On the oscillation spectrum of a magnetized core in a giant star

Michel Rieutord with Sébastien Deheuvels

Institut de Recherche en Astrophysique et Planétologie, France

M. Rieutord On the oscillation spectrum of a magnetized core in a giant star

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Many giant stars observed with Kepler do not show, or weakly show, dipolar l=1, modes.

Fuller, Cantiello, Stello, Garcia and Bildsten (2015) proposed that this "depression" is caused by magnetic fields hidden in the core of these stars.

Two questions :

- The analysis showing that this is possible is local (WKB). Can we confirm this by global models?
- If yes, can we get more information on the magnetic fields of giants, intensity, geometry, origin ?

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FIGURE : From Fuller et al. 2015

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Consider a $1.5M_{\odot}$ MESA-model at the right stage of evolution (thanks Sebastien !)

$$10 \le \rho \le 10^5 \text{g/cm}^3$$
$$\nu \sim 10^2 \text{cm}^2/\text{s} \qquad \eta \sim 10^2 \text{cm}^2/\text{s} \qquad \kappa \sim 10^9 \text{cm}^2/\text{s}$$

$$\mathcal{P} \sim 10^{-7}$$
 $\mathcal{P}_m \sim 1$

Viscous and ohmic dissipation negligible

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- We assume ρ ≃ Cst not very realistic, but necessary for a first step in a tough problem.
- We impose a uniform $\vec{B} = B_0 \vec{e}_z$ magnetic field
- We insert heat absorbers to stably stratify the core

A (10) × (10) × (10)

Sketch of the configuration



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Perturbations $\vec{v}, \vec{b}, \delta p, \dots$ obey

$$\begin{cases} \partial_t \vec{v} = -\vec{\nabla} \delta p / \rho + (\delta \rho / \rho) \vec{g} + v \Delta \vec{v} + (\vec{\nabla} \times \vec{b}) \wedge \vec{B}_0 / \rho \mu_0 \\ \vec{\nabla} \cdot \vec{v} = 0 \\ \vec{\nabla} \cdot \vec{b} = 0 \\ \partial_t \vec{b} = \vec{\nabla} \times (\vec{v} \wedge \vec{B}_0) + \eta \Delta \vec{b} \\ \partial_t \delta T + \vec{v} \cdot \vec{\nabla} T_0 = \kappa \Delta \delta T \end{cases}$$
(1)

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Dimensionless form

The equations of perturbations read :

$$\vec{\nabla} \cdot \vec{u} = \vec{\nabla} \cdot \vec{b} = 0$$

$$\partial_{\tau}\vec{b} = \vec{\nabla} \times (\vec{u} \wedge \vec{e}_z) + \mathbf{E}_{\eta}\Delta\vec{b}$$

$$\partial_{\tau}\theta + ru_r = \mathbf{E}_{\mathbf{t}}\Delta\theta$$
$$\partial_{\tau}\vec{u} = -\vec{\nabla}p + \theta\vec{r} + \mathbf{E}_{\nu}\Delta\vec{u} + \mathcal{A}\left(\vec{\nabla}\times\vec{b}\right)\times\vec{e}_z$$

where we set

$$\mathcal{A} = \left(\frac{V_a}{NR}\right)^2, \qquad E_v = \frac{v}{NR^2}, \quad E_t = \frac{\kappa}{NR^2}, \quad E_\eta = \frac{\eta}{NR^2}$$

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Diffusion :

$$E_{\nu} \sim 10^{-18}, \qquad E_{\eta} \sim 10^{-18}, \qquad E_t \sim 10^{-11}$$

Numerically, we choose

$$E_{\nu} \sim 10^{-10}, \qquad E_{\eta} \sim 10^{-10}, \qquad E_t \sim 10^{-7}$$

Brunt-Väisälä frequency : the maximum value is around $10^4 \mu$ Hz if $B_0 = 2 \times 10^5$ G then $\mathcal{A} = 10^{-8}$. if $B_0 = 2 \times 10^7$ G then $\mathcal{A} = 10^{-4}$.

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Evolution of gravity modes eigenvalues with B_0



Evolution of gravity modes When perturbed by \vec{B}_0



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We impose that the excited acoustic modes shake the core-envelope interface at their frequency. Some energy is dissipated in the core. How does it vary with the intensity of the magnetic field ?

Answer from the Boussinesq model



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propagation diagram



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- Alfvén modes are high frequency modes, Gravity modes are low frequency modes
- When they meet/interact their wavelength is the largest possible, thus magneto-gravity modes may not be very good at dissipating energy.
- The Boussinesq model shows that indeed damping is reduced and frequency shifted...
- The field might not be at the largest scale (the radius) but at the scales left by ohmic diffusion. In 700Myrs, all scales below 0.035R_c are erased. However, in dynamo generated field large scales are dominant ...
- The mechanism which weakens the $\ell = 1$ -modes of the giants need further investigations : \vec{B} -fields may have a more subtle effect than just being absorbers.
- A real difficulty : damping the modes without touching their frequency...

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Thus, magnetic fields may do the job, may be not

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Se non è vero, è ben trovato !