.

A measurement of the *global dynamics* has never been attempted. This would offer a direct, entirely model-independent route towards $H(z)$.

Observing the Expansion

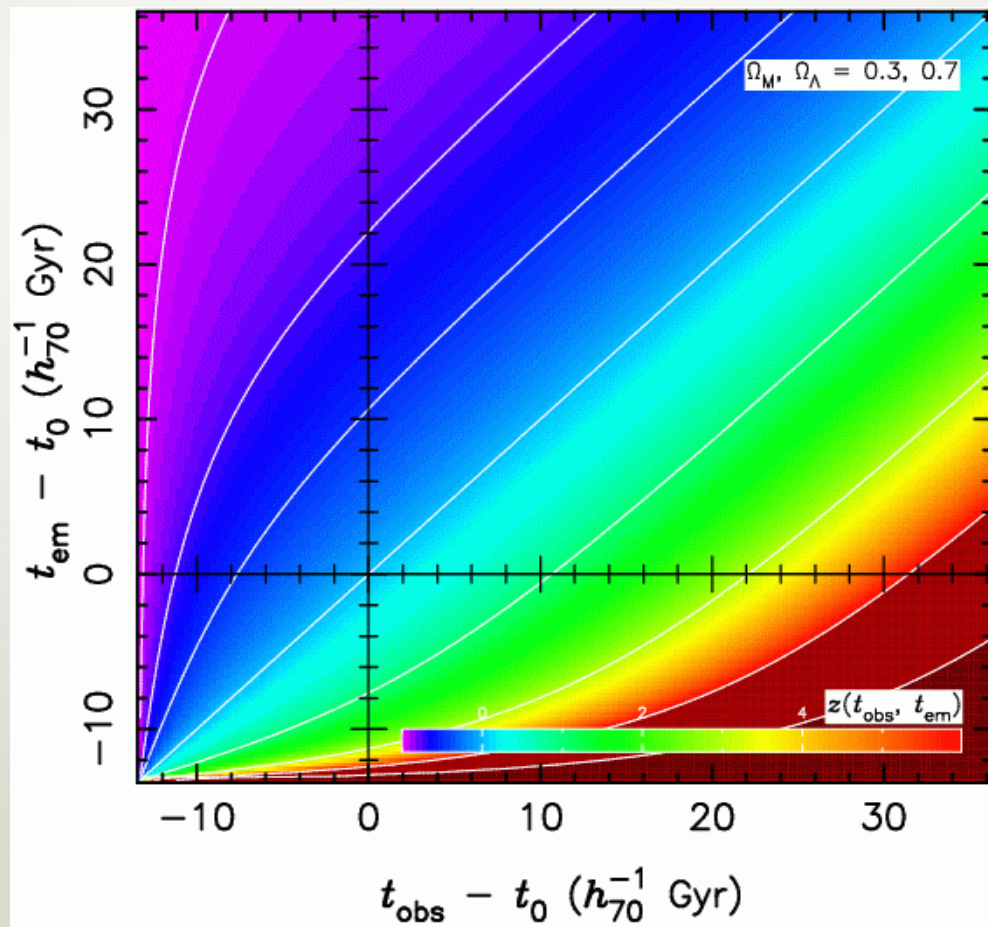
A photon emitted by some object at comoving distance χ at time t_{em} and observed at t_{obs} suffers a redshift of:

$$1 + z(t_{obs}, t_{em}) = \frac{a(t_{obs})}{a(t_{em})}$$

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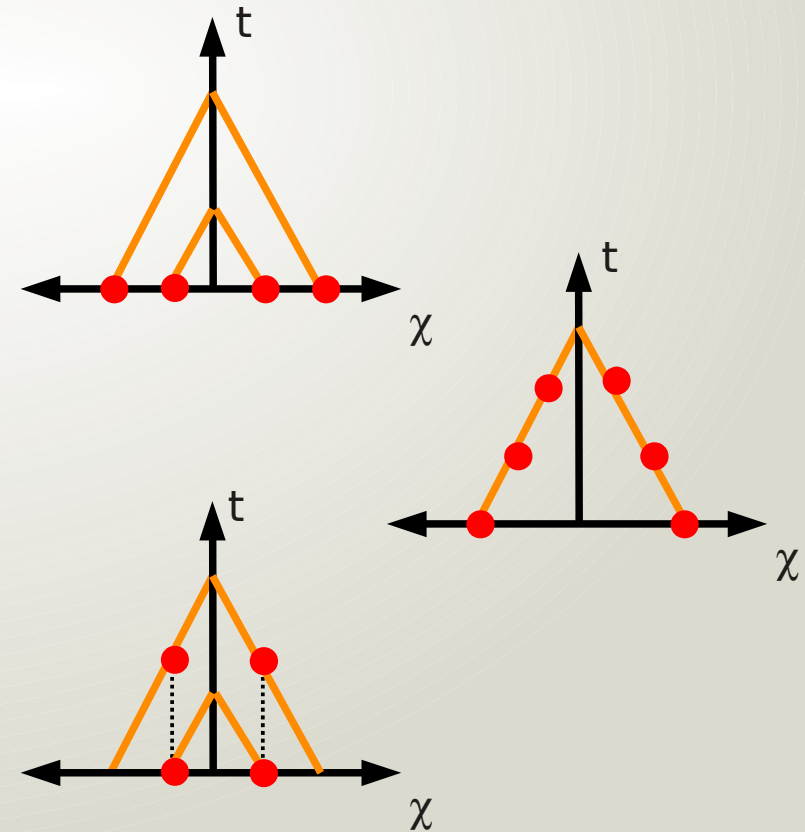
$$1 + z(t_{obs}, t_{em}) = \frac{a(t_{obs})}{a(t_{em})}$$

Three ways to look at this equation:

1. $1 + z(t_{obs}) = \frac{a(t_{obs})}{a(t_{em})}$ χ varies

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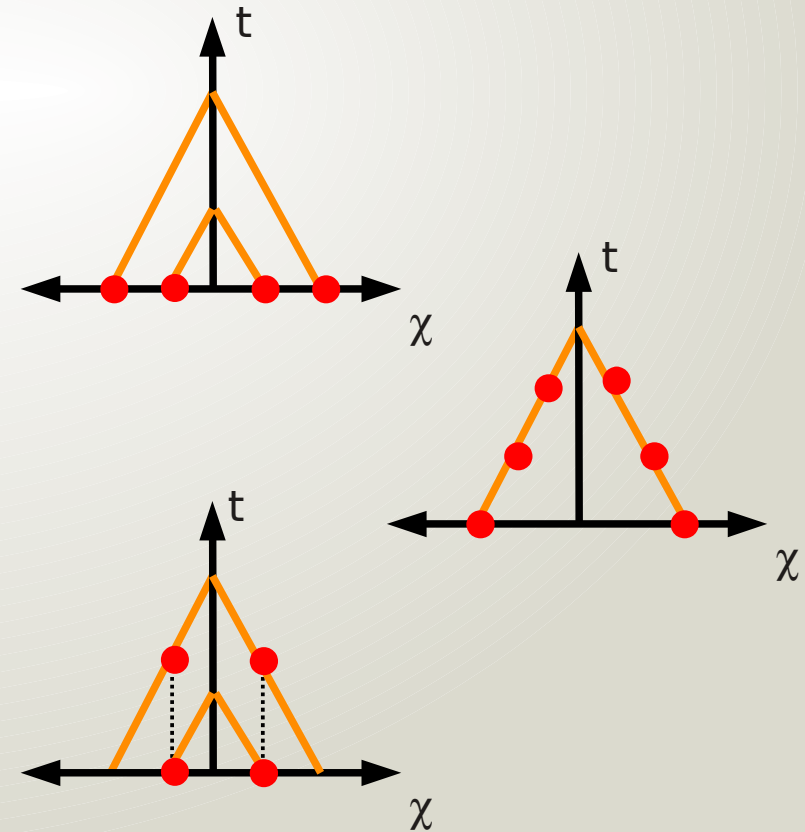
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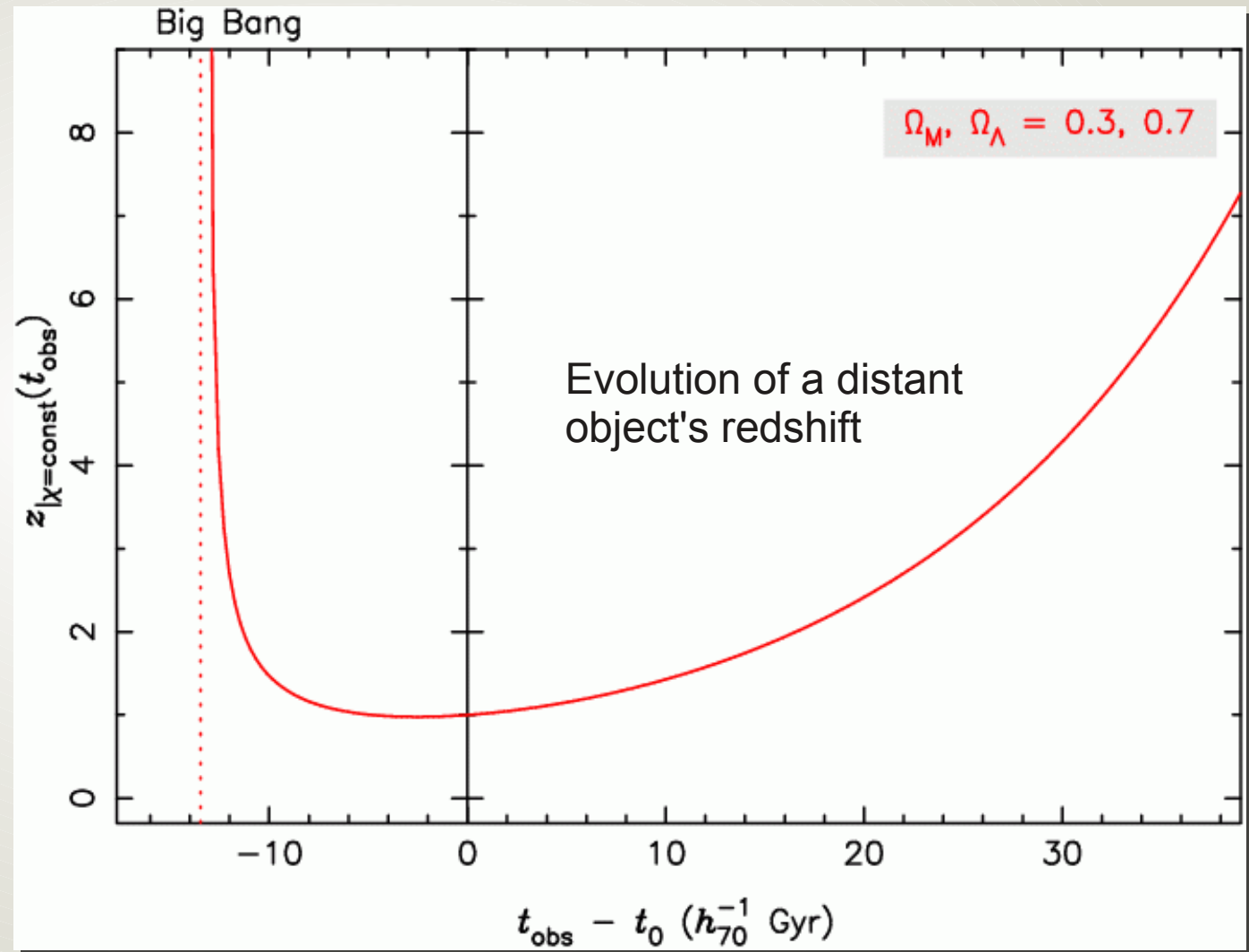
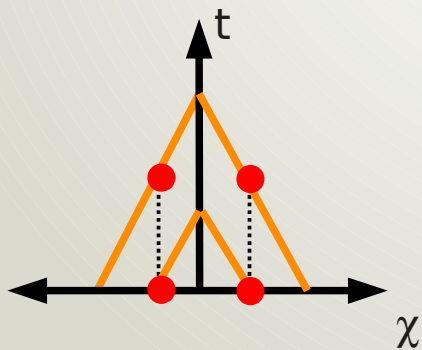
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Evolving redshifts

The evolution of an object's redshift with time contains the entire expansion history.

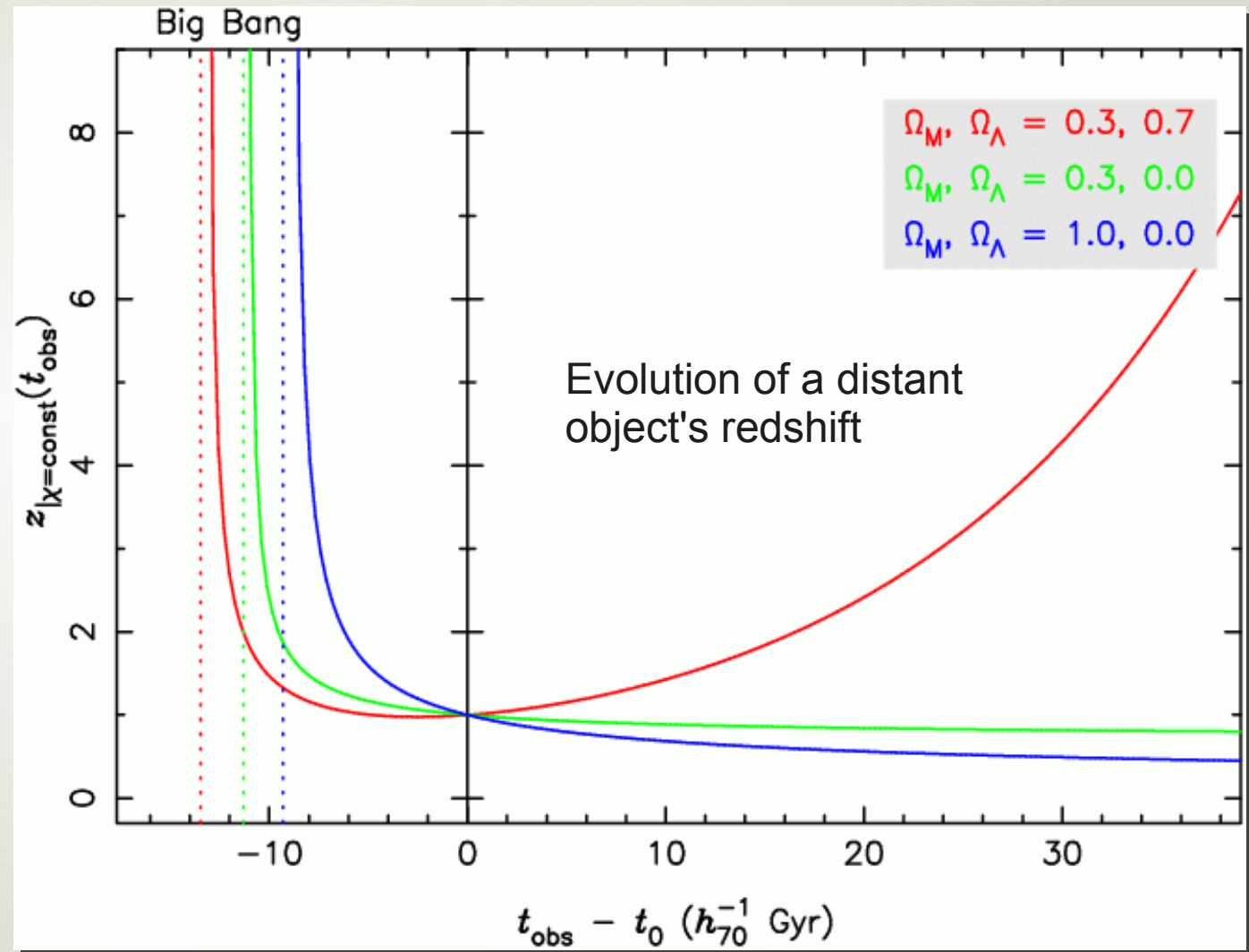
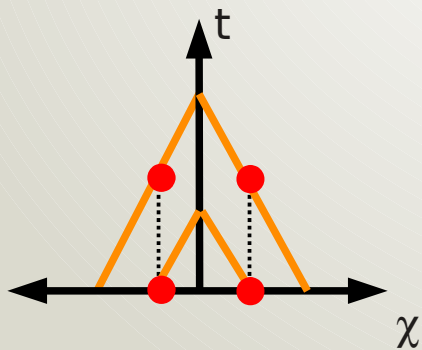
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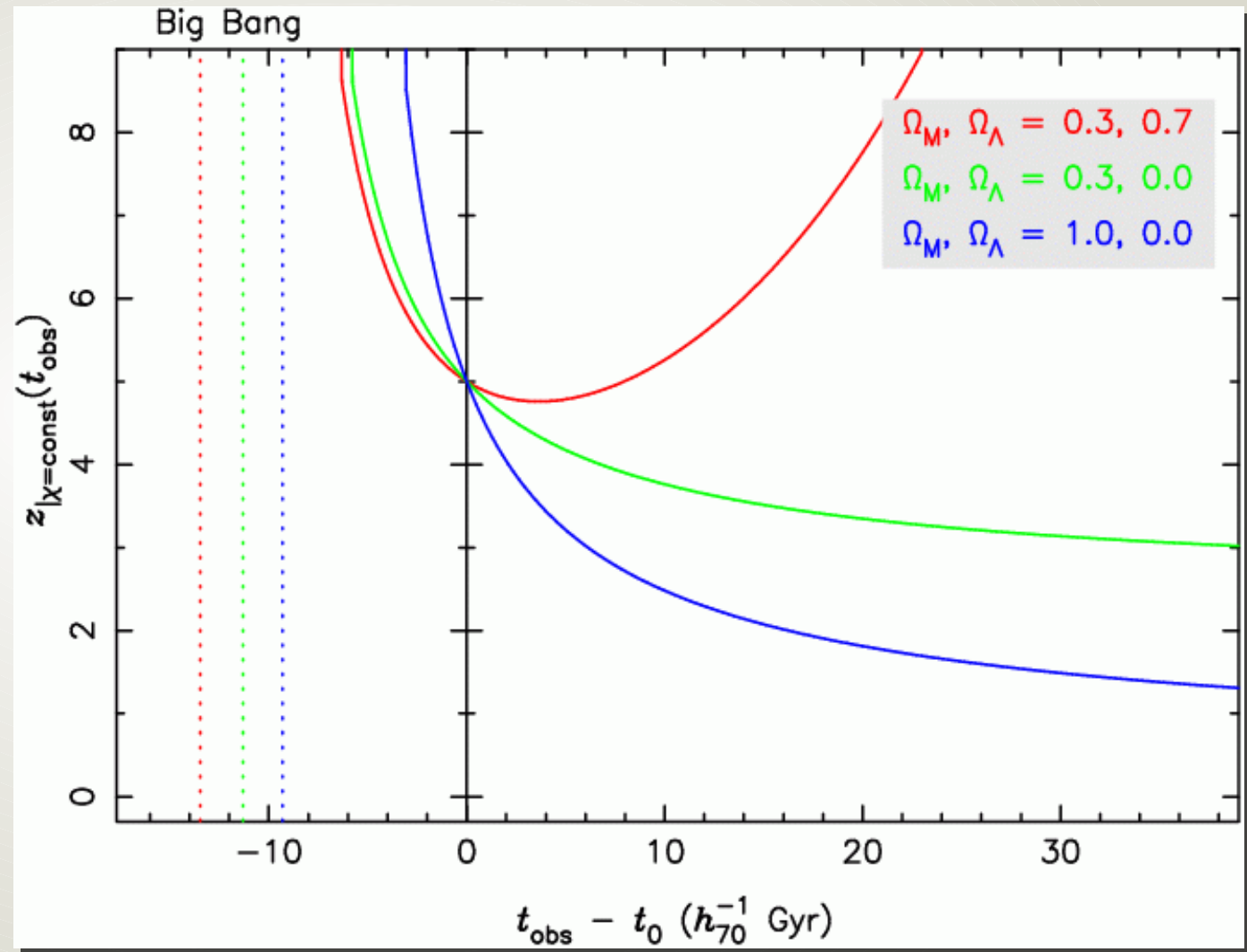
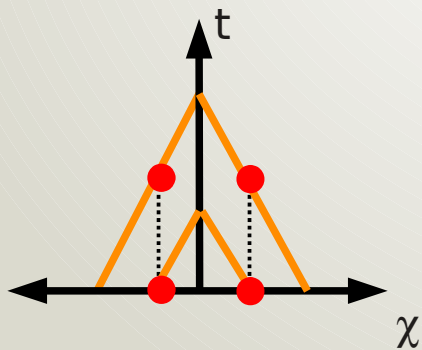
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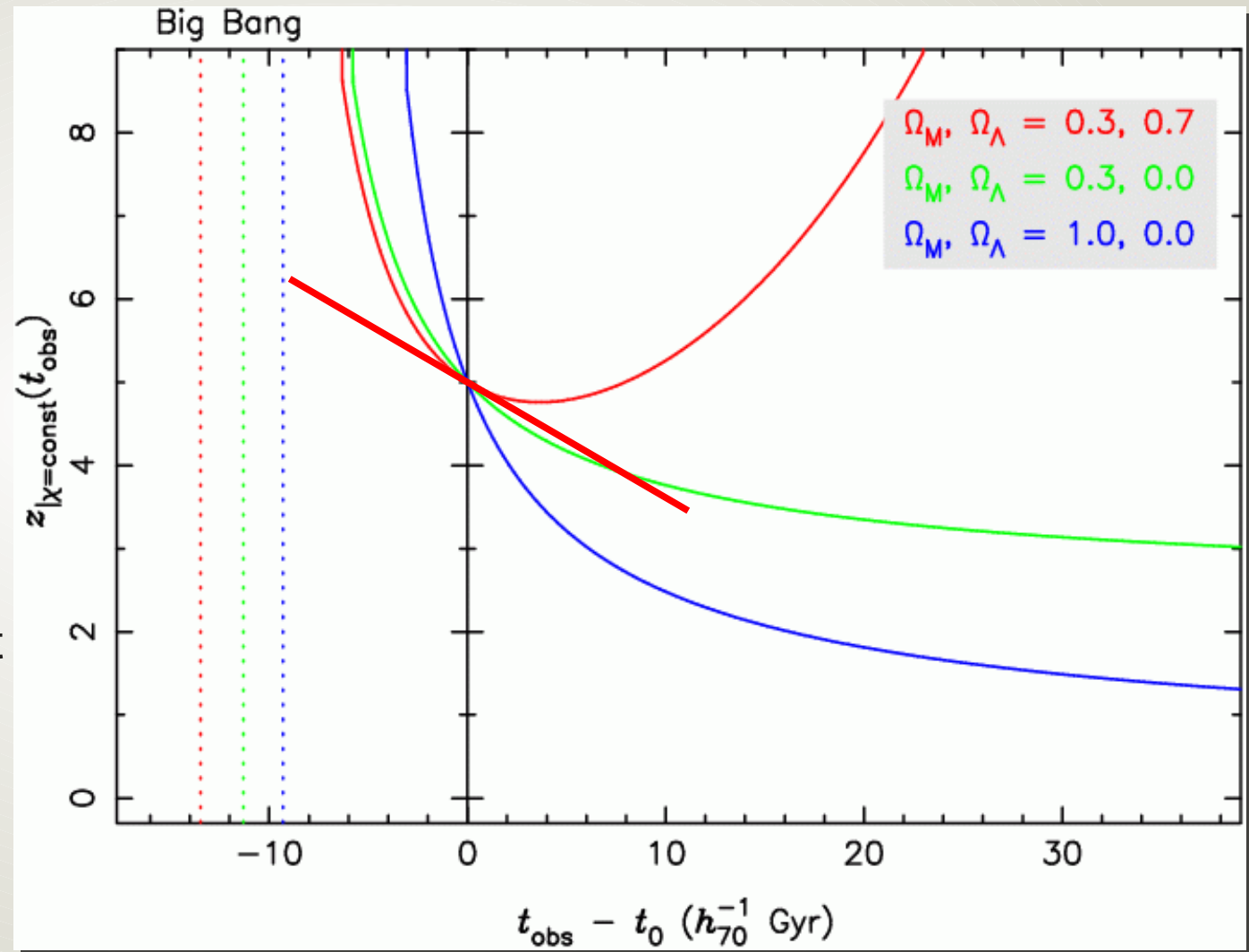


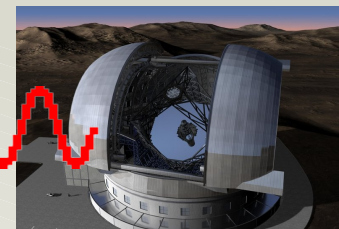
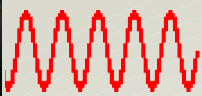
Evolving redshifts

To use $z(t_{\text{obs}})$ to reconstruct the expansion history we need to observe for Gyrs!
Alternative: measure

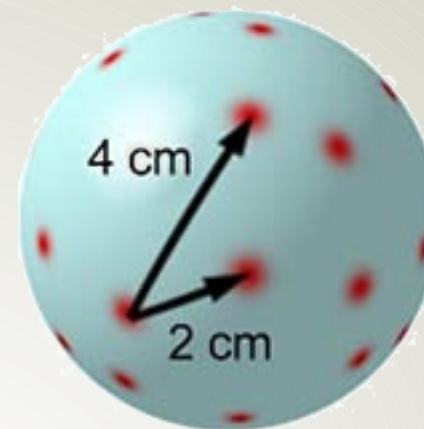
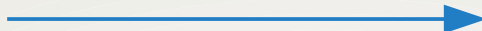
$$\frac{dz}{dt_{\text{obs}}}$$

dz/dt = change of redshift as a function of time.

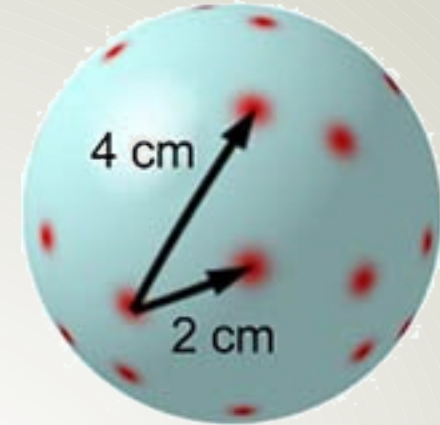
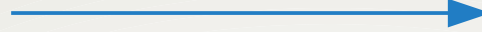
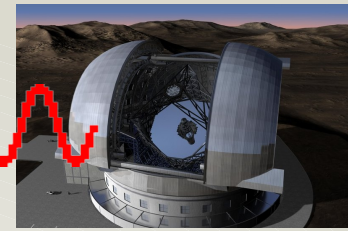
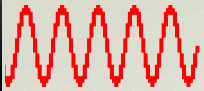




5 billion years ago

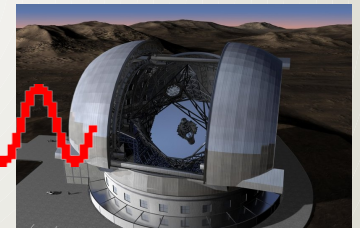
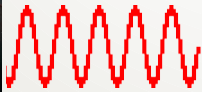


Today

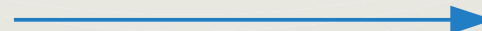


5 billion years ago

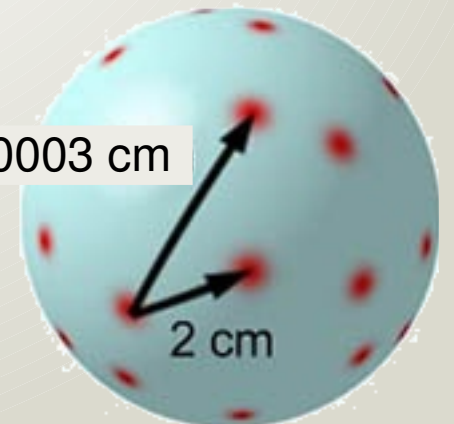
Today



2.0000001 cm

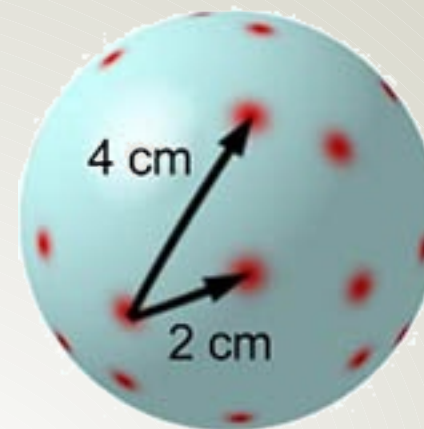
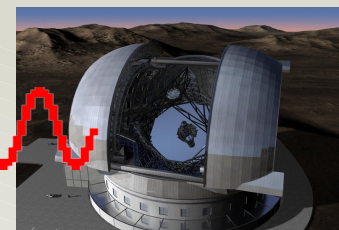
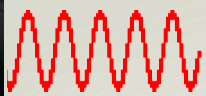


4.0000003 cm



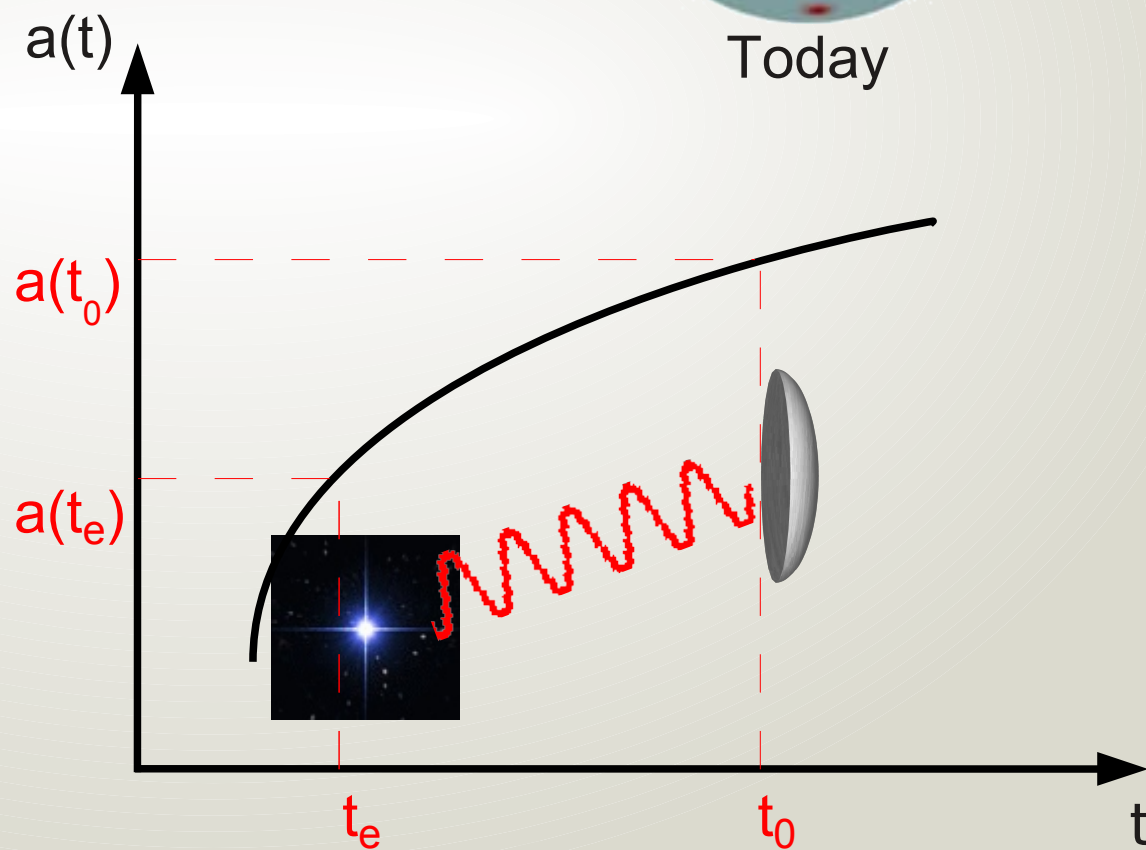
5 years later

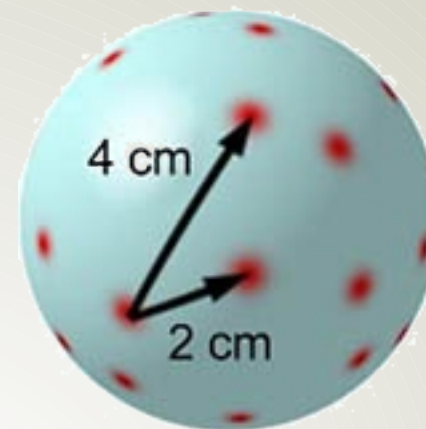
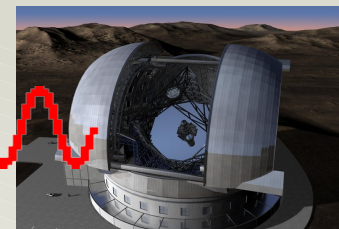
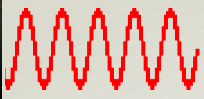
Today + 10 years



5 billion years ago

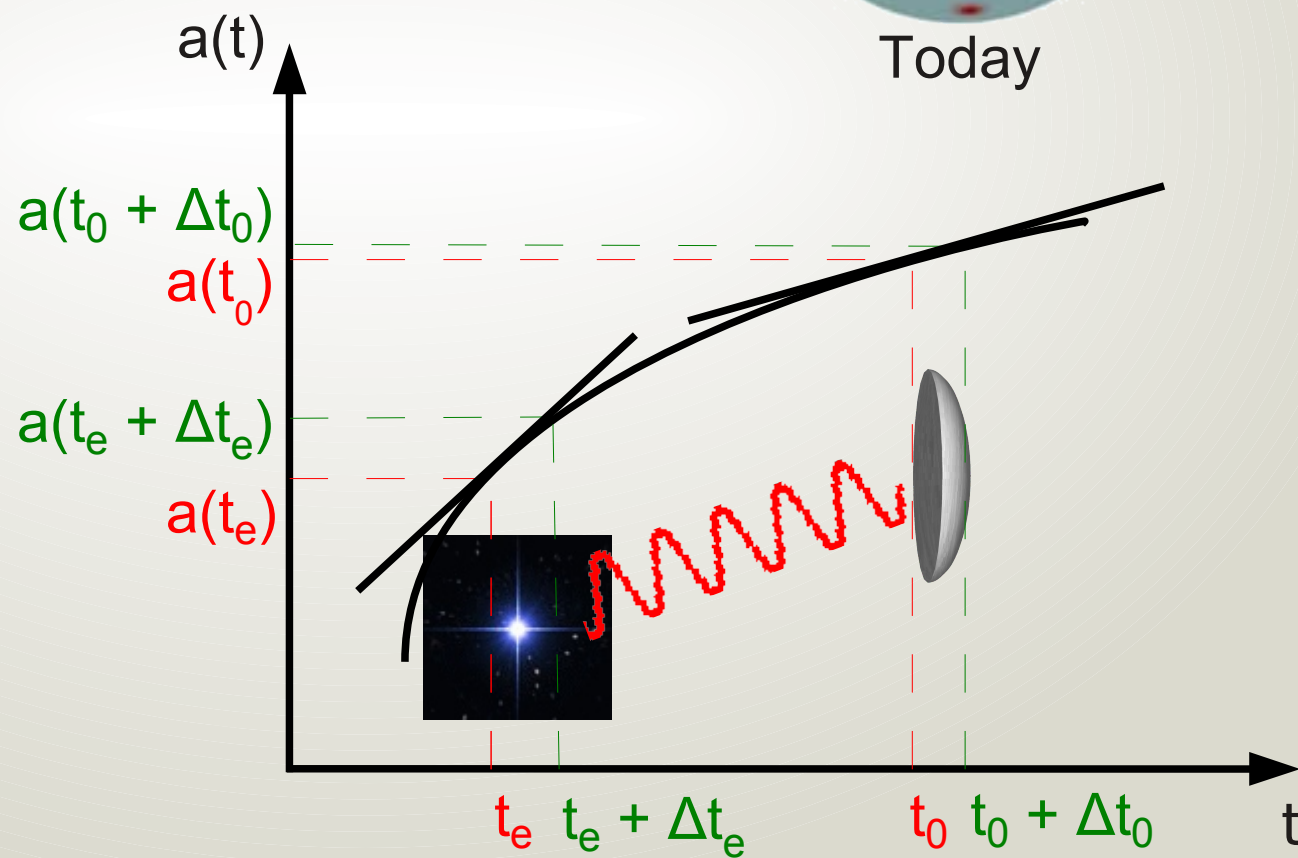
Today





5 billion years ago

Today



What is dz/dt_0 ?

$$1+z = \frac{a(t_0)}{a(t_e)}$$

$$\frac{d}{dt_0} \left[1+z = \frac{a(t_0)}{a(t_e)} \right]$$

$$\begin{aligned} \frac{dz}{dt_0} &= \frac{\dot{a}(t_0)}{a(t_e)} - \frac{a(t_0)}{a(t_e)^2} \dot{a}(t_e) \frac{dt_e}{dt_0} \\ &= (1+z) \frac{\dot{a}(t_0)}{a(t_0)} - (1+z) \frac{\dot{a}(t_e)}{a(t_e)} \frac{1}{1+z} \end{aligned}$$

$$\frac{dz}{dt_0} = (1+z) H_0 - H(z)$$

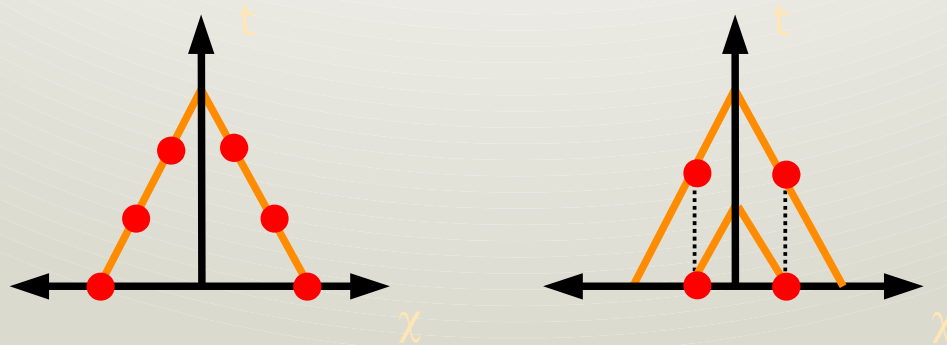
Cosmic Dynamics

The de- or acceleration of the universal expansion rate between epoch z and today causes a small drift in the observed redshift as a function of time:

$$\dot{z} = (1+z)H_0 - H(z)$$

Two remarkable features:

- For this equation to be valid you only need:
 - gravity can be described by a metric theory
 - homogeneity and isotropy
- The redshift drift does not deduce the evolution of the expansion by mapping out our present-day past light-cone but directly measures the evolution by comparing our past light-cones at different times.



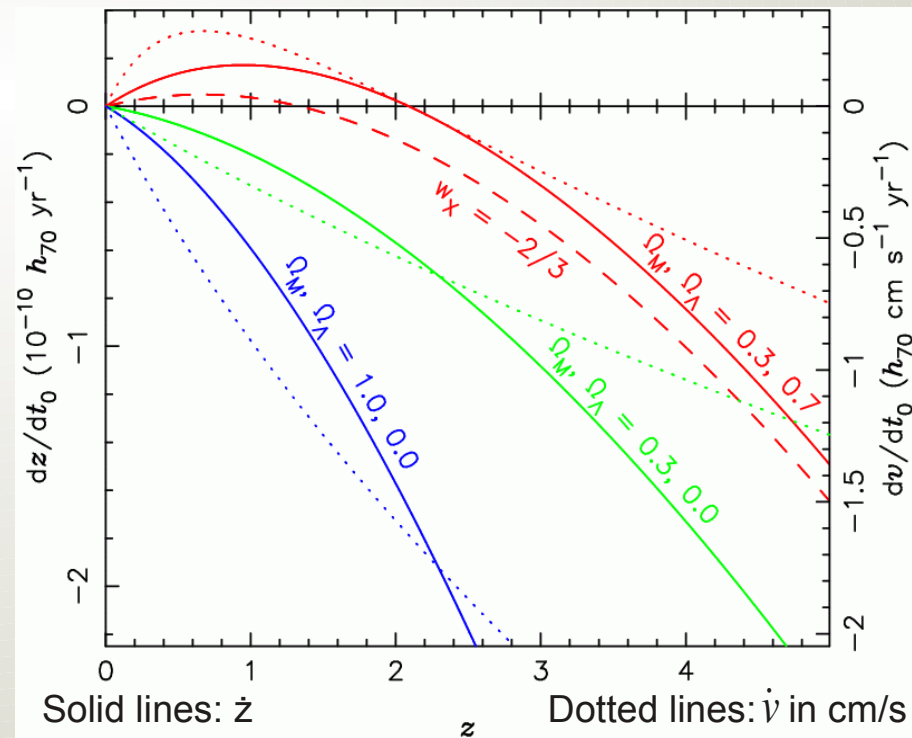
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Measuring $\dot{z}(z)$:

- Allows us to watch, in real time, the universe changing its expansion rate.
- Most direct and model-independent route to the expansion history and acceleration.
- First non-geometric measurement of the global FRW metric.
- Tests whether the geometry and dynamics of spacetime are determined by the 'same' stress-energy tensor.
- Independent confirmation and quantification of accelerated expansion.
- $H(z)$ determination in a redshift range inaccessible to other methods.



Size of the signal

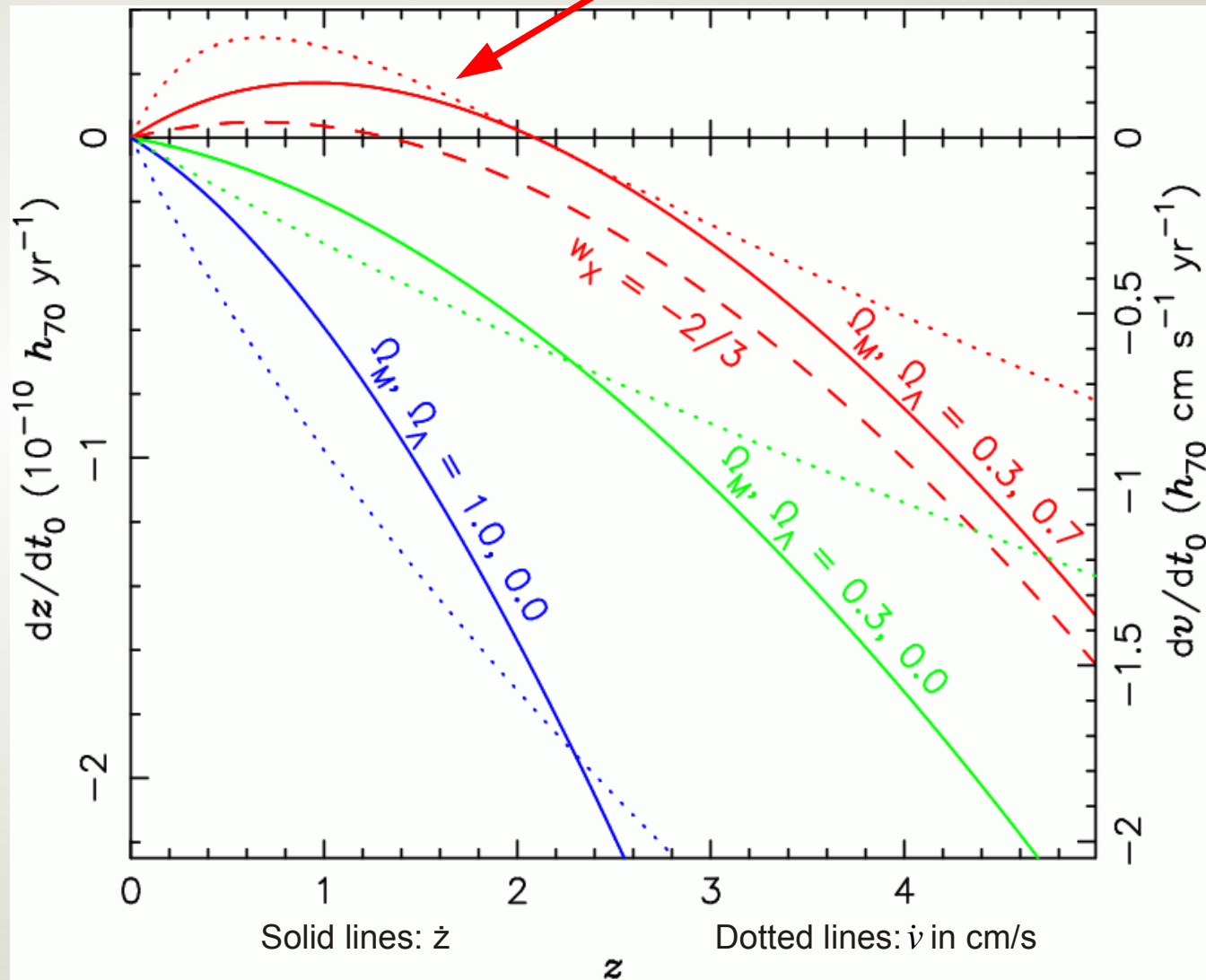
If $\Delta t = 10$ years then:

- $\Delta z \sim 10^{-9}$
- $\Delta \lambda = \lambda_{\text{rest}} \Delta z$
 $\sim 10^{-6} \text{ \AA}$
 $\sim 10^{-4} \text{ pixel}$
 $\sim 1 \text{ nm on CCD}$
- $\Delta v = c \Delta z / (1+z)$
 $\sim 6 \text{ cm/s}$

→ Tiny signal!

BUT: HARPS has already achieved a long-term accuracy of $\sim 1 \text{ m/s}$ with $\sim 10 \text{ cm/s}$ accuracy over a few hours.

Signature of $\Lambda > 0$

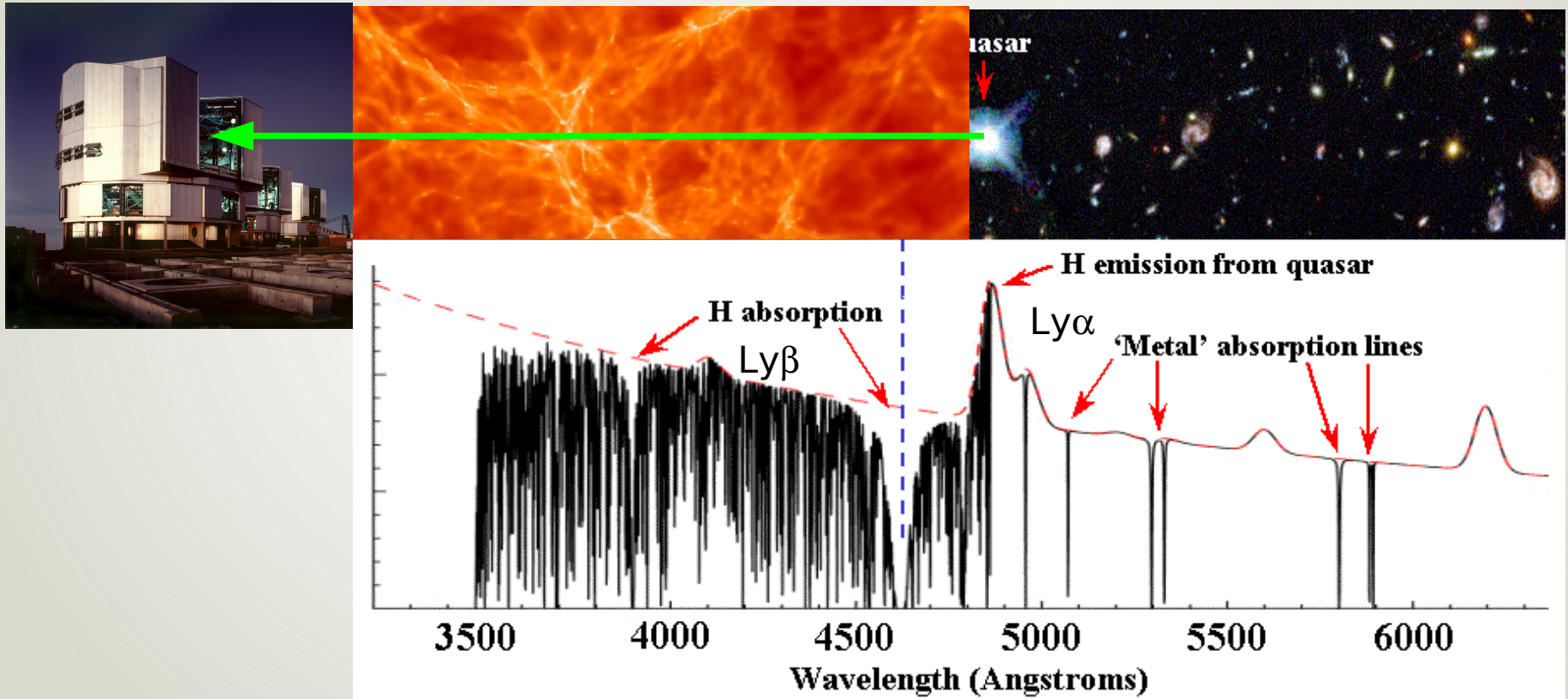


How can we measure the redshift drift?

The precision with which a velocity shift of a spectrum can be determined depends on:

- The number and sharpness of available spectral features.
- The S/N at which they are recorded, i.e.
 - the brightness of the source(s),
 - the size of the telescope,
 - the total system efficiency,
 - the exposure time.

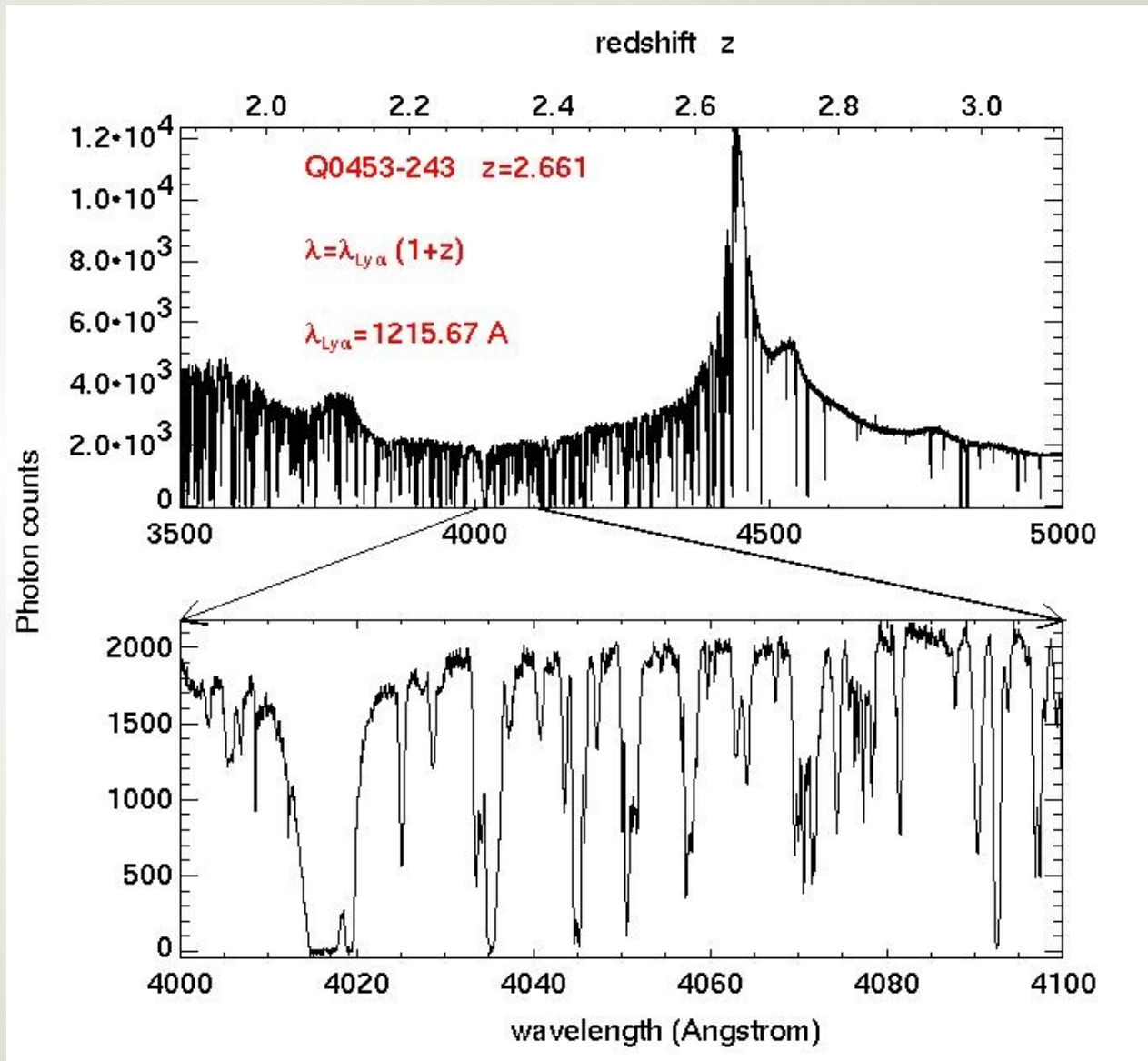
Measuring dz/dt in the IGM



by John Webb

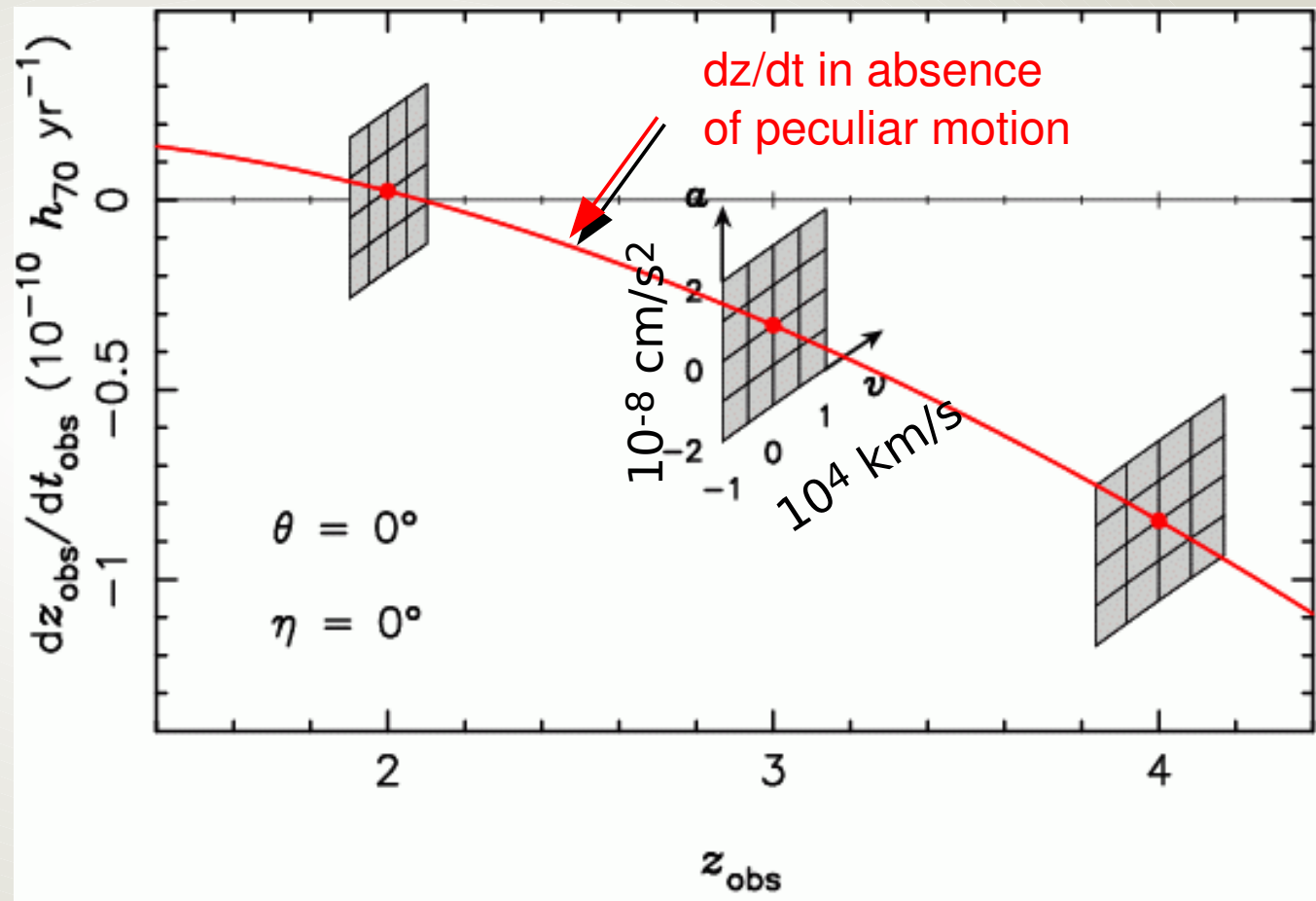
The Lyman α Forest

- ✓ QSOs are the brightest sources at any redshift.
- ✓ QSOs exist over all redshifts, $0 < z < 6$.
- ✓ Each line of sight to a background QSO shows $\sim 10^2$ Ly α lines.
- ✓ The Ly α forest is an excellent tracer of the Hubble flow (small peculiar motions).
- ✗ Line widths are 15-50 km/s. (Metal line widths are of order 1 km/s but reside in deeper potential wells).



Effect of peculiar motion

- The effect of peculiar motion should be compared to the size of the error on an *individual* \dot{z} measurement.
- Peculiar motion is only problematic when using a small number of high-precision measurements.
- No problem when using QSO absorption lines, even if the absorbing gas lies in a deep potential well.

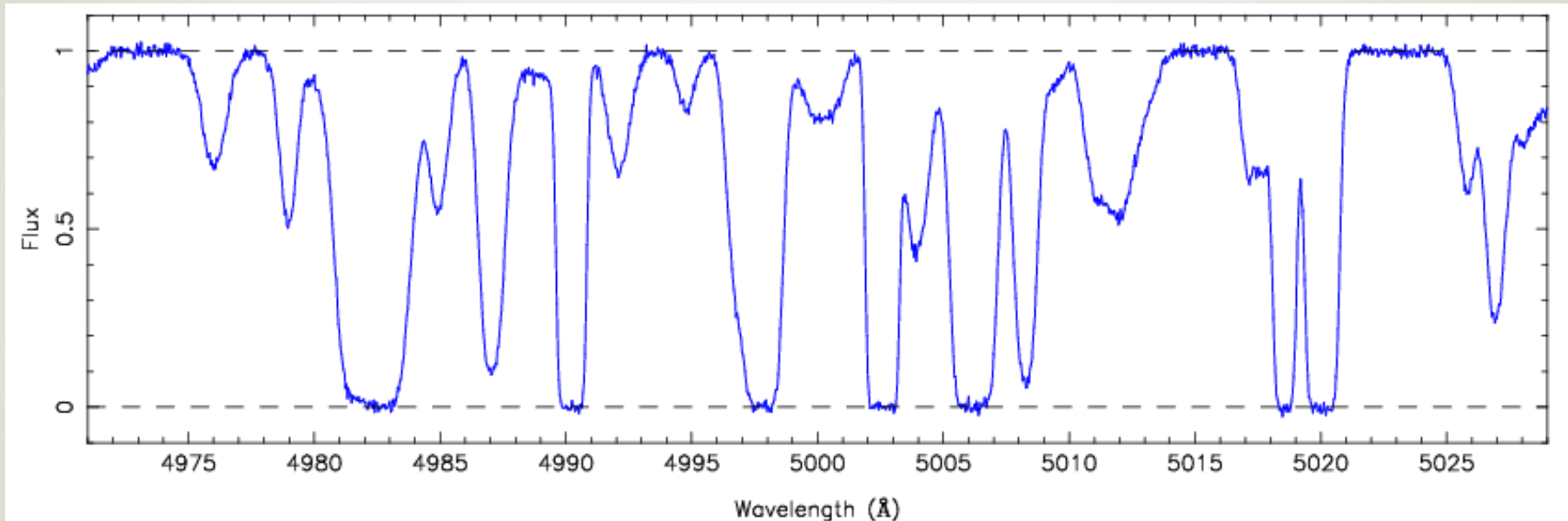


Liske et al. (2008)

→ The Ly α forest traces the Hubble flow!

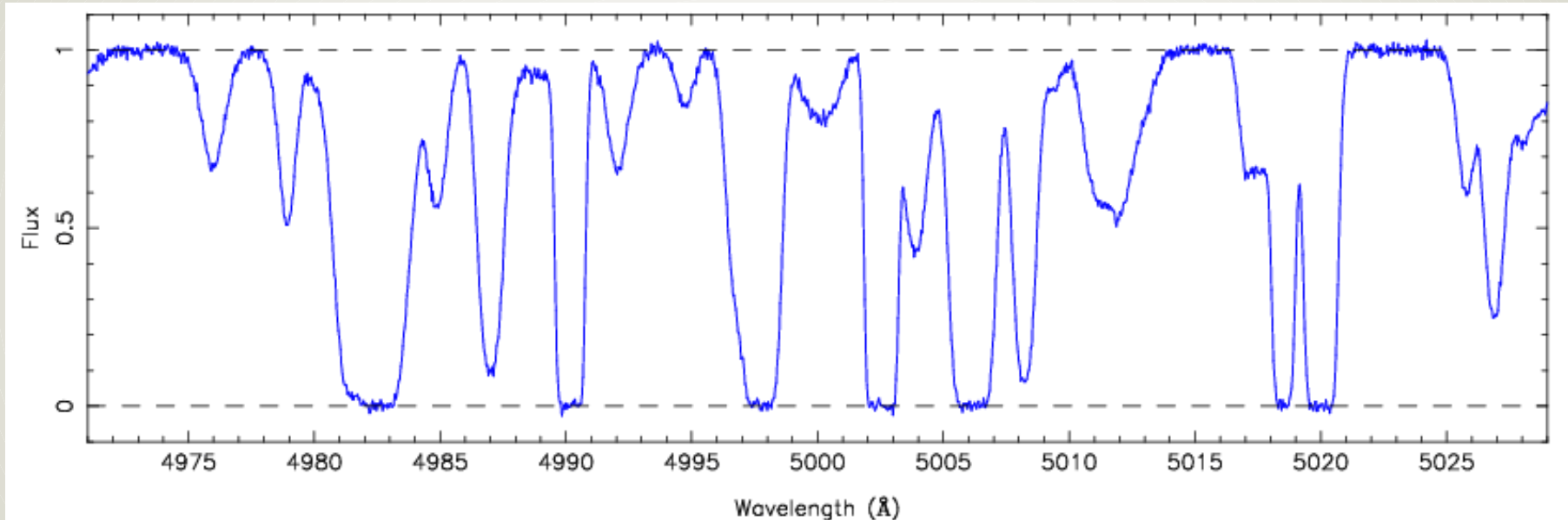
Observing dz/dt in the Ly α Forest

Simulation of the Ly α forest at $z \sim 3$:



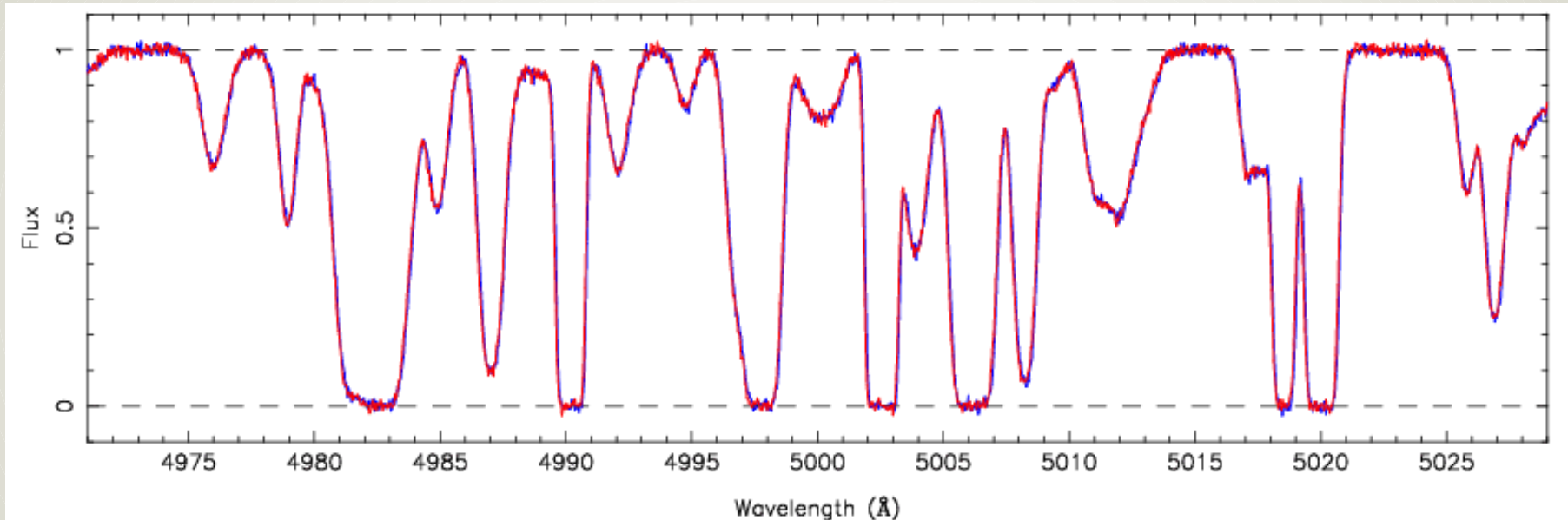
Observing dz/dt in the Ly α Forest

Simulation of the Ly α forest at $z \sim 3$:



Observing dz/dt in the Ly α Forest

Simulation of the Ly α forest at $z \sim 3$:



$$\Delta t = 10^6 \text{ years!}$$

How can we measure the redshift drift?

The precision with which a velocity shift of a spectrum can be determined depends on:

- The number and sharpness of available spectral features.
- The S/N at which they are recorded, i.e.
 - the brightness of the source(s),
 - the size of the telescope,
 - the total system efficiency,
 - the exposure time.



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Can we collect enough photons?

Can we collect enough photons to achieve the required radial velocity accuracy?

QSOs from latest compilations (including SDSS):

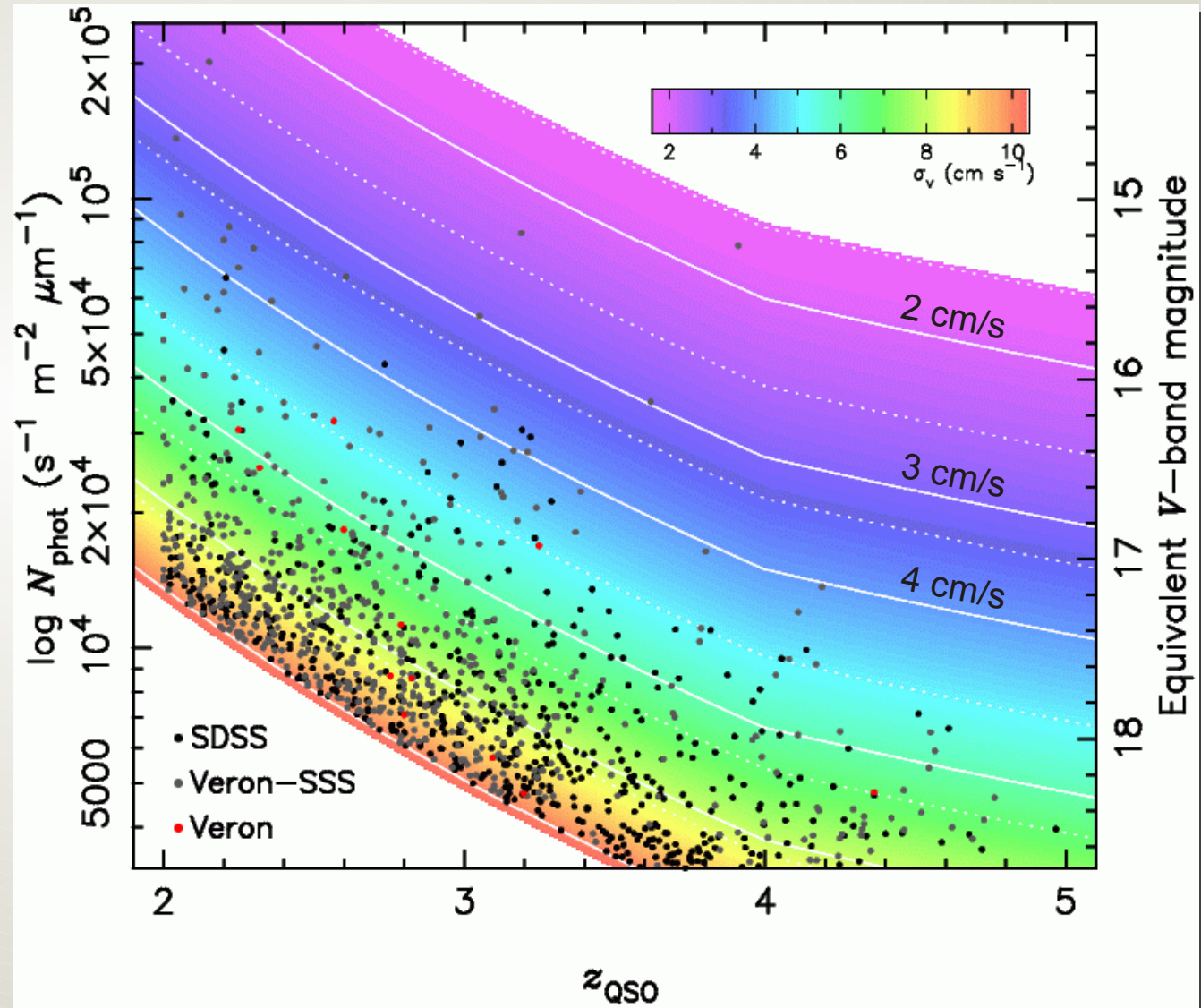
Lines of constant σ_v assume:

$D = 39$ m

efficiency = 25%

$t_{\text{exp}} = 2000$ h

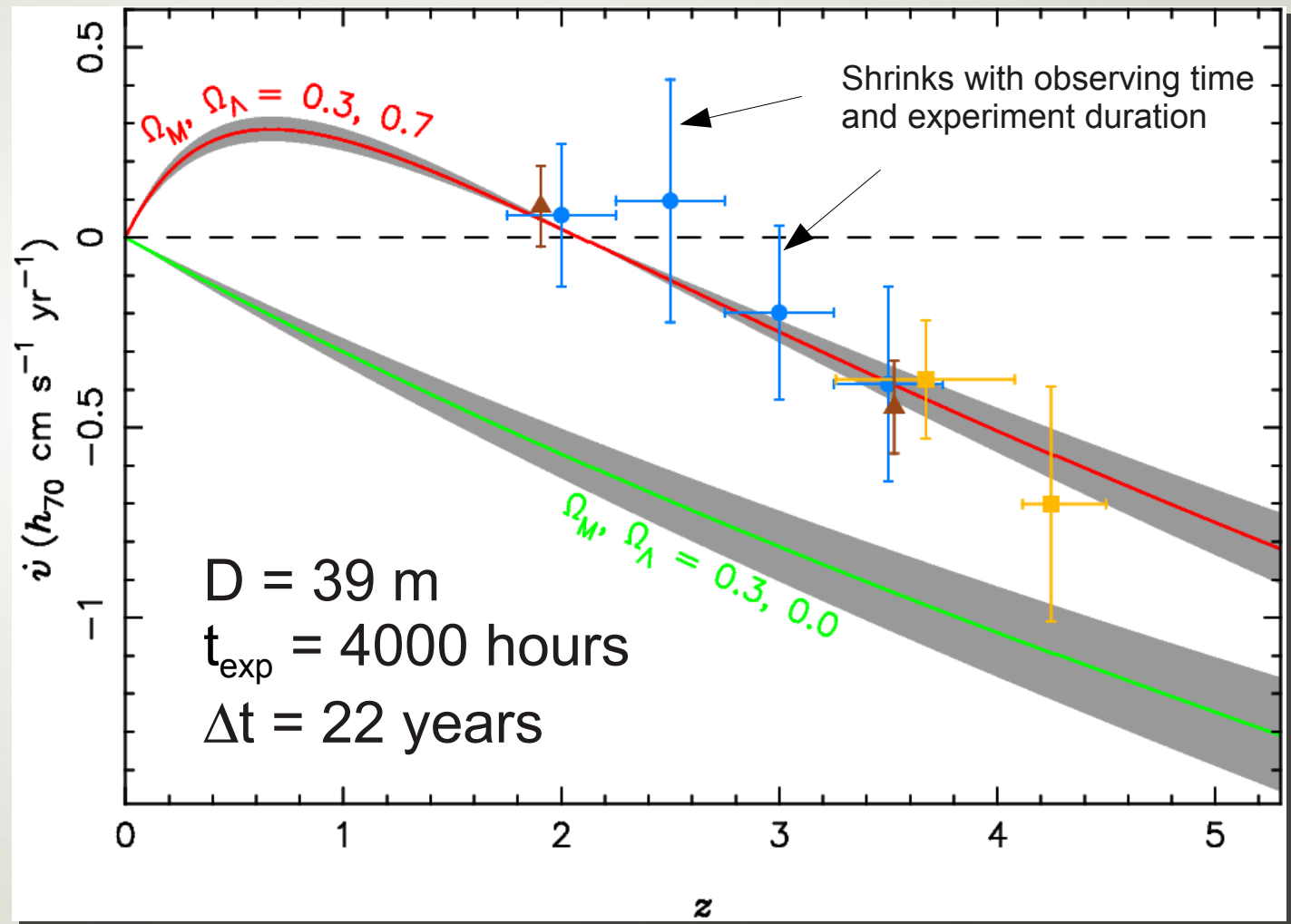
Yes: 18 known QSOs with $2 < z < 5$ are bright enough to achieve a radial velocity accuracy of 4 cm/s using 2000 hours on a 39-m ELT.



Simulation Results

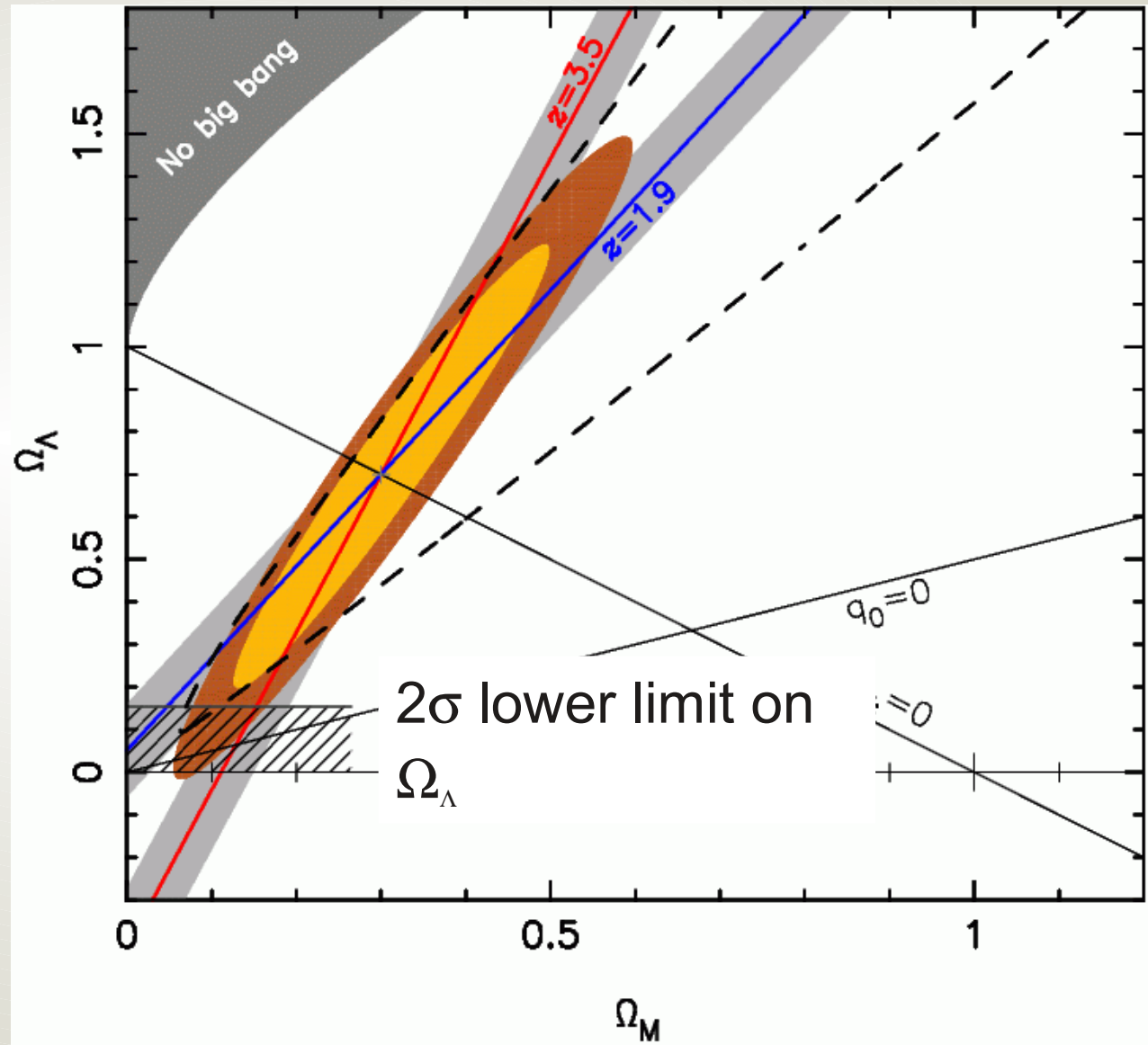
4000 h on a 39-m ELT over 22 years will deliver any *one* of these sets of points.

Different sets correspond to different target selection strategies.



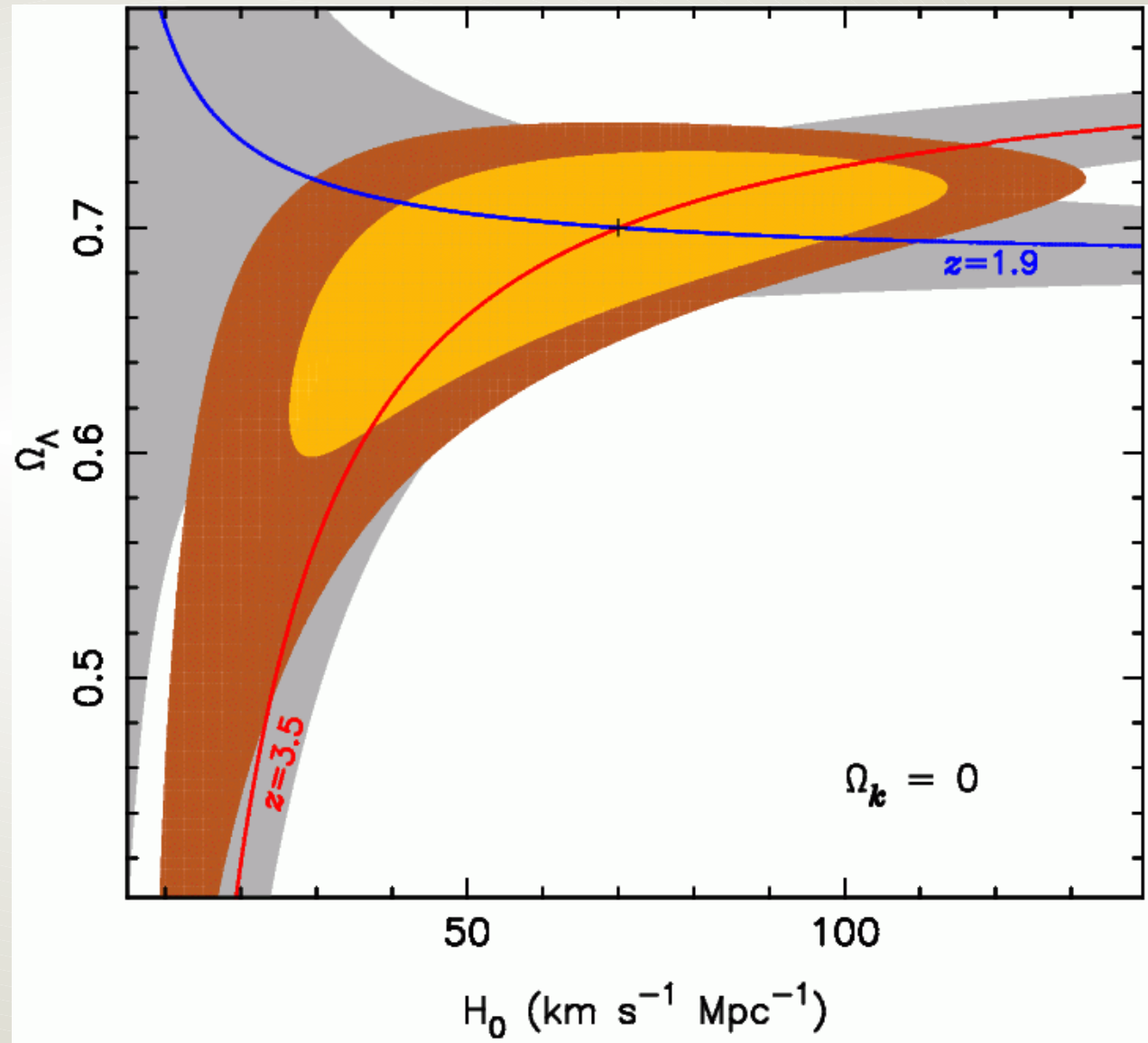
Constraints on Cosmology

- 4000 hours over 22 years will unequivocally prove the existence of dark energy without assuming flatness, using any other cosmological constraints or making any other astrophysical assumption whatsoever.
- Provides independent confirmation of SNIa results, using a different method and a complementary redshift range.



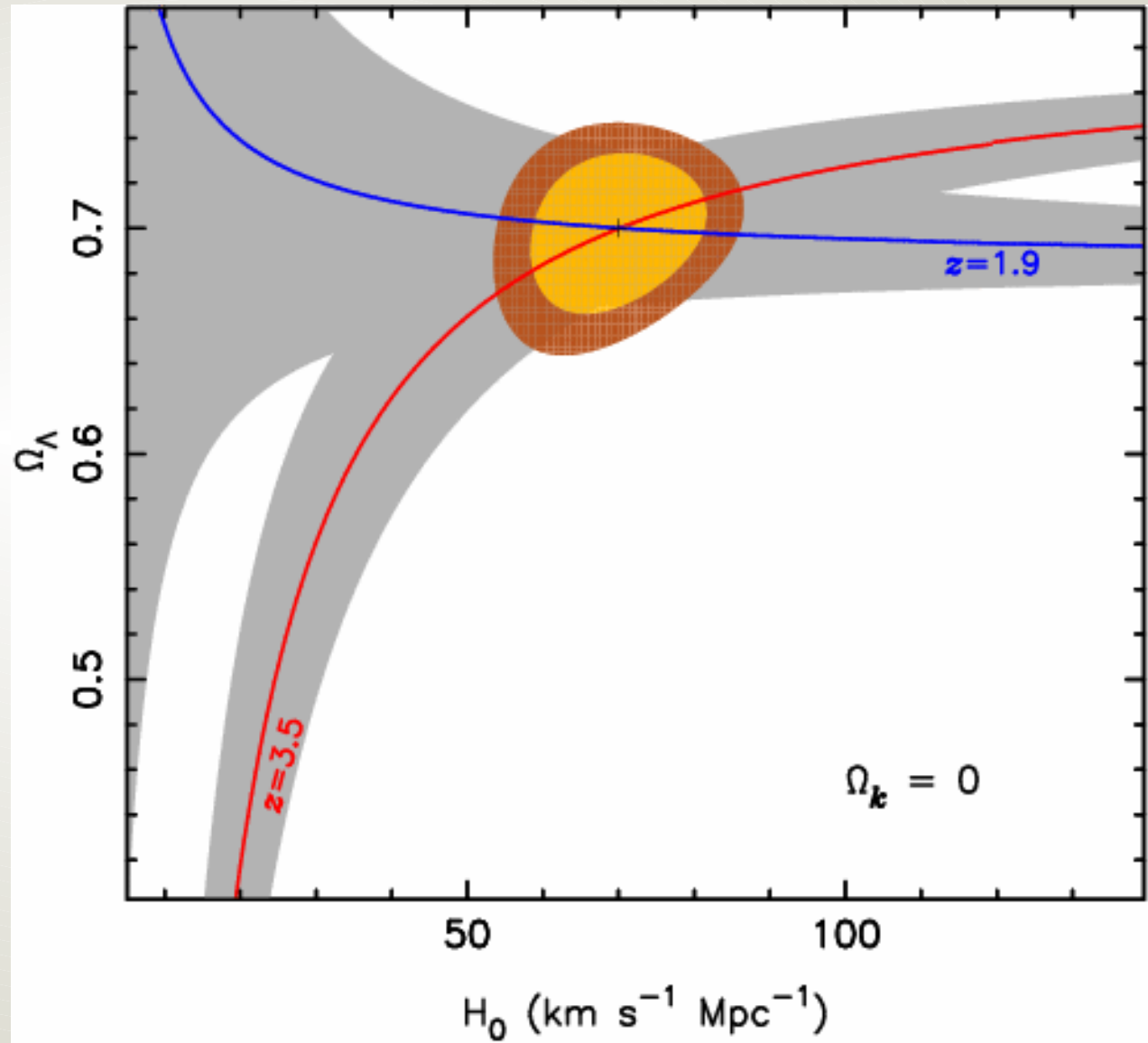
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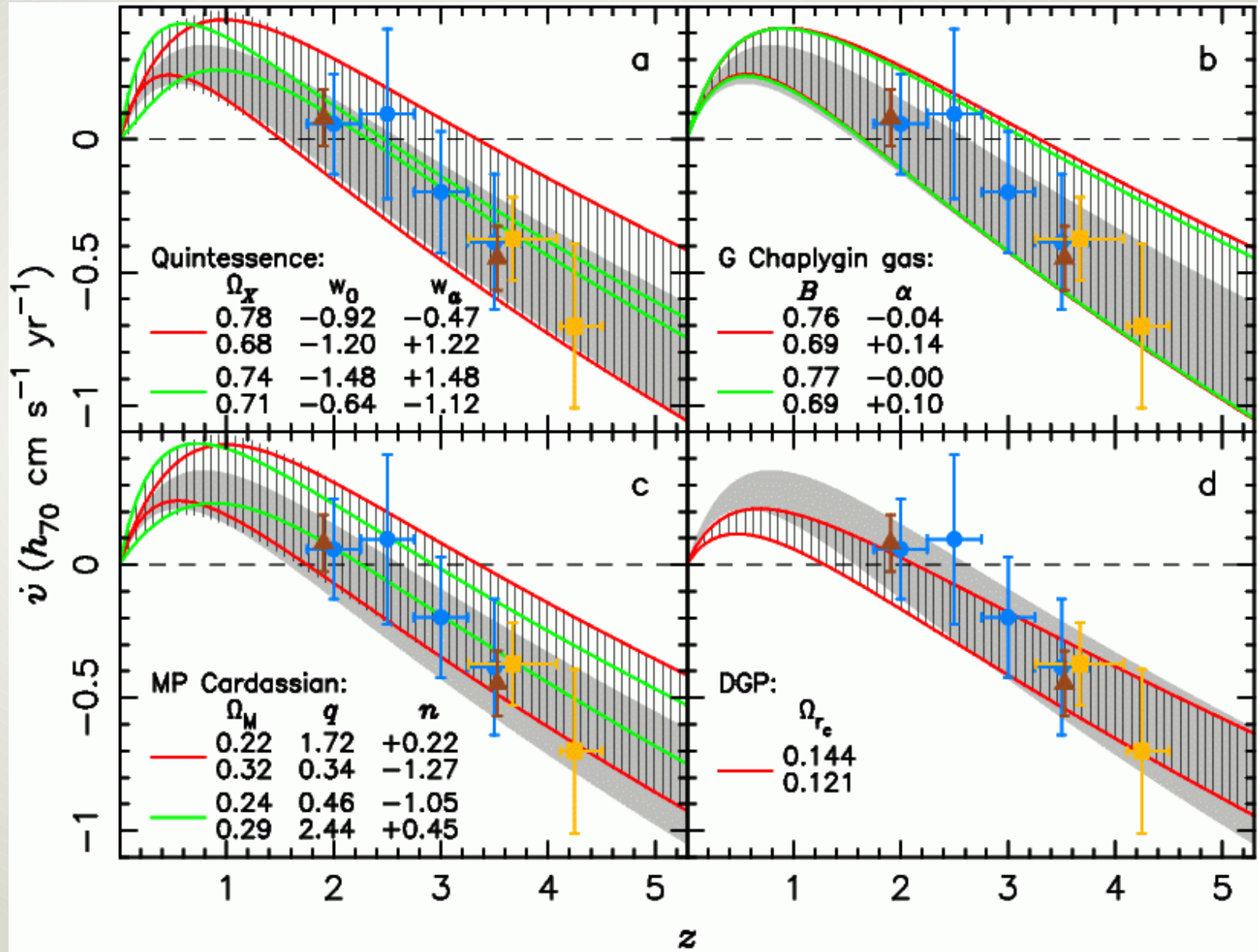
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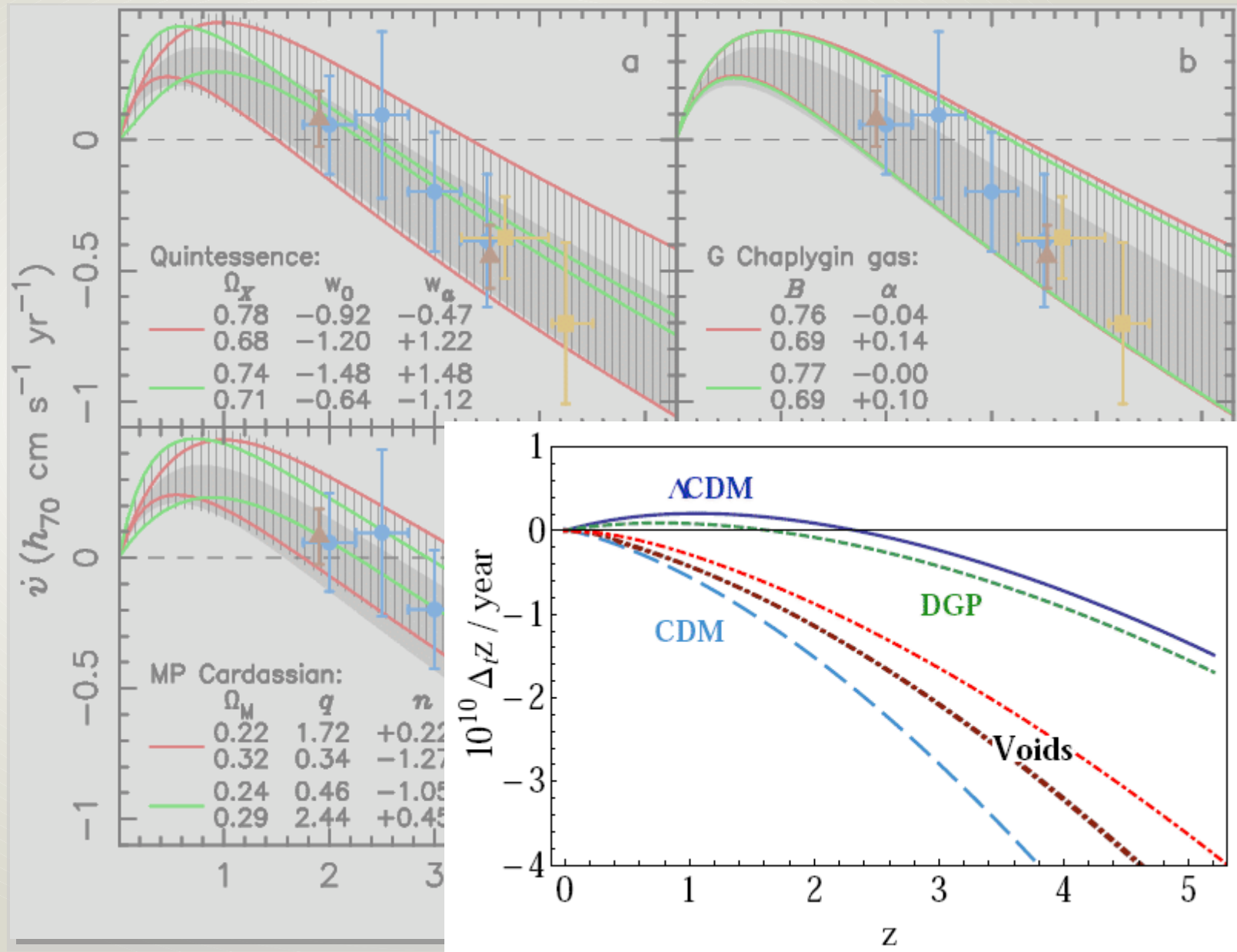
Constraints on non-standard models

Assuming flatness and a fixed H_0 the hashed regions show the allowed dz/dt ranges after the models have been constrained by SNIa, CMB and BAO data (Davis et al. 2007).



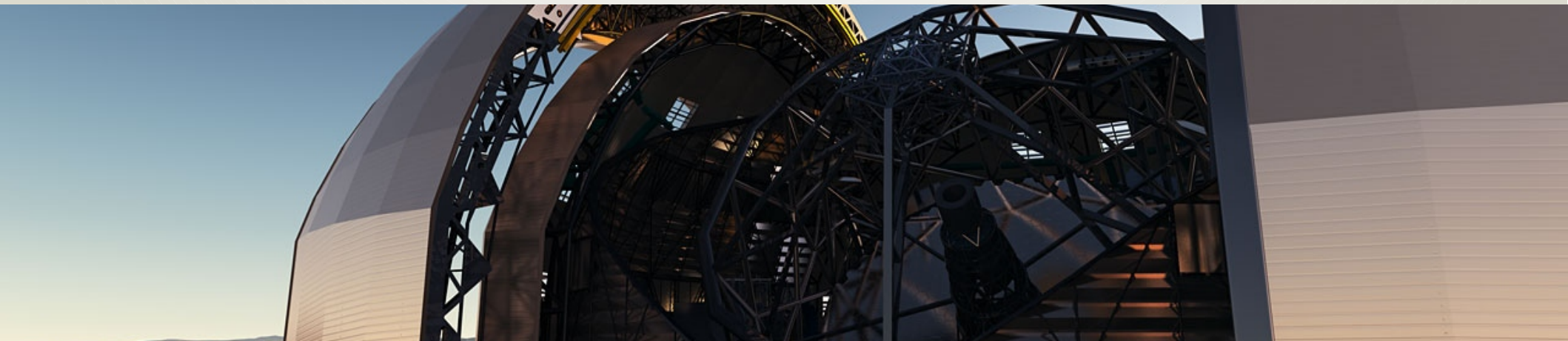
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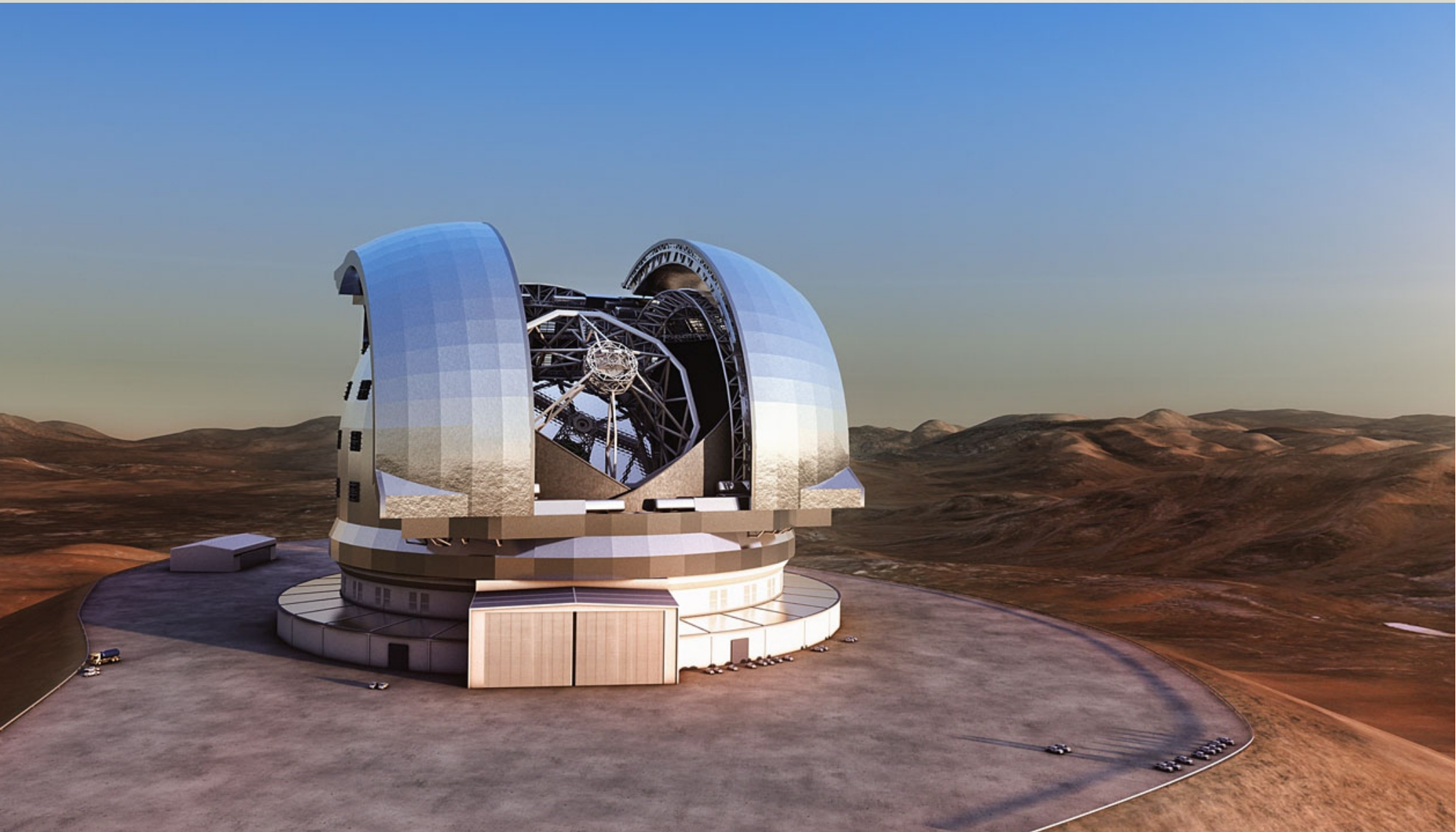


Redshift Drift Summary

- The evolution of the redshift of cosmological sources as a function of time is a direct, dynamical signal of the de/acceleration of the Universe's expansion.
- The E-ELT will offer us the first opportunity to measure the redshift drift (over a timescale of ~ 20 yrs), resulting in a unique measurement of the expansion history:
 - Allows us to watch, in real time, the universe changing its expansion rate.
 - Most direct and model-independent route to the expansion history and acceleration.
 - First non-geometric measurement of the global FRW metric.
 - Requires no priors and is independent of other cosmological experiments.
 - Independent confirmation and quantification of accelerated expansion.
 - $H(z)$ determination in a cosmic epoch inaccessible to other methods.
 - Does not involve or rely on any astrophysics (such as the [unknown] evolution of the sources used).
 - Keeps on giving: signal grows linearly with time \rightarrow very cost effective.



E-ELT
=
**Extremely Exciting Long Term
science**



Is it affordable?

4000 h is an impressive time request for any telescope. However:

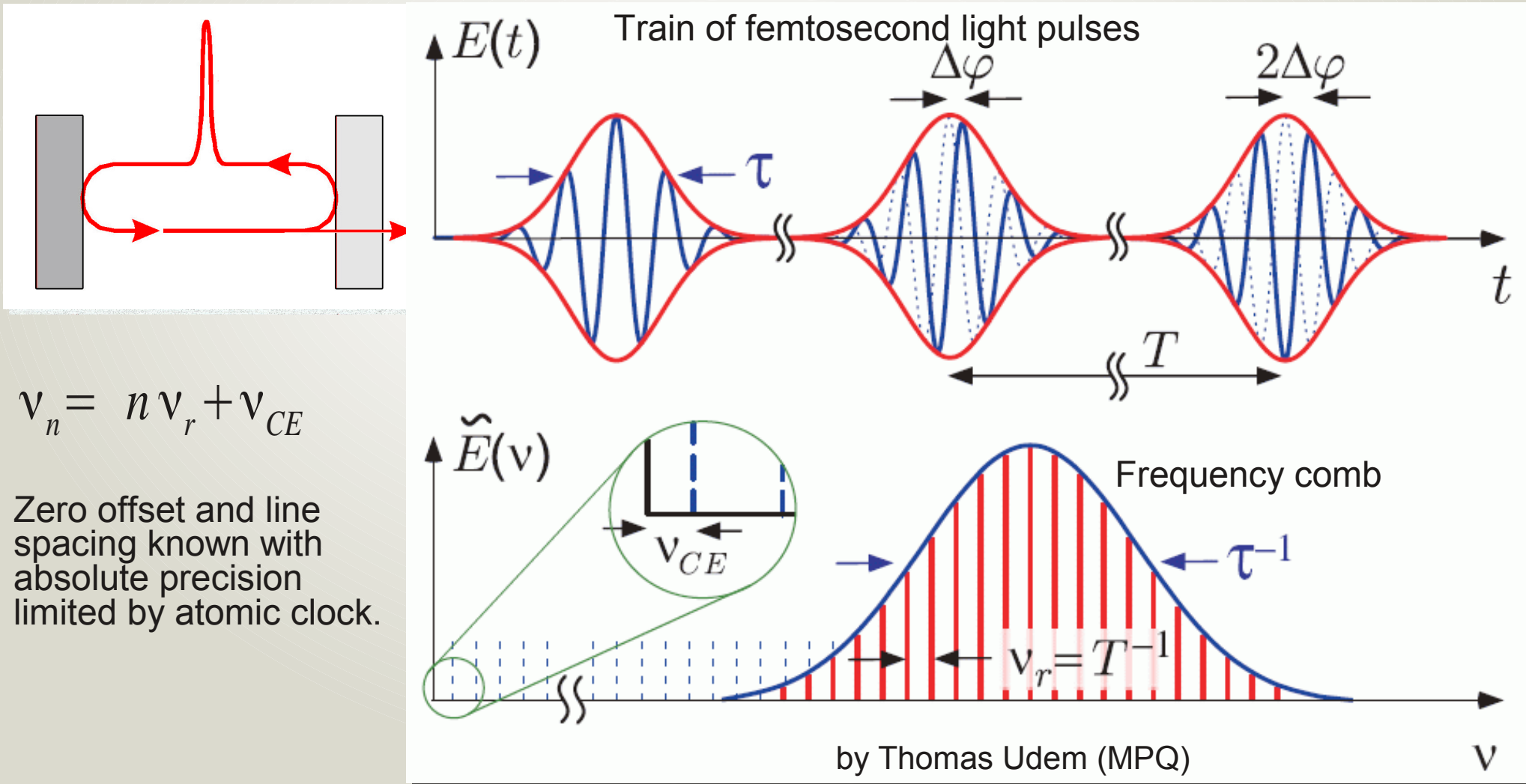
- **The total time is distributable (to some extent)**
4000 h / 20 yr = 20 nights per year
- **Comparable to past investment**
VLT/UVES has invested ~3000 hours on QSO spectroscopy.
- **Synergy with other ELTs**
Assuming appropriate instrumentation, data from all ELTs could be combined.
- **Immediate science with the same data**
 - Cosmological variation of fundamental constants
 - $T_{\text{CMB}}(z)$
 - Primordial deuterium abundance
 - Metallicity evolution of the low-density IGM
 - Tomography of the IGM

Wavelength Calibration

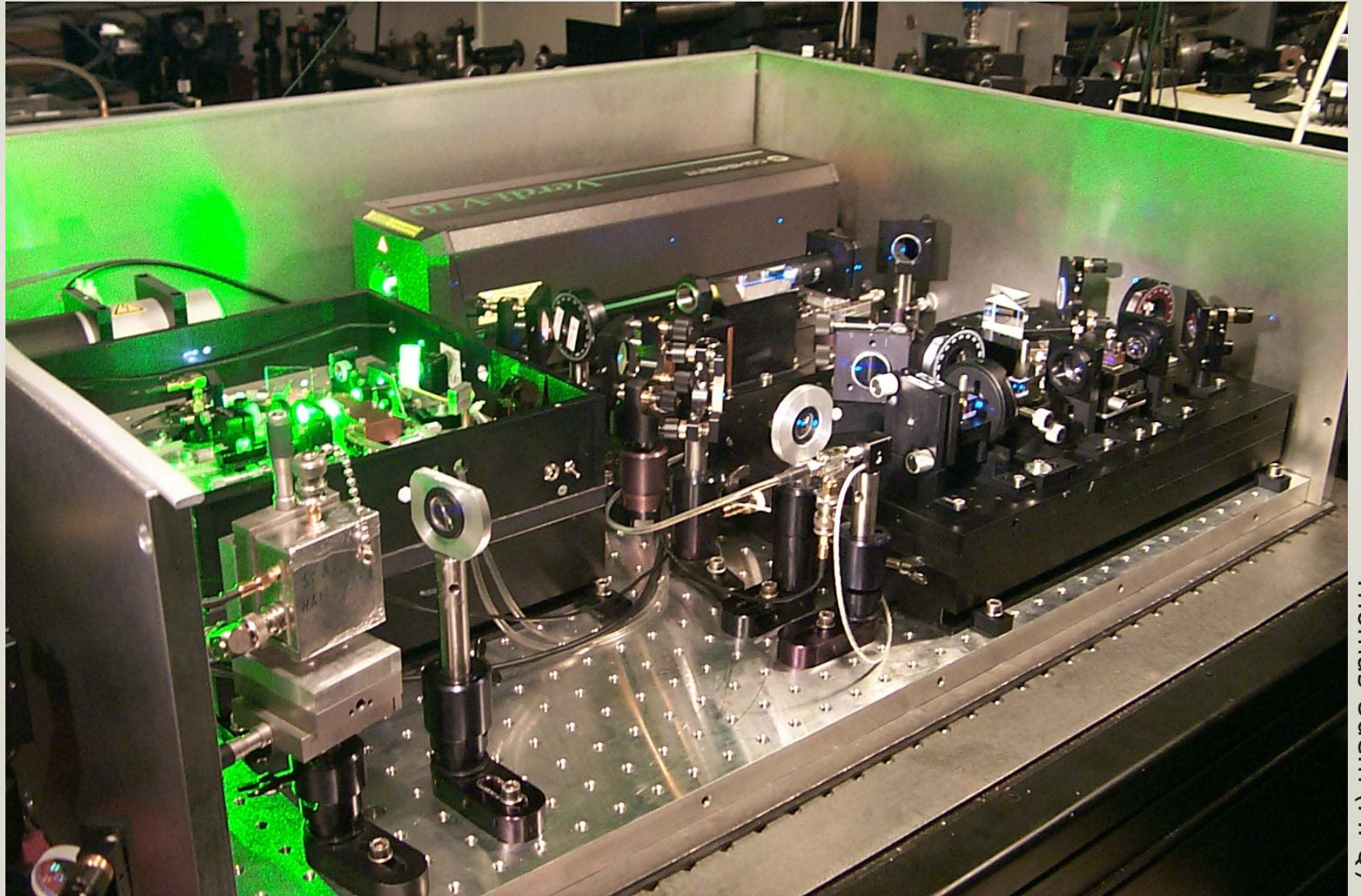
Desired characteristic	ThAr	I ₂ cell	LFC
From fundamental physics	✓	✓	✓
Individually unresolved	Mostly	✓	✓
Resolved from each other	✗	✗	✓
Uniformly spaced	✗	✗	✓
Cover optical range	✓	✗	?
Uniform intensity	✗	✗	?
Long-term stability	✗	?	✓
Maintain object S/N	✓	✗	✓
Exchangeable	✓	✓	✓
Easy to use	✓	✓	?
Reasonably low cost	✓	✓	✓

Laser Frequency Comb

- Optical or NIR laser producing a train of monochromatic femtosecond light pulses.
- Pulse repetition rate is controlled by an atomic clock.
- Produces a spectrum of evenly spaced δ -functions (frequency comb) whose absolute wavelengths are known to a precision limited only by the atomic clock.



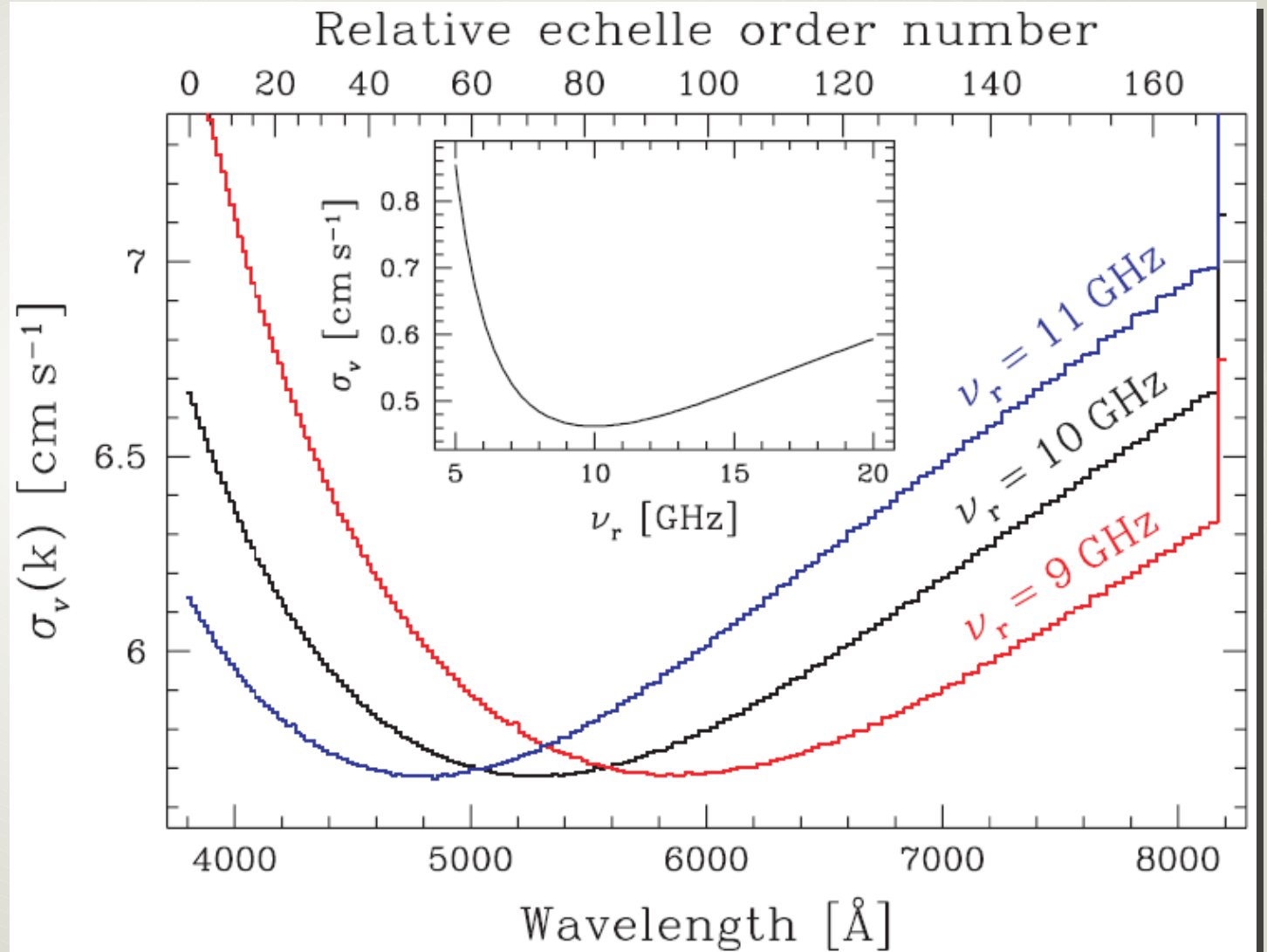
Laser Frequency Comb



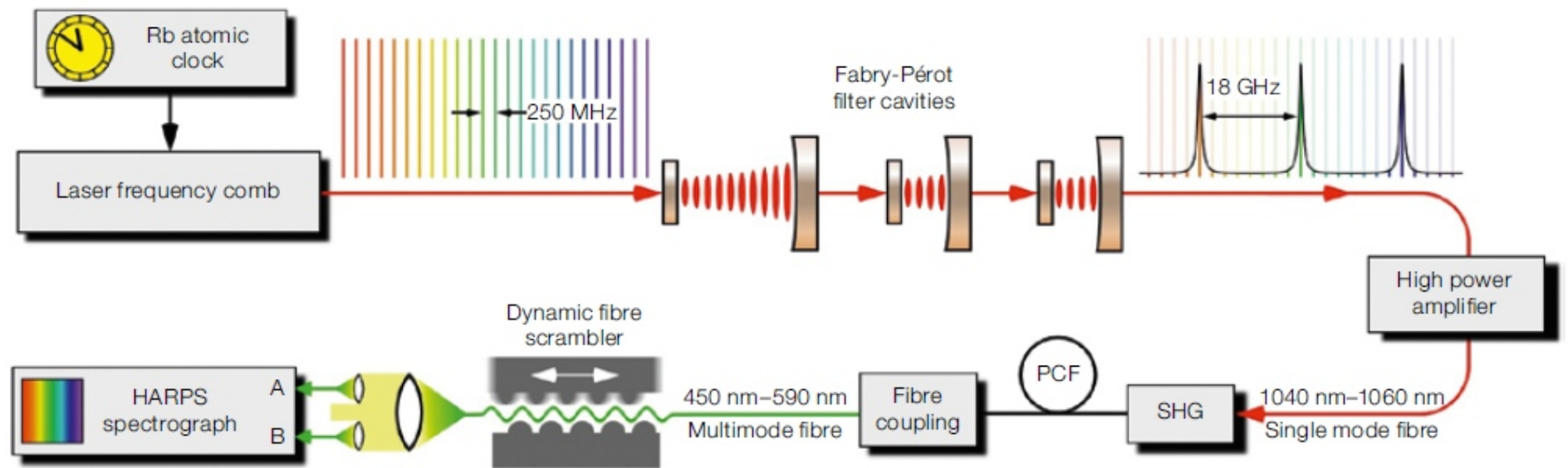
Simulation Results

Photon-limited
wavelength calibration
precision is ~ 0.5 cm/s.

Optimal pulse repetition
rate is 10-20 GHz.

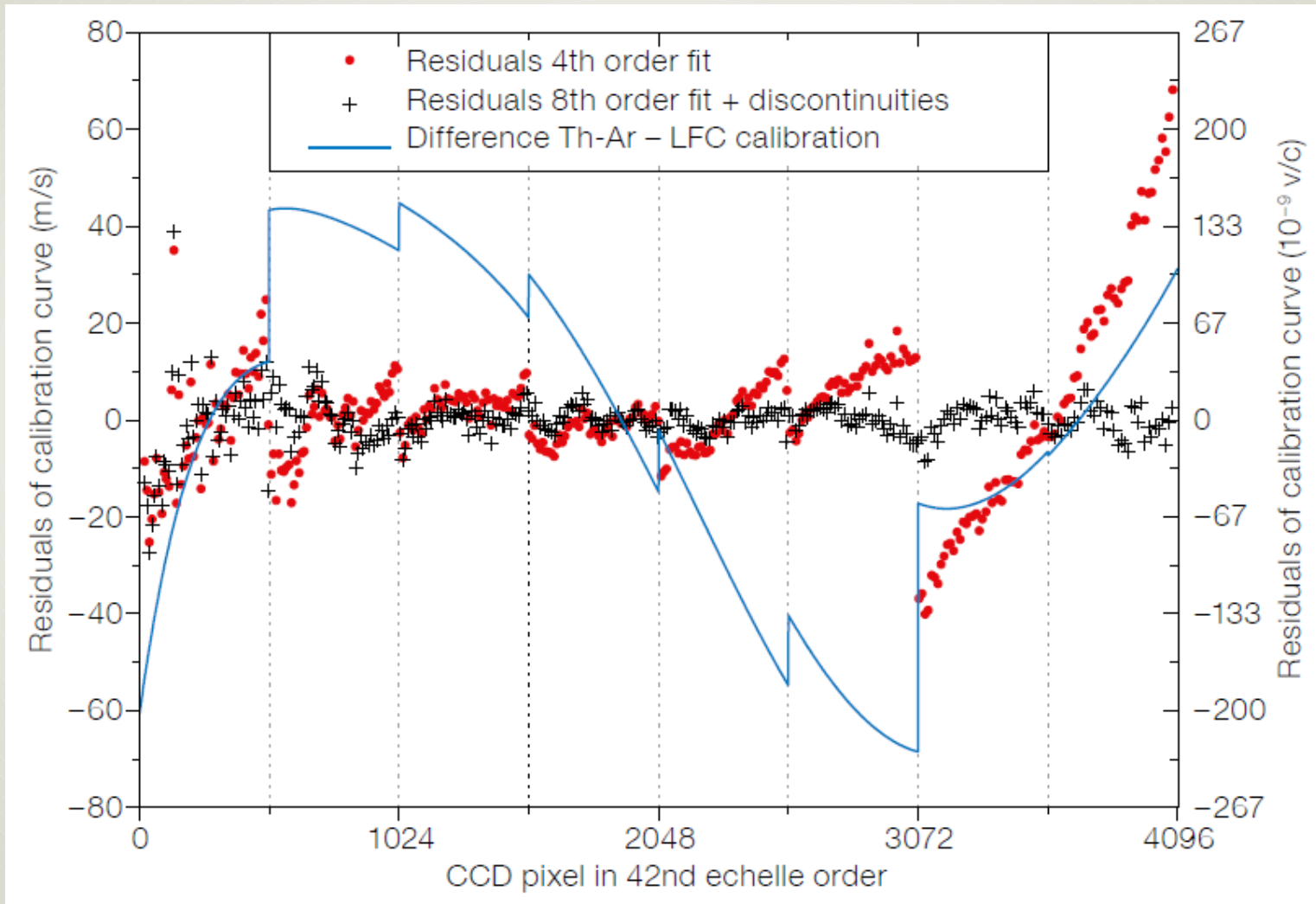


LFC on HARPS @ ESO 3.6 m



Lo Curto et al. (2012)

LFC on HARPS @ ESO 3.6 m



Lo Curto et al. (2012)